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Validity of PEC Approximation for On-Body Propagation

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Abstract—Many articles on on-body propagation assumes that the human body can be approximated by a perfect electric conductor (PEC) instead of the actual constitutive parameters of the human body, which is that of a lossy dielectric. This assumption is investigated in this article through comparison of the scattering of a plane wave at oblique incidence by a PEC and a lossy dielectric cylinder. The investigation shows that the validity of the assumption depends on the polarization of the plane wave, the angle of incidence, and the region of interest.

I. INTRODUCTION

The area of on-body propagation has been of high interest in recent years. Wireless communications for on-body devices serves as one of the enabling factors in the emerging commercial area of wearables.

Numerous investigations have been made on the on-body propagation and many models have been proposed. In many of these the body is approximated by a perfect electric conductor (PEC) of some shape, [1]–[5]. Often these models justify directly or indirectly this assumption based on [6]. But since [6] is concerned only with the disturbance in the channel between two off-body antennas this seems to be a poor validation.

The purpose of this work is to investigate if modeling the body as a PEC is a valid assumption. Furthermore, if the assumption might be valid under certain other assumptions.

The on-body propagation will be investigated by the use of the exact solution to the fields around an infinitely long cylinder. The fields will be calculated for a PEC cylinder and for a human body average cylinder. The material parameters of the human body are documented in [7]. The average material parameters in [8] is used ($\epsilon_r = 39.2$ and $\sigma = 1.8 \text{ S/m}$). The cylinder has a radius of $a = 80 \text{ mm}$, which is similar to the radius of the leg or the radius of the head. The fields are calculated for both TM and TE polarization and at three different angles of incidence.

II. THEORY

The infinitely long lossy dielectric cylinder is located in free space. A plane wave is incident and propagating in the direction given by $\hat{d} = (0, \sin \alpha, \cos \alpha)$ as seen in Fig. 1. The propagation constant in the cylinder and in free space

Fig. 1. The geometry of the infinitely long dielectric cylinder model.

are k_1 and k , respectively. Two other propagation constants are defined as $k' = k \sin \alpha$ and $k'_1 = \sqrt{k_1^2 - k^2 \cos^2 \alpha}$. The electric and magnetic fields outside the cylinder incident by a transverse magnetic (TM) plane wave are given by the following formulas obtained from [9], [10]:

$$E_\rho = -jk \cos \alpha \frac{\partial u}{\partial \rho} - \frac{j\mu\omega}{\rho} \frac{\partial v}{\partial \phi} \quad (1)$$

$$E_\phi = -\frac{jk \cos \alpha}{\rho} \frac{\partial u}{\partial \phi} + j\mu\omega \frac{\partial v}{\partial \rho} \quad (2)$$

$$E_z = k'^2 u \quad (3)$$

$$H_\rho = \frac{jk^2}{\mu\omega\rho} \frac{\partial u}{\partial \phi} - jk \cos \alpha \frac{\partial v}{\partial \rho} \quad (4)$$

$$H_\phi = -\frac{jk^2}{\mu\omega} \frac{\partial u}{\partial \rho} - \frac{jk \cos \alpha}{\rho} \frac{\partial v}{\partial \phi} \quad (5)$$

$$H_z = k'^2 v \quad (6)$$

where a common factor of $e^{-jkz \cos \alpha}$ has been suppressed as well as a time dependency of $e^{j\omega t}$. To obtain E_ρ and E_ϕ inside the cylinder the permeability μ must be replaced by

Fig. 2. Electric field strength (V/m) plot for TM plane wave incident on a human-body average (a, b, c) and PEC (d, e, f) cylinder at different angles.

μ_1 . To obtain E_z and H_z inside the cylinder the propagation constant k' must be replaced by k'_1 . To obtain H_ρ and H_ϕ inside the cylinder μ and k must be replaced by μ_1 and k_1 in the terms involving u and remain the same for the terms involving v . For $\rho > a$ the scalar potentials u and v are given by [9]:

$$u = \frac{-E_0 \sin \alpha}{k'^2} \sum_{n=-\infty}^{\infty} \left(J_n(k'\rho) - a_n H_n^{(2)}(k'\rho) \right) e^{-jn\phi} \quad (7)$$

$$v = \frac{-E_0 \sin \alpha}{k'^2} \sum_{n=-\infty}^{\infty} c_n H_n^{(2)}(k'\rho) e^{-jn\phi} \quad (8)$$

and for $\rho < a$ by:

$$u = \frac{-E_0 \sin \alpha}{k'^2} \sum_{n=-\infty}^{\infty} b_n J_n(k'_1 \rho) e^{-jn\phi} \quad (9)$$

$$v = \frac{-E_0 \sin \alpha}{k'^2} \sum_{n=-\infty}^{\infty} d_n J_n(k'_1 \rho) e^{-jn\phi} \quad (10)$$

where a_n, b_n, c_n , and d_n are constants chosen to ensure the continuity of the tangential components of the E-fields and

H-fields across the boundary. The constants are given by:

$$a_n = \frac{J_n(k'a)}{\Delta H_n^{(2)}} \left\{ \left[\frac{\mu_1}{\mu k'_1 a} \frac{J'_n(k'_1 a)}{J_n(k'_1 a)} - \frac{1}{k' a} \frac{H_n^{(2)'}(k' a)}{H_n^{(2)}(k' a)} \right] \cdot \left[\frac{k_1^2}{k^2} \frac{\mu}{\mu_1 k'_1 a} \frac{J'_n(k'_1 a)}{J_n(k'_1 a)} - \frac{1}{k' a} \frac{J'_n(k' a)}{J_n(k' a)} \right] - \left[\frac{1}{(k' a)^2} - \frac{1}{(k'_1 a)^2} \right]^2 n^2 \cos^2 \alpha \right\} \quad (11)$$

$$b_n = \frac{k'^2}{k_1'^2} \left[J_n(k' a) - a_n \frac{H_n^{(2)}(k' a)}{J_n(k'_1 a)} \right]$$

$$c_n = -\frac{2kn \cos \alpha}{\Delta \mu \omega \pi (k' a)^2} \frac{1}{H_n^{(2)}(k' a)} \left[\frac{1}{(k' a)^2} - \frac{1}{(k'_1 a)^2} \right]$$

$$d_n = \frac{k'^2}{k_1'^2} c_n \frac{H_n^{(2)}(k' a)}{J_n(k'_1 a)}$$

where

$$\Delta = \left[\frac{\mu_1}{\mu k'_1 a} \frac{J'_n(k'_1 a)}{J_n(k'_1 a)} - \frac{1}{k' a} \frac{H_n^{(2)'}(k' a)}{H_n^{(2)}(k' a)} \right] \cdot \left[\frac{k_1^2}{k^2} \frac{\mu}{\mu_1 k'_1 a} \frac{J'_n(k'_1 a)}{J_n(k'_1 a)} - \frac{1}{k' a} \frac{H_n^{(2)'}(k' a)}{H_n^{(2)}(k' a)} \right] - \left[\frac{1}{(k' a)^2} - \frac{1}{(k'_1 a)^2} \right]^2 n^2 \cos^2 \alpha.$$

Fig. 3. Electric field strength (V/m) plot for TE plane wave incident on a human-body average (a, b, c) and PEC (d, e, f) cylinder at different angles.

The fields for a transverse electric (TE) plane wave is obtain by interchanging E and H and replacing μ by $-\frac{k^2}{\omega^2\mu}$ and multiplying the fields with $\frac{k}{\mu\omega}$. Note that there is a sign error on c_n in [9].

For the numerical calculation the infinite series need to be truncated at a suitable number N , where the series have converged. Experience shows that truncating at $N = k\rho + 20$ yields very accurate results. To check the validity of the numerical code the fields at the boundary between the cylinder and free space was calculated and the tangential component was found to be continuous. The incident field strength was chosen to be 1 V/m.

III. NUMERICAL RESULTS

The results are seen for the TM polarization and the TE polarization in Fig. 2 and Fig. 3, respectively.

For the TM polarization, for normal incidence, $\alpha = 90^\circ$, it is seen that the fields are almost the same for the PEC and the human-body cylinder. The only slight difference is the front reflection. Gradually, as the angle of incidence α decreases the difference becomes larger as most clearly seen for the plots of the electric field at an incidence angle of $\alpha = 30^\circ$. It is seen that the fields on the back side are higher for the human-body cylinder than for the PEC cylinder. The field strength on the back side increases for lower angles of incidence. In front of

the cylinder the reflection is seen to be much more powerful for the PEC cylinder than for the human-body cylinder.

For the TE polarization the fields are more different at normal incidence, i.e., $\alpha = 90^\circ$. The reflection is seen to be stronger in front of the PEC cylinder than for the human-body cylinder. On the back side the fields are significantly higher for the PEC cylinder. As the angle of incidence decreases the differences between the PEC and human-body cylinder increases. The fields on the back side are significantly higher for the PEC cylinder.

In general, the TM polarization seem to give quite similar results for the PEC and human-body cylinder whereas the difference is slightly larger for the TE polarization. It is noted that the situation in [6] is actually TM polarized and therefore there is a good correspondence with the results found in this reference. Furthermore, the differences in both cases become more pronounce when the angle of incidence decreases.

IV. CONCLUSION

The accuracy of the use of PEC as an approximation of the human body has been investigated. It was found that for normal incidence and TM polarization the human body is very well modeled by PEC. At normal incidence and TE polarization the human body is quite well modeled by PEC. As the angle of incidence become more oblique, the

differences between PEC and actual human-body parameters become more profound. For very oblique angles of incidence and TE polarization the validity of approximating the human-body by PEC is questionable.

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