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3C 273 WITH NuSTAR: UNVEILING THE ACTIVE GALACTIC NUCLEUS

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ABSTRACT

We present results from a 244 ks NuSTAR observation of 3C 273 obtained during a cross-calibration campaign with the Chandra, INTEGRAL, Suzaku, Swift, and XMM-Newton observatories. We show that the spectrum, when fit with a power-law model using data from all observatories except INTEGRAL over the 1–78 keV band, leaves significant residuals in the NuSTAR data between 30 and 78 keV. The NuSTAR 3–78 keV spectrum is well described by an exponentially cutoff power law ($\Gamma = 1.646 \pm 0.006$, $E_{\text{cutoff}} = 202^{+31}_{-34}$ keV) with a weak reflection component from cold, dense material. There is also evidence for a weak (EW = 23 ± 11 eV) neutral iron line. We interpret these features as arising from coronal emission plus reflection off an accretion disk or distant material. Beyond 80 keV INTEGRAL data show clear excess flux relative to an extrapolation of the active galactic nucleus model fit to NuSTAR. This high-energy power law is consistent with the presence of a beamed jet, which begins to dominate over emission from the inner accretion flow at 30–40 keV. Modeling the jet locally (in the NuSTAR + INTEGRAL band) as a power law, we find that the coronal component is fit by $\Gamma_{\text{AGN}} = 1.638 \pm 0.045$, $E_{\text{cutoff}} = 47 \pm 15$ keV, and jet photon index by $\Gamma_{\text{jet}} = 1.05 \pm 0.4$. We also consider Fermi/LAT observations of 3C 273, and here the broadband spectrum of the jet can be described by a log-parabolic model, peaking at ~2 MeV. Finally, we investigate the spectral variability in the NuSTAR band and find an inverse correlation between flux and $\Gamma$.

Key words: quasars: individual (3C 273) – X-rays: individual (3C 273)

1. INTRODUCTION

At a redshift of $z = 0.158$ (Schmidt 1963), 3C 273 is the nearest high-luminosity quasar and has been extensively studied at all wavelengths since its discovery in 1963 (for a review, see Courvoisier 1998). It is radio-loud and highly variable across nearly all energies (Soldi et al. 2008), with a jet showing apparent superluminal motion.

At radio to millimeter and at $\gamma$-ray energies, flares from the relativistic jet dominate the variability of 3C 273 (Türler et al. 2000; Abdo et al. 2010). In the optical–UV band there is a bright excess (blue hump) that is spectrally complicated and suggestive of two independently varying components: a more rapidly varying component originating from thermal reprocessing from an accretion disk, and a slower component that could be synchrotron emission unrelated to the radio–millimeter jet (Paltani et al. 1998). As observed in many other active galactic nuclei (AGNs), there is a soft excess in the low-energy ($<2$ keV) X-ray band possibly due to Comptonized UV photons (Page et al. 2004). Correlations between the UV and low-energy X-rays (Walter & Courvoisier 1992; Page et al. 2004) have been noted in the past, supporting this interpretation. However, recent observations have failed to detect correlated variability (Chernyakova et al. 2007; Soldi et al. 2008), leaving the interpretation of the optical/UV excess uncertain.

There is evidence of an intermittently weak iron line in the X-ray spectrum, which appears to be broad ($\sigma \sim 0.6$ keV, EW $\sim 20–60$ eV), occasionally neutral (Turner et al. 1990; Grandi & Palumbo 2004; Page et al. 2004), and sometimes ionized (Yaqoob & Serlemitsos 2000; Kataoka et al. 2002). The line is too faint to reliably trace its variability as a function of flux with current instrumentation.

Above 2 keV and up to MeV energies, previous observations report a hard power-law spectrum, as is common for jet-dominated AGNs (blazars). Over the 30 yr that 3C 273 has been reliably monitored, there appears to be a long-term spectral evolution underlying the short-term variations. The source was in its softest observed state in 2003 June (photon index $\Gamma \sim 1.82 \pm 0.01$), a value of $\Delta \Gamma \sim 0.3–0.4$ above what was measured in the 1980s ($\Gamma \sim 1.5$), and since then the source has hardened again to a value of $\Gamma \sim 1.6–1.7$ (Chernyakova et al. 2007).
Finally, 3C 273 is a strong γ-ray emitter; it has been
to routinely detected by Fermi, showing several interesting γ-ray
flares (Abdo et al. 2010). Broadband observations including
Fermi and INTEGRAL reported by Esposito et al. (2015) imply
that the broadband spectrum of the jet emission is not a single
power law, but instead it can be described by a log-parabola
model peaking around 3–8 MeV.

Because of the presence of the soft excess and the iron line, it
has been postulated that the X-ray emission arises from a mix
of thermal and nonthermal processes: an AGN component from
the inner accretion flow (i.e., emission from the corona possibly
accompanied by disk reflection) and a jet component. Grandi &
Palumbo (2004) fitted the broadband spectrum from BeppoSAX
with a cold reflector irradiated by an isotropic coronal X-ray
source for the AGN component (Magdziarz & Zdziarski 1995)
and a power law for the jet component. They were able to
obtain reasonable fits to a restricted model where the coronal
continuum photon index and cutoff energy were held
constant, excluding the thermal disk photons by a corona has
no convincing correlation between the variations above and
below 20 keV. As of yet no clear signature above 10 keV of a
disk-related component in the form of a refection feature and/or
upscattering of the thermal disk photons by a corona has
been found. We show that this is no longer the case and present
here overlapping contemporaneous observations of 3C 273
obtained by the six observatories: Chandra, INTEGRAL,
NuSTAR, Suzaku, Swift, and XMM-Newton. We place particular
emphasis on the NuSTAR observation, which was about six
times longer than the rest and for the first time allows a study of
the 3–78 keV region without cross-calibration concerns.

2. DATA ANALYSIS

On UT 2012 July 17, the six observatories Chandra
(Weisskopf et al. 2002), INTEGRAL (Winkler et al. 2003),
NuSTAR (Harrison et al. 2013), Suzaku (Mitsuda et al. 2007),
Swift (Gehrels et al. 2004), and XMM-Newton (Jansen et al.
2001) observed 3C 273 as part of a cross-calibration
campaign organized by the International Astronomical
Consortium for High Energy Calibration (IACHEC). Figure 1
shows the overlapping GTI times and length of each
observation, and Table 1 lists the observation IDs and exposure
times for the respective observatories.

For Chandra we used CIAO 4.6.1 and CALDB
4.6.1.1. The data were taken with gratings configuration ACIS+HETG and
reprocessed using the CIAO chandra_repro reprocessing
script. We combined orders 1–3 for the HEG and MEG arm
separately and binned the data at 30 counts.

For INTEGRAL we used the standard OSA 10.0 data
reduction package. We only used data from IBIS/ISGRI
(Lebrun et al. 2003; Ubertini et al. 2003) as data from JEM-X
and SPI did not allow us to constrain the spectral shape

http://web.mit.edu/iACHEC/

significantly. We extracted the IBIS data on a science window
(ScW) by ScW basis using our own scripts, which are based on
the IBIS Analysis User Manual procedure. We combined the
spectra using spe_pick according to observation date.

For NuSTAR we used HEASoft 6.15.1 and CALDB
20131223. The data were processed with all standard settings
and source counts extracted from a 30″ radius circular region.
Background was taken from the same detector.

For Suzaku we used CALDB: HXD (20110913), XIS
(20140203), and XRT (20110630). The observation was taken
in 1/4 window mode, and we used 100″ radius circular regions
for the FI detectors (XIS0, 3) and a 140″ region for the BI
detector (XIS1), such that the regions were centered on the
source but were restricted to the operational portions of the
detectors.

For Swift we used HEASoft 6.15.1 and XRT CALDB 2014
February 04. The data were taken in “PHOTON” mode and
were reduced using xrtpipeline. Spectra were extracted
from an annulus region with inner radius of 5″ and outer radius
of 30″ to correct for pileup. The two observations were
combined and the spectra binned at 50 counts.

For XMM-Newton we used SAS v. 13.5.0 with CALDB
2014 January 31. The data were taken in “Small Window”
mode, and to correct for pileup in the MOS, we excised counts
from an annulus region with inner radius of 15″ and outer
radius of 45″. For the PN we extracted from a circular region
of radius 45″.

We analyzed the high-energy γ-ray data from the Fermi
Large Area Telescope (LAT), using the software package
ScienceTools v100p5 and the latest calibration standard
Pass 8. Since 3C 273 was found in a relatively low γ-ray state, this
required integrating the γ-ray flux over a timescale of 50 days
(MJD 56,100–56,150), which is significantly longer than the
6-day NuSTAR campaign. We used the instrument response
function P8R2_SOURCE_V6 (front and back), including
galactic diffuse emission model gll_iem_v06, isotropic back-
ground model iso_P8R2_SOURCE_V6_v06, and background
point sources within 15° from 3C 273 taken from the 2FGL
catalog. Events were extracted from a region of interest within
10° from 3C 273, and we applied standard selection cuts for the
South Atlantic Anomaly (SAA) avoidance, and the zenith angle
cut z < 100°.

Figure 1. Length and coverage of the six overlapping 3C 273 observations. We
do not enforce strict simultaneity between the low Earth orbit observatories
NuSTAR, Suzaku, and Swift, and we extract the full range for all observatories
—except for INTEGRAL—between the Suzaku start and stop times as indicated
by the shaded region in Figure 2.
3. SPECTRAL FITTING

For the entire data analysis we use the XSPEC version 12.8.2 analysis software (Arnaud 1996) and present 90% confidence limits unless otherwise stated.

3.1. Broadband Fitting

To investigate the overall broadband spectrum, we fit data to all observatories except INTEGRAL (to be discussed in the next section) for an interval with reasonable temporal overlap. Figure 1 shows the observation times for the various missions. We chose the interval spanning the Suzaku observation for extraction of data since there is reasonable overlap of part of this interval with all observatories except INTEGRAL. Owing to the relative phasing of the SAA passages and occultation periods among the low Earth orbit observatories (NuSTAR, Suzaku, and Swift), we do not enforce strict simultaneity within this window. The NuSTAR light curve in Figure 2 (binned at 25 ks, live-time and vignetting corrected) shows that during the Suzaku window, marked by the shaded region, the flux changed only gradually by ~10%. We therefore do not expect sudden flux changes to have occurred between occultation/SAA periods. We extracted the full time range for all instruments except NuSTAR, where we truncated the observation by applying the Suzaku start and stop times. Because the majority of the INTEGRAL observation lies outside the near-simultaneous window, where the flux is significantly different, we excluded INTEGRAL from this broadband fit.

We fitted an absorbed power-law model using Wilms abundances (Wilms et al. 2000) and Verner cross sections (Verner et al. 1996), freezing the hydrogen column to the Galactic value of $1.79 \times 10^{20}$ cm$^{-2}$ (Dickey & Lockman 1990). We keep this value fixed in the remainder of the fits. To avoid complications with the soft excess, we limited our fitting range between 1 and 78 keV, though we note that during these observations the soft excess appears to have been very modest. We allowed all normalizations to float; the relative cross-normalization terms and the relative slope errors between instruments are discussed in detail in K. K. Madsen (2015, in preparation) by the IACHEC consortium.

Figure 3 shows the results of the best fit power-law: $\Gamma = 1.647 \pm 0.003$. Formally the fit is good, with $\chi^2_{\text{red}} = 1.018$ (7815 dof). However, visual inspection of the high-energy residuals above 20 keV in NuSTAR shows obvious systematic deviations from a power law. This good formal fit is due to the very large number of bins between 2 and 8 keV. Replacing the power law with a cutoff power law, $\text{cutoffpwrlw}$ (XSPEC), reduces the residuals, and we find $\Gamma = 1.624 \pm 0.006$, $E_{\text{cutoff}} = 291^{+55}_{-50}$ keV ($\chi^2_{\text{red}} = 1.013$). To examine any further features, we fitted a cutoff power law to NuSTAR alone, ignoring the energy ranges 5–8 keV and 10–50 keV, and plot the NuSTAR ratio of the model to data alone in Figure 4. The shape of these residuals is reminiscent of reflection of the primary continuum off dense material, with an iron line and a Compton hump.

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Residuals of the data to model of the fits between NuSTAR and INTEGRAL for a pexrav (top) and pegpwrlw (bottom) fit. Only NuSTAR FPMA is shown for ease of viewing.

3.2. Fitting NuSTAR and INTEGRAL

We investigated joint NuSTAR and INTEGRAL observations over a broader time range, dispensing with the low-energy instruments. We applied the INTEGRAL GTIs from the July observation to the NuSTAR data set. Fitting the data with a pexrav alone results in INTEGRAL diverging above 80 keV. A power law provides a reasonable broadband fit, but with strong NuSTAR residuals between 20 and 78 keV as shown in Figure 5. Following the idea that the X-ray spectrum is in reality a superposition of an AGN component (i.e., coronal emission plus disk reflection) with a jet component, we fitted the data with the pexrav model for the AGN component and a flux-pegged power law, pegpwrlw (XSPEC), for the jet.

The two models prove to be degenerate. To partially break this degeneracy, we investigated three INTEGRAL data sets taken in 2011 December, 2012 June, and 2012 July. We fitted the spectra with pegpwrlw from 30 to 250 keV. The individual fits are summarized in Table 2, with the spectral energy distributions (SEDs) shown in Figure 6. The spectral slope changes with time, but the flux difference decreases with increasing energy and converges at ~120 keV. We also draw attention to the dip in the June and July observations between 60 and 80 keV; the error bars are large, but this could be interpreted as the AGN turnover seen by NuSTAR.
Notes.

\(^a\) Energy range 40–100 keV in units of \(10^{12}\) erg cm\(^{-2}\) s\(^{-1}\).

\(^b\) Energy range 80–150 keV in units of \(10^{12}\) erg cm\(^{-2}\) s\(^{-1}\).

The SED of the jet component is a log-parabola \(N(E) = N_0 \log P (E/E_0)^{\alpha - \beta \log(E/E_0)}\), which has the advantage of adding only one parameter to the overall model. Parameter \(\alpha\) corresponds to the local photon index at the photon energy of \(E_0\), and here we choose \(E_0 = 100\) keV. We added this model to the best-fit INTEGRAL model from Table 3, which we kept fixed, and we fitted it jointly to the NuSTAR, INTEGRAL, and Fermi/LAT data (Figure 8). We obtained a good fit with \(\chi^2 = 1.171\), with the results \(\alpha = 1.428 \pm 0.011\) and \(\beta = 0.091 \pm 0.002\). This places the SED peak of the jet component at \(E_{\text{peak}} = (2.28 \pm 0.21)\) MeV, consistent with the results of Esposito et al. (2015).

We note that the INTEGRAL data points are systematically above the fit. But a perfect fit is not expected, since 3C 273 was found in a relatively low \(\gamma\)-ray state, which required integrating the \(\gamma\)-ray flux over a timescale of 50 days, far longer than the 6-day NuSTAR campaign. Short-timescale variability might therefore have put the jet in a higher flux state during the NuSTAR period than on average. The fit, however, shows that the inferred slope of the jet is indeed reasonable.

### 3.3. The Iron Line

Although the detection of an iron line in 3C 273 has in the past been firm, it remains faint and at times even absent. To maximize our signal, we use all of the 244 ks of the NuSTAR observation and just NuSTAR alone. We restricted the energy range between 3 and 10 keV. In this limited band an absorbed power law provides a good fit, but with a slight excess between 5 and 7 keV as can be observed in Figure 4. We added a Gaussian component with a fixed rest energy of 6.4 keV, and the quality of fit improves from \(\chi^2/\text{dof} = 419/345\) without to \(\chi^2/\text{dof} = 398/343\) with. A Monte Carlo simulation shows this feature to be significant in excess of 99.97%.

The width of the line is broad, \(\sigma_{\text{width}} = 0.65 \pm 0.3\) keV, and we measure an equivalent width of \(\text{EW} = 23 \pm 11\) eV, which is within the envelope of what has been previously reported (Yaqoob & Serlemitsos 2000; Kataoka et al. 2002; Grandi & Palumbo 2004; Page et al. 2004).

We note that the magnitude of the iron emission is close to the calibration limit (±1%), and as a separate test we fixed the
line energy to 6.4 keV and allowed the redshift to vary. We recovered the correct redshift within errors. We also replaced the line with a relativistic disk line model but did not find significant improvement to the fit.

3.4. Continuum Models

As discussed in the previous sections, the 3C 273 continuum data can be well fit with an absorbed cutoff power law with weak reflection (a pexrav model) up to the NuSTAR upper energy limit at 78 keV. However, inclusion of the INTEGRAL data reveals that the best fit between the two observatories requires an additional power law for the jet component. Unfortunately, INTEGRAL does not cover the entire NuSTAR observation, and the rise of the jet component is not visible in the NuSTAR data alone. We therefore proceed with our investigation of the continuum by using the full 244 ks NuSTAR observation, but we make the assumption that the jet component is constant for the entire time range and takes on the values given in Table 3. We realize this to be an approximation, and therefore to contrast, we also fit the data without the jet component to bracket parameters.

The results of the continuum fits are shown in Table 4 and the residual of data to model in Figure 9. We first fit the continuum with a pexrav model with and without a Gaussian iron line included and compare with a pexmon (Nandra et al. 2007), which self-consistently includes the iron line. The two models yield very similar parameters, with a photon index $\Gamma = 1.63 \pm 0.01$ and $E_{\text{cutoff}} \sim 260 \pm 35$ keV for the model without the jet and $\Gamma = 1.66 \pm 0.01$ and $E_{\text{cutoff}} \sim 50 \pm 2$ keV with the jet. The relative reflection for both cases is around $R \sim 0.1$.

To probe a more physical model for the coronal continuum, we fitted with a CompTT model (Titarchuk 1994) and included the reflection component using pexmon. The CompTT model calculates the emergent spectrum resulting from Comptonization in an electron plasma, taking into account relativistic effects. It has models for a disk and spherical plasma geometry, and we adopted the disk geometry. We assumed solar abundances, fixed the inclination angle at 35°, and used as soft photon temperature from the disk 10 eV based on the temperature–luminosity relation from Makishima et al. (2000) for a 10% accretion efficiency and $10^9 M_\odot$ black hole (Courvoisier 1998). We fixed the values of the reflection component to the values found for the pexmon fit, checking that the shape of the direct spectrum of the pexmon model matched that of the CompTT direct spectrum, and included only the reflection part of the model.

Table 4 shows the results from fitting the CompTT disk model to the NuSTAR data with and without a jet component (with jet parameters again fixed to match the INTEGRAL $E > 80$ keV data as provided in Table 3). The inclusion of the jet component significantly changes the best-fit parameters. Without the jet, we find values of the plasma temperature of

![Figure 7](image1.png)

**Figure 7.** INTEGRAL and NuSTAR fit. (a) Fit from Table 3 (black/red = FPMA/FPMB, green = ISGRI), (b) Ratio between jet and AGN component for three flux settings of the jet (given in units of $10^{-12}$ erg cm$^{-2}$ s$^{-1}$). (c) Contours of input spectrum of the pexrav. (d) Contours of the photon index of the AGN against the jet component.

![Figure 8](image2.png)

**Figure 8.** Broadband SED for 3C 273 where the AGN component (pexrav) has been kept fixed to the values from Table 3 and the jet fitted with a log-parabola that peaks at $E_{\text{peak}} = 2.28 \pm 0.21$ MeV.
temperature where the model is unfortunately not considered reliable owing to the large amount of scattering (Poutanen & Svensson 1996).

3.5. Variability

The NuSTAR light curve (Figure 2) shows that 3C 273 went through flux changes of 10%–30% during the observation. We calculated the noise-subtracted variance fraction of the mean, $F_{\text{var}}$, using the formalism of Edelson et al. (2002), and Figure 10 shows that $F_{\text{var}}$ statistically does not vary as a function of energy during our observation.

We investigated the variability of the spectrum by decomposing the observation into five time bins, GTI A–GTI E (start and stop times are listed in Table 5), as illustrated in Figure 11. The motivation behind this partitioning is to cover the sections where the flux is roughly constant. We chose to continue with the pexrav model and fitted both with and without the jet component, assuming that it was frozen to the values given in Table 3. The two models and the parameters are summarized in Table 5.

We find that by enforcing a jet component, we obtain lower $E_{\text{cutoff}}$ energy since more bending is required when the spectrum is diluted by the jet, and the jet component systematically softens the index. The jet component also systematically increases the relative reflection for the same reason.

Because of the coupling between the photon index, cutoff energy, and relative reflection, it is hard to determine the correlation between flux and slope for the AGN component. To approach this question in a model-independent manner, we decomposed the spectrum into four energy bands, 3–10 keV, 10–20 keV, 20–40 keV, and 40–78 keV, and fit each section with a phenomenological power law, calculating the flux for the intervals. The results are shown in Figure 12. The highest flux point marked with an open symbol is GTI A. We fitted a linear function to all data points (solid line) and to GTI B–GTI E, excluding the GTI A point (dashed line).

The linear Pearson correlation (Press et al. 1992) favors in all but the lowest energy range the GTI B–GTI E (P1 series) over GTI A–GTI B–GTI E (P2 series): $P1/P2 = -0.73/ -0.86$ for 3–10 keV, $P1/P2 = -0.79/ -0.64$ for 10–20 keV, $P1/P2 = -0.64/0.14$ for 20–40 keV, and $P1/P2 = -0.70/ -0.62$ for 40–78 keV. In the lowest-energy band we find that the photon index is inversely correlated with the flux, becoming harder for higher fluxes. This correlation persists through all bands when excluding GTI A, but it starts to break down in the 10–20 keV band when GTI A is included. Since the source during GTI A was very stable compared to later times, it might suggest that the source was in a different state altogether.

4. DISCUSSION

From the radio to $\gamma$-ray bands the SED of 3C 273 is dominated by nonthermal radiation from the jet; however, evidence for emission from the inner accretion flow has been argued based on the presence of a weak iron line and soft excess in the X-ray and the big blue bump in the optical/UV, which are all common features of AGNs (Kataoka et al. 2002; Grandi & Pulumbo 2004; Pietrini & Torricelli-Ciamponi 2008). Evidence for hard X-ray emission from the corona in the form of a cutoff power law and reflection features above 10 keV has been elusive, perhaps owing to the available

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**Table 4**

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Without Jet</th>
<th>With Jet</th>
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<tr>
<td>$\chi^2_{\text{red}}$</td>
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<td>1.039</td>
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<tr>
<td>$\Gamma$</td>
<td>1.63 ± 0.01</td>
<td>1.66 ± 0.01</td>
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<tr>
<td>$E_{\text{cutoff}}$</td>
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<td>52 ± 2 keV</td>
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<td>Relative reflection</td>
<td>0.07 ± 0.03</td>
<td>0.2 ± 0.05</td>
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<table>
<thead>
<tr>
<th>Parameter</th>
<th>pexrav$^a$</th>
<th>pexrav$^a$+zgauss</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\chi^2_{\text{red}}$</td>
<td>1.033</td>
<td>1.035</td>
</tr>
<tr>
<td>$\Gamma$</td>
<td>1.63 ± 0.01</td>
<td>1.66 ± 0.01</td>
</tr>
<tr>
<td>$E_{\text{cutoff}}$</td>
<td>260 ± 33 keV</td>
<td>52 ± 2 keV</td>
</tr>
<tr>
<td>Relative reflection</td>
<td>0.07 ± 0.03</td>
<td>0.2 ± 0.03</td>
</tr>
<tr>
<td>Line energy (fixed)</td>
<td>6.4 keV</td>
<td>6.4 keV</td>
</tr>
<tr>
<td>Line width $\sigma$</td>
<td>0.62 ± 0.02</td>
<td>0.87 ± 0.07 keV</td>
</tr>
<tr>
<td>EW</td>
<td>15 ± 11 eV</td>
<td>15 ± 11 eV</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Parameter</th>
<th>pexmon$^a$</th>
<th>CompTT$^a$-pexmon</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\chi^2_{\text{red}}$</td>
<td>1.040</td>
<td>1.145</td>
</tr>
<tr>
<td>Geometry</td>
<td>disk</td>
<td>disk</td>
</tr>
<tr>
<td>$T_e$ (fixed)</td>
<td>10 eV</td>
<td>10 eV</td>
</tr>
<tr>
<td>Plasma temp, kT</td>
<td>133 ± 0.03 keV</td>
<td>12.0 ± 0.3 keV</td>
</tr>
<tr>
<td>$\tau_p$ (fixed)</td>
<td>0.33 ± 0.08</td>
<td>2.77 ± 0.06</td>
</tr>
<tr>
<td>$\Gamma$ (fixed)</td>
<td>1.63</td>
<td>1.64</td>
</tr>
<tr>
<td>$E_{\text{cutoff}}$ (fixed)</td>
<td>257 keV</td>
<td>53 keV</td>
</tr>
<tr>
<td>Relative reflection (fixed)</td>
<td>0.04</td>
<td>0.08</td>
</tr>
</tbody>
</table>

**Note.**

$^a$ Solar abundances.

$kT_e = 247^{+69}_{-64}$ keV, while with the jet included the temperature decreases to $kT_e = 12.0 ± 0.3 ± 1$ keV and increases the optical depth from $\tau_p = 0.15^{+0.04}_{-0.04}$ without the jet to $\tau_p = 2.77 ± 0.06$ with the jet. As shown (Figure 9), the fit without the jet is excellent and shows no indication of a jet component. The fit with the jet displays residuals at high energies, indicating that our assumption of a constant jet is not correct, or that the constraints from INTEGRAL are not sufficient to properly constrain this component. Allowing the jet to fit freely, we can find better solutions, but these results push the CompTT model up against its valid phase space and are thus questionable. We can, however, place an upper limit of $kT_e \sim 133^{+42}_{-20}$ keV on the electron temperature, since inclusion of the jet in any form decreases the value from the no-jet scenario. Conversely, the optical depth has a lower limit of $\tau_p = 0.33 ± 0.1$. Given that the uncertainties in the CompTT physical parameters are completely driven by the uncertainties in the level of the jet flux, we do not investigate alternative coronal geometries further.

We attempted to fit with a CompPS model (Poutanen & Svensson 1996), which is a coronal plasma model that includes reflection off a cold disk in the manner of the pexrav, but the returned fits were within a regime of high optical depth and low
instrument sensitivities or because the source was observed in states with high jet flux, which could have obscured the AGN signatures.

The NuSTAR data show that the 3–78 keV spectrum cannot be fit simply by an absorbed power law, but the spectrum requires a rollover above $\sim$20 keV. Modeling the NuSTAR data alone, we find evidence for a weak reflected continuum consistent with reflection from a cold accretion disk or distant material. We confirm the presence of a weak iron line, and assuming that it is neutral, we measure its parameters to be $\sigma = 0.65 \pm 0.3$ keV with an EW = 23 $\pm$ 11 eV, consistent with previous measurements (Yaqoob & Serlemitsos 2000; Kataoka et al. 2002; Grandi & Palumbo 2004; Page et al. 2004). If we include a reflection component, we can account for the weak line and excess above 10 keV (the compton hump) with a reflection fraction $R = 0.07 \pm 0.03$ and $E_{\text{cut-off}} = 262 \pm 34$ keV.

Figure 9. Ratio plots of data to model for the fits in Table 4. Not shown is the pexrav+zgauss since the improvement in these residuals is in the iron-line region covered in Figure 4.

Figure 10. Noise-subtracted variance fraction of the mean for the four energy bands: 3–10 keV, 10–20 keV, 20–40 keV, and 40–78 keV. Statistically there is no evidence of a difference in variability as a function of energy.

Considering NuSTAR together with INTEGRAL data for a temporally overlapping window, we show that the broadband (3–150 keV) spectrum cannot be explained with a single spectral model. Extrapolating the NuSTAR best-fit AGN model above 80 keV reveals a clear excess in INTEGRAL that we associate with emission from the jet. If we attempt to fit the combined INTEGRAL and NuSTAR spectrum with a model including the AGN components and an additional power law to describe the jet emission, we find strong degeneracies and no unique fit. We therefore fix the flux of the jet component from INTEGRAL 80–150 keV data and investigate the range of fits to the broadband spectrum corresponding to the range of jet flux uncertainty. To further limit degeneracies, we fix the relative reflection to $R = 0.15$. We find that the photon index for the coronal component ranges from 1.6 to 1.7, with a cutoff energy of 30–70 keV. The jet photon index is poorly constrained, ranging from 0.5 to 1.5. Taking into account the high-energy $\gamma$-ray spectrum from the Fermi/LAT, and assuming a log-parabolic shape of the jet component, we find a local photon index at 100 keV of about 1.4. We caution that the complete lack of data in the 1–100 MeV range means that the broadband shape of the jet component is uncertain. Nevertheless, the log-parabola model predicts a modest luminosity of the SED peak at 2 MeV. If the photon index at 100 keV was in the range 0.5–1, the implied peak luminosity of the jet component would need to be significantly higher, and that would be more demanding from the point of view of jet energetics. We note that a similar log-parabola model added to an AGN component was found to be successful by Esposito et al. (2015), who analyzed the Swift, INTEGRAL, and Fermi data for 3C 273 in different spectral states. In any case, we find that the jet component begins to dominate over the AGN component in the 30–50 keV range.
The cutoff energy range of 50–70 keV found with the jet included is at the low end compared to other sources where a cutoff energy has been measured by NuSTAR (Ballantyne et al. 2014; Brenneman et al. 2014; Marinucci et al. 2014; Baloković et al. 2015; Matt et al. 2015); however, we note that considerable uncertainty arises owing to the poor constraints on the jet emission. If we exclude the jet component, an upper limit for the energy cutoff is $E_{\text{cutoff}} \sim 260 \pm 40$ keV (see Table 4).

We explore a physical model for the AGN component and fit with the coronal Comptonizing coronal electron plasma (Titarchuk 1994). If we fix the jet component parameters at the best-fit values given in Table 3, we find an unusually low plasma temperature of $kT = 12 \pm 1$ keV, and high optical depth $\tau_p = 2.77 \pm 0.06$. Without the jet component included we find $kT = 133^{+82}_{-42}$ keV and optical depth $\tau_p = 0.33 \pm 0.1$. Since the jet contribution is poorly constrained, we can only conclude that the actual values lie between these extremes.

The reflection fractions we derive for all models we consider are low ($R = 0.02–0.2$; see Tables 4 and 5). These indicate very weak reflection, as is often found in broad-line radio galaxies (BLRGs; Wozniak et al. 1998; Eracleous et al. 2000; Zdziarski & Grandi 2001; Sambruna et al. 2009; Ballantyne et al. 2014). Possible explanations for the weak reflection include high inner disk ionizations (Ballantyne et al. 2002), a change in inner disk geometry (Eracleous et al. 2000; Lohfink et al. 2013), obscuration of the central accretion flow by the jet (Sambruna et al. 2009), black holes with retrograde spin (Evans et al. 2010), and dilution of the X-ray spectrum by jet emission (Grandi et al. 2002). Since we still only find a small reflection fraction even when taking into account the jet, it is likely not due to dilution of the jet. It has instead been suggested that a
geometrically thick accretion disk, which obscures part of itself from the external source (Paltani et al. 1998), could be responsible for the low reflection fraction. Considering the possibility of a high optical thickness of the corona, it could also be that the corona itself is smearing out the reflection signature.

3C 273 is well known for its variability. During the NuSTAR observation, it went through several flux changes. Our study found no statistical evidence for a change in variability as a function of energy in the NuSTAR band. The disentanglement and correlation of the jet and AGN component with respect to flux remain unclear. In the low-energy band (3–10 keV) we find an inverse correlation between flux and slope, hardening the spectrum with increasing flux. This correlation seems to persist through all bands when excluding the highest flux bin, but it appears to break down above the 10 keV band when including the high flux bin. This inverse correlation in the low-energy band has been previously observed in the 2–10 keV flux band by Kataoka et al. (2002) during their 1999–2000 campaign with RXTE. However, during their 1996–1997 campaign this correlation was no longer present. They attributed this change to the emergence of the AGN component during 1999–2000 when the source was at a higher flux than in 1996–1997. We postulate that it may rather be that a higher flux level of the jet is responsible for the hardening of the spectrum. An inverse correlation is not common for AGNs, where spectra typically get steeper with increasing flux (but we note that the opposite is generally true for blazars, both those associated with quasars and the lineless variety; respective examples are 3C 279, Hayashida et al. 2015; and Mkn 421, Baloković et al. 2013). This could indicate that the variability is instead driven by the jet flux, hardening the spectrum as the jet flux contributes more strongly. The highest flux bin corresponding to GTI A does break the correlation, but considering the change in the light curve itself, there is an indication that the source transitioned from one state to another at the end of GTI A.

If we assume the two-component scenario, it appears that the AGN was bright at the beginning, clearly outshining the jet. Subsequently, lower flux softened the index at all energies, but it is not possible to tell what combination of jet and AGN flux levels went into generating the variability since the two components are not expected to be correlated (Soldi et al. 2008).

A long-standing issue in 3C 273 is that the Comptonization of UV photons predicts a correlation to exist between UV and X-ray flux, which was claimed to have been observed in the past by Walter & Courvoisier (1992) but later found absent by Chernyakova et al. (2007). In the jet/AGN component scenario the superposition of two uncorrelated model components would affect the reliability of a correlation between UV and X-rays, and the disappearance of it may simply therefore be a matter of jet domination.

The spectral structure inferred by the consideration of the NuSTAR and INTEGRAL data could possibly also be explained by invoking a two-component inverse Compton model arising in a relativistic jet. In such a scenario, the MeV-to-GeV range emission is likely to be dominated by scattering of the isotropic radiation external to the jet (as in External Radiation Compton, or ERC; see, e.g., Sikora et al. 1994), while the slight “hump” detected by NuSTAR and modeled by us as Compton reflection from the accretion disk would be due to synchrotron self-Compton (SSC) emission, with the target photons being the synchrotron photons in the jet itself (Maraschi et al. 1992). We consider such a scenario less likely: the NuSTAR data (as well as previous X-ray observations) detect a broad Fe K line, which is a signature of reprocessing from the accretion disk, which is likely to be accompanied by the reflection component as in our modeling. Finally, as argued by Esposito et al. (2015), the timing analysis performed by those authors argues for the presence of the Seyfert component in addition to the jet component. However, any firm conclusions regarding the origin of the jet radiation—namely, the relative dominance of the SSC versus ERC component—are not possible with the data presented here.

5. CONCLUSION

In the coordinated observing campaign in 2012 we observed 3C 273 in a state where the jet flux was relatively weak. This made it possible to observe spectral signatures from coronal emission and constrain the weak Compton reflection in the hard X-ray band. We can separate the spectral components from the inner accretion flow from the jet flux in the 3–150 keV band quantitatively only by making an explicit assumption about jet flux level. To truly separate the two components and constrain the physical parameters of the corona, and/or address the possibility of SSC emission from the jet masking as reflection, further simultaneous observations with NuSTAR, INTEGRAL, and Fermi of the source in a state with low levels of flux from the jet are required.

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Facilities: Fermi, Chandra (CXO), INTEGRAL, NuSTAR, Suzaku, Swift, and XMM.

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