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3-D Imaging using Row–Column-Addressed 2-D Arrays with a Diverging Lens

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Abstract—It has been shown that row–column-addressed (RCA) 2-D arrays can be an inexpensive alternative to fully addressed 2-D arrays. Generally imaging with an RCA 2-D array is limited to its forward-looking volume region. Constructing a double-curved RCA 2-D array or applying a diverging lens over the flat RCA 2-D array, can extend the imaging field-of-view (FOV) to a curvilinear volume without increasing the aperture size, which is necessary for applications such as abdominal and cardiac imaging. Extended FOV and low channel count of double-curved RCA 2-D arrays make it possible to have 3-D imaging with equipment in the price range of conventional 2-D imaging. This study proposes a delay-and-sum (DAS) beamformation scheme specific to double-curved RCA 2-D arrays and validates its focusing ability based on simulations. A synthetic aperture imaging (SAI) sequence with single element transmissions at a time, is designed for imaging down to 14 cm at a volume rate of 88 Hz. The curvilinear imaging performance of a $\lambda/2$ -pitch 3 MHz 62+62 RCA 2-D array is investigated as a function of depth, using a diverging lens with f-number of -1. The results of this study demonstrate that the proposed beamforming approach is accurate for achieving correct time-of-flight calculations, and hence avoids geometrical distortions.

I. INTRODUCTION

An $N \times N$ element 2-D array can be operated utilizing only $2N$ connections, when a row–column or cross-electrode addressing scheme is used [1]–[3]. This is contrary to the N^2 connections needed, when conventionally addressing the elements. In general, a row–column-addressed (RCA) array is a 2-D matrix array, which is addressed via its row- and column indices. Effectively, this makes two 1-D arrays arranged orthogonal to each other. As an example, a $256 + 256$ RCA array will have 512 elements. A 2-D matrix array of equivalent size would have 65,536 elements, over a factor of 7 more than the current state-of-the-art X6-1 PureWave xMATRIX probe from Phillips (Eindhoven, Netherlands) that has 9212 elements [4]. This exhibits the potential of having very large RCA 2-D arrays with low channel count and real-time capabilities.

It has been demonstrated in several studies [1]–[3] that row–column technology is a realistic alternative to the state-of-the-art matrix probes, especially as a low-cost alternative. However, one major issue with the RCA arrays is that they can only emit acoustic energy directly below the array and in a cross-shape to the sides. For applications such as abdominal imaging, it is

relevant to have a probe with a large aperture capable of phased array imaging. True volumetric phased array imaging is possible with RCA arrays, provided that the array is double curved to spread the energy during transmit [1]. However, manufacturing curved transducer elements is challenging for both capacitive micromachined ultrasonic transducer (CMUT) and piezoelectric transducer (PZT) technologies. Another approach to spread the acoustic energy is by using a double curved diverging acoustic lens on top of the RCA array [5]. Using a lens makes it easier to fabricate curved transducers, as it is not needed to manufacture curved elements, and also making a lens is a well-tested technology. An in-depth study of the possibilities of this approach is therefore the main goal of this study.

It is investigated whether curvilinear volumetric imaging is possible with an RCA array equipped with a diverging lens. A dedicated beamformer is developed and the performance is evaluated using Field II [6], [7] simulations. The quality assessments of the B-mode images, *i.e.*, spatial resolution and contrast resolution, are carried out on the simulated data using SAI technique. The remainder of the paper is organized as follows: The current limitations with flat RCA arrays and different approaches to disperse the acoustic energy are described in Section II. Section III presents a DAS beamformer for a double curved RCA array. In section IV, a detailed overview of the simulation setup is presented. Section V explains the simulation results and the final section concludes the paper.

II. CURVED RCA 2-D ARRAYS

To spread the acoustic energy of a line-element, the element has to be curved to generate a diverging wave. However, manufacturing double curved transducers is very challenging. Alternatively, the defocusing of the waves can be made by using a fixed electronic delay profile along each flat line-element to generate a diverging wavefront. This can be seen as using micro-beamforming with a fixed first stage. Another simpler approach to spread the acoustic energy is by using a double curved diverging acoustic lens on top of the RCA array [5]. A concave diverging lens can be designed with a material having a lower speed of sound compared to the human tissue, so that it has higher thickness around the corners and the sides of the array, and less thickness close to center of the array.

Alternatively, a convex diverging lens can be made from a material with a higher speed of sound compared to the human tissue, which is preferred for a better contact surface. A flat diverging lens can also be made by using a combination of two different materials, one with higher and other one with lower speed of sound compared to the human tissue, which results in a flat surface for good contact.

Note that, for the same aperture size, lower lens $f_{\#}$ values for the lens corresponds to larger thicknesses of the lens, and therefore the attenuation becomes higher through the lens material. Thus, there is a trade-off between field-of-view (FOV) and attenuation. The $f_{\#}$ is defined as the ratio between the focal distance to the lens diameter. For example the delay profile can be in a range of $0 \mu\text{s}$ to $3.5 \mu\text{s}$ for a lens with $f_{\#} = -0.7$ and a speed of sound of 1400 m/s , which corresponds to a thickness range of 0 mm to 5 mm . A suitable material for a lens could be Sylgard 160 (PDMS) with a density of 1580 kg/m^3 and a speed of sound of 950 m/s and attenuation of $0.4f^{1.4} \text{ dB/mm}$, where f is the operating frequency in MHz. Therefore, for an operating frequency of 3 MHz the maximum attenuation is 6.14 dB at the largest thickness [8]. This might be compensated by doubling the amplitude of the excitation pulse.

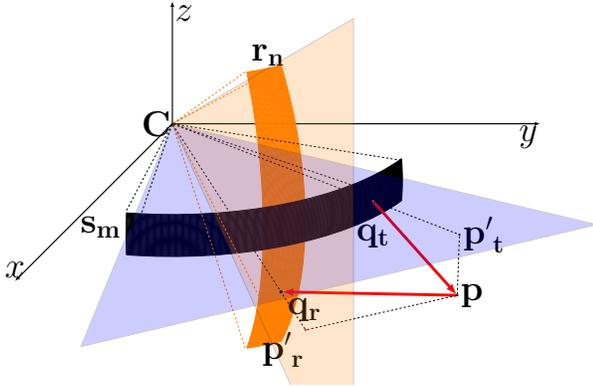


Fig. 1. The time-of-flight of a wavefront is given by the shortest distance from the source s_m to the point being focused p and back to the receiving element r_n , divided by the speed of sound. The points q_r and q_t are the closest contact points on the transmit and receive elements s_m and r_n to the point p .

III. DAS BEAMFORMING WITH CURVED RCA 2-D ARRAYS

The time-of-flight (ToF) of a wavefront is given by the shortest distance from the arc source s_m to the point being focused p , and back to the receiving element r_n , divided by the speed of sound. Using the notations from Fig. 1 this can be written as:

$$\text{ToF}_m(n, \mathbf{p}) = \frac{d(s_m, \mathbf{p}) + d(r_n, \mathbf{p})}{c}, \quad (1)$$

where c is the speed of sound in the medium, n is an index from 1 to the number of receive line elements N , and m is the emission index. The function $d(\cdot, \cdot)$ calculates the shortest distance between an arc and a point in space, which will be defined in the remainder of this section.

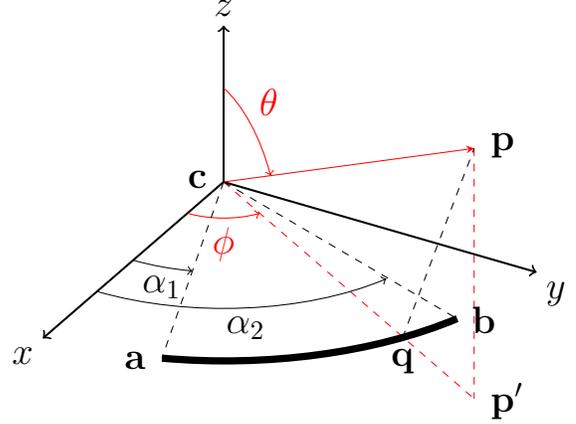


Fig. 2. Distance between a point p and an arc \widehat{ab} is calculated using (4).

The arc segment from point a to point b with center c is termed \widehat{ab} assuming the center at origin. This is illustrated in Fig. 2. The projection of point p onto the plane passing through the arc \widehat{ab} and its center c is termed p' and is determined by the usual equation for projection. To determine if the vector cp' is in between vector ca and vector cb , we define the normalized cross products \hat{l}_a and \hat{l}_b as

$$\hat{l}_a = \frac{\mathbf{cp}' \times \mathbf{ca}}{\|\mathbf{cp}'\| \|\mathbf{ca}\|}, \quad \hat{l}_b = \frac{\mathbf{cp}' \times \mathbf{cb}}{\|\mathbf{cp}'\| \|\mathbf{cb}\|}. \quad (2)$$

Depending on the location of the point p , vectors \hat{l}_a and \hat{l}_b can be either \hat{j} or $-\hat{j}$, where \hat{j} is the unit vector of the z -axis. \hat{l}_a and \hat{l}_b have different signs, when $\alpha_1 \leq \phi \leq \alpha_2$, and they have the same sign, when $\alpha_2 \leq \phi$ or $\phi \leq \alpha_1$. Here α_1 , α_2 , and ϕ are the angles between the x -axis and the vectors ca , cb , and cp , respectively, as shown in Fig. 2.

When \hat{l}_a and \hat{l}_b have different signs, *i.e.* $\hat{l}_a = \hat{j}$ and $\hat{l}_b = -\hat{j}$, or $\hat{l}_a = -\hat{j}$ and $\hat{l}_b = \hat{j}$, the standard formula for the distance between an arc and a point can be used:

$$d = \sqrt{\|\mathbf{pp}'\|^2 + (\|\mathbf{cp}'\| - R)^2}, \quad (3)$$

where R is the curvature of the arc and equal to $\|\mathbf{ca}\|$ or $\|\mathbf{cb}\|$.

When \hat{l}_a and \hat{l}_b have the same signs, *i.e.* $\hat{l}_a = \hat{j}$ and $\hat{l}_b = \hat{j}$, or $\hat{l}_a = -\hat{j}$ and $\hat{l}_b = -\hat{j}$, the shortest distance from the arc segment to the point is the distance from the closest end of the arc segment (a or b) to the point p . The following therefore determines the minimum distance between the point p and the arc segment \widehat{ab} :

$$d(\widehat{ab}, \mathbf{p}) = \begin{cases} \sqrt{\|\mathbf{pp}'\|^2 + (\|\mathbf{cp}'\| - R)^2} & \text{if } \begin{cases} \hat{l}_b = -\hat{j} \text{ and } \hat{l}_a = \hat{j} \\ \hat{l}_b = \hat{j} \text{ and } \hat{l}_a = -\hat{j} \end{cases} \\ \|\mathbf{ap}\| & \text{if } \hat{l}_b = -\hat{j} \text{ and } \hat{l}_a = -\hat{j} \\ \|\mathbf{bp}\| & \text{if } \hat{l}_b = \hat{j} \text{ and } \hat{l}_a = \hat{j} \end{cases} \quad (4)$$

Using (4), the distances $d(\widehat{\mathbf{s}}_m, \mathbf{p})$ and $d(\widehat{\mathbf{r}}_n, \mathbf{p})$ can now be determined. The focused signal at point \mathbf{p} is calculated by summing all receive signals at the time instances given by (1):

$$z_m(\mathbf{p}) = \sum_{n=1}^N a_{\text{elec}}(n, \mathbf{p}) y_{m,n}(\text{ToF}_m(n, \mathbf{p})), \quad (5)$$

where N is the number of receive elements, a_{elec} is the electronic receive apodization, and $y_{m,n}(t)$ is the measured signal from emission m on the receive element n at time t .

IV. SIMULATION AND MEASUREMENT SETUP

In this work, Field II [6], [7] is used for all simulations. The simulated receive signals are beamformed using a MATLAB (MathWorks Inc., Massachusetts, USA) implemented the proposed DAS beamformer for curved RCA arrays. The simulation parameters of a RCA 62+62 element 2-D array are shown in Table I. The receive array is rotated 90° with respect to the transmit array. Field II is set up to use lines to describe the apertures and each line-element is divided into square mathematical sub-elements with a side length of $\lambda/4$. To remove the otherwise apparent secondary echoes originating from the either ends of arc shaped elements, two roll-off apodization regions are placed at both ends of each element [9]. The length of each apodization region was equal to 15 times the pitch of the array. Each mathematical sub-element in both transmit and receive arrays is delayed according to the lens delay profile.

Table I
TRANSDUCER AND SIMULATION PARAMETERS.

Parameter name	Notation	Value	Unit
Number of elements	–	62+62	–
Center frequency	f_0	3.0	MHz
Speed of sound	c	1480	m/s
Wave length	λ	493.3	μm
Array pitch -x	d_x	$\lambda/2 = 246.6$	μm
Array pitch -y	d_y	$\lambda/2 = 246.6$	μm
Sampling frequency	f_s	120	MHz
Emission pulse	–	2-cycles, Hann-weighted	–
Lens focal ratio	$f_\#$	-1	–

V. RESULTS AND DISCUSSION

In Fig. 3 the pulse-echo energy as a function of lateral position for different lens $f_\#$ is illustrated. For the flat array, the pulse-echo energy drops by moving away from the forward looking region of the array. This drop of the energy is due to the diffraction of the sound waves. At around 20 mm lateral position the pulse-echo energy drops by 40 dB, when no lens is used. However, by using a diverging acoustic lens on top of an RCA 2-D array a larger FOV can be attained. The FOV can be adjusted by using different $f_\#$ values for the lens. The main advantage of using a diverging lens or designing a curved 2-D array, is to disperse the transmit and receive fields, so that they overlap in a larger area. By using a diverging lens with $f_\# = -1$, the overlapped transmit and receive region increases to $\pm 26.5^\circ$ in both directions and the energy is maintained

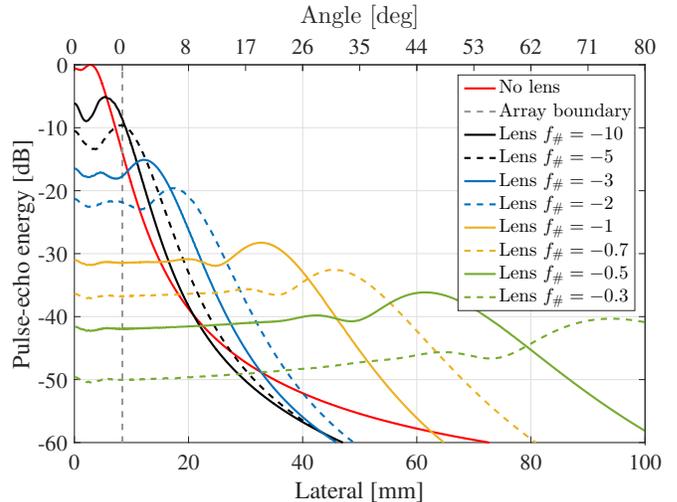


Fig. 3. Comparison of the pulse-echo energy for different lens $f_\#$ values. The points are located on a line at 80 mm away from the surface of the transducer.

within the FOV (yellow solid line in Fig. 3). By using all the elements in the transmit and placing the transmit focus in front of the array, this drop of the energy might be compensated partially.

Fig. 4 is illustrating three cross-planes (azimuth, elevation and C-plane) of a phantom with point targets simulated and beamformed with the proposed DAS beamforming method, with and without a diverging lens. The point targets are located along elevation dimension from -30 mm to 30 mm in an axial range of 5 mm to 95 mm. It can be seen from the figures that by using a diverging lens the FOV is extended compared to the flat RCA array.

To study the PSF characteristics as a function of lateral angle and radial distance, a point scatterer is imaged by sweeping it from 0° to 40° in lateral plane with steps of 10° at radial distances from 10 mm to 60 mm from the center of the array. At each radial distance and lateral angle the full-width at half-maximum (FWHM) and cystic resolution (CR) values are calculated over a volume of $10 \text{ mm} \times 10 \text{ mm} \times 10 \text{ mm}$ surrounding the point target. Fig. 5 is illustrating the measured FWHM and CR values as a function of depth and lateral angle. Moving away from the center of the elements towards the edges, the transmit wavefronts contact each other at a sharper point compared with the contact point at the center. This can be observed in Fig. 5 that, by moving away from the center towards the higher angular position in lateral plane, the elevation FWHM values become smaller while the CR values become larger. On the other hand, the lateral FWHM values stay constant, this is due to the intersection of the wavefronts in receive direction which is at the center of the receive elements and therefore the elevation FWHM values stay constant for all lateral angles.

VI. CONCLUSION

In this paper, the imaging performance of a curved 62+62 RCA 2-D array with a diverging lens, is quantitatively demon-

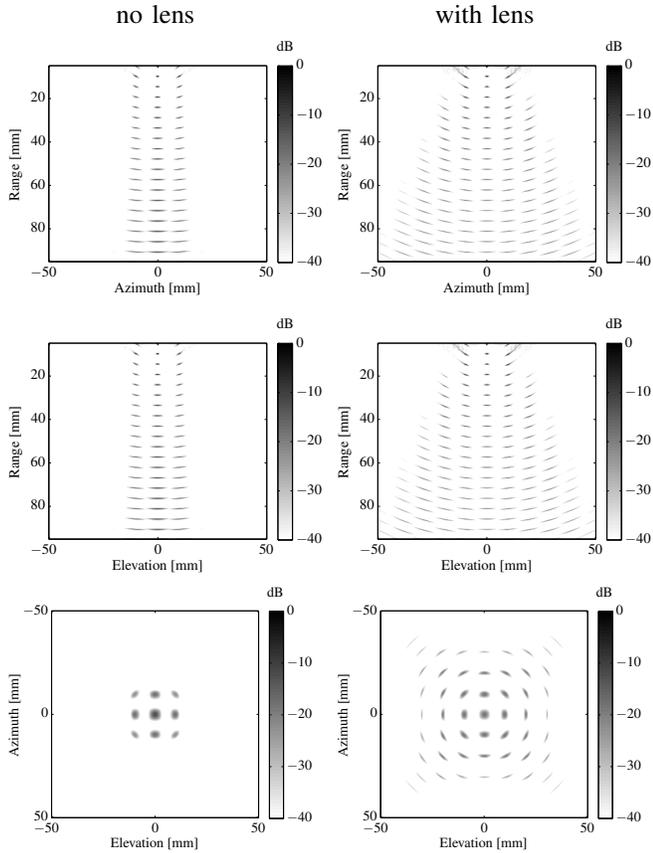


Fig. 4. Three cross-planes (azimuth, elevation, and C-plane) of a phantom with point scatterers imaged with and without a diverging lens ($f_{\#} = -1$) normalized to their maximum values, are shown at a dynamic range of 40 dB. The C-planes are at depth of 47 mm.

strated. A SAI sequence with single element transmissions at a time, was designed for imaging down to 14 cm at a volume rate of 88 Hz. The capabilities of a curved RCA 2-D array to effectively focus in both transmit and receive are investigated. A suitable DAS beamformer was introduced and implemented. Simulated results confirm that using a diverging lens can increase the imaging FOV and also that it is possible to perform dynamic transmit-receive focusing throughout the curvilinear FOV. Thereby, the inherent imaging limitation with flat RCA 2-D arrays, *i.e.*, its forward looking rectilinear FOV, is overcome by using a diverging lens. Overall, having a low channel count and a large FOV, offers the potential to fabricate arrays with large aperture sizes, which is important for abdominal scans. Thus by using a curved RCA 2-D array, 3-D imaging with equipment in the price range of conventional 2-D imaging is possible.

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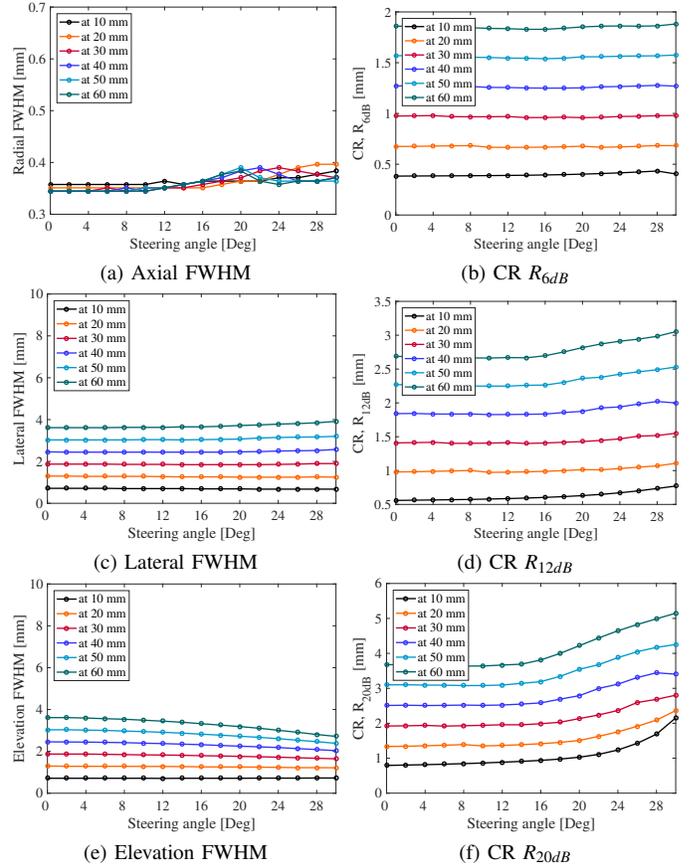


Fig. 5. CR and FWHM values calculated for point targets located at radial distances from 10 mm to 60 mm as a function of different azimuth steering angle away from the central forward-looking axis of the array.

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