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TESS Discovery of a Transiting Super-Earth in the pi Mensae System

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orbital period of 6.27 days. Radial-velocity data from the High-Accuracy Radial-velocity Planet Searcher and Anglo-Australian Telescope/University College London Echelle Spectrograph archives also displays a 6.27 day periodicity, confirming the existence of the planet and leading to a mass determination of $4.82 \pm 0.85 M_{\oplus}$. The star's proximity and brightness will facilitate further investigations, such as atmospheric spectroscopy, asteroseismology, the Rossiter–McLaughlin effect, astrometry, and direct imaging.

Key words: planetary systems - planets and satellites: detection - stars: individual (HD 39091, TIC 261136679)

1. Introduction

The mission of the *Transiting Exoplanet Survey Satellite* (*TESS*; Ricker et al. 2015) is to search for transiting planets as small as Earth around the nearest and brightest stars. Four 10 cm optical telescopes are used to repeatedly image wide fields and monitor the brightness of suitable stars. The data are then searched for periodic dips that could be caused by transiting planets. The spacecraft was launched on 2018 April 18 and began the sky survey on July 25. Here, we report on the discovery of a small transiting planet around a bright star π Men.

 π Men (also known as HD 39091) is a naked-eye G0V star at a distance of 18.27 \pm 0.02 pc (Gaia Collaboration et al. 2018) with a mass of 1.1 M_{\odot} and a radius of 1.1 R_{\odot} . Doppler monitoring by Jones et al. (2002) and Wittenmyer et al. (2012) revealed a planet (π Men b) with a mass about 10 times that of Jupiter, an orbital period of 5.7 yr, and an orbital eccentricity of 0.6. With a visual apparent magnitude of 5.67, the star is a prime target for the *TESS* survey. It is one of several hundred thousand pre-selected stars for which data will be available with 2 minute time sampling, as opposed to the 30 minute sampling of the full image data set.

This Letter is organized as follows. Section 2 presents the *TESS* photometric data that led to the detection of the new planet π Men c, as well as the archival radial velocity data that confirm the planet's existence. Section 3 describes our methods for determining the system parameters, including the mass and radius of the star and planet. Section 4 discusses some possible follow-up observations that will be facilitated by the star's brightness and proximity to Earth.

2. Observations and Data Reduction

2.1. TESS Photometry

The *TESS* survey divides the sky into 26 partially overlapping sectors, each of which is observed for approximately one month during the two-year primary mission. π Men is located near the southern ecliptic pole in a region where six sectors overlap, implying that it is scheduled to be observed for a total of 6 months. This Letter is based on data from Sector 1 (2018 July 25–August 22), during which π Men was observed with charge-coupled device (CCD) 2 of Camera 4.

The data were processed with two independently written codes: the MIT Quick Look Pipeline (partially based on fitsh; Pál 2009), which analyzes the full images that are obtained with 30 minute time sampling; and the Science Processing Operations Center pipeline, a descendant of the *Kepler* mission pipeline based at the NASA Ames Research Center (Jenkins et al. 2010), which analyzes the 2 minute data that are obtained for pre-selected target stars. For π Men, both pipelines detected a signal with a period of 6.27 days, an amplitude of about 300 ppm, a duration of 3 hr, and a flat-bottomed shape consistent with the light curve of a planetary transit.

Previous surveys taught us that transit-like signals sometimes turn out to be eclipsing binaries that are either grazing, or blended with a bright star, causing the amplitude of the signal to be deceptively small and resemble that of a planet (e.g., Cameron 2012). In this case, the signal survived all of the usual tests for such "false positives." There is no discernible secondary eclipse, no detectable alternation in the depth of the transits, and no detectable motion of the stellar image on the detector during the fading events.⁴⁵

After identifying the transits, we tried improving on the light curve by experimenting with different choices for the photometric aperture, including circles as well as irregular pixel boundaries that enclose the blooming stellar image. Best results were obtained for the aperture shown in Figure 1. Also shown are images of the field from optical sky surveys conducted 30–40 yr ago, long enough for the star to have moved about an arcminute relative to the background stars. This allows us a clear view along the line of sight to the current position of π Men, which is reassuringly blank: another indication that the transit signal is genuine and not an unresolved eclipsing binary. The other stars within the photometric aperture are too faint to cause the 300 ppm fading events.

The top panel of Figure 2 shows the result of simple aperture photometry. Most of the observed variation is instrumental. There may also be a contribution from stellar variability, which is expected to occur on the 18 day timescale of the rotation period (Zurlo et al. 2018). To remove these variations and permit a sensitive search for transits, we fitted a basis spline with knots spaced by 0.3 days, after excluding both 3σ outliers and the data obtained during and immediately surrounding transits. We then divided the light curve by the best-fitting spline.

The middle panel of Figure 2 shows the result. The scatter is 142 ppm per 2 minute sample, and 30 ppm when averaged into 6 hr bins, comparable to the highest-quality *Kepler* light curves. The gap in the middle of the time series occurred when observations were halted for data downlink. The other gap occurred during a period when the spacecraft pointing jitter was higher than normal. We also excluded the data from the 30–60 minute intervals surrounding "momentum dumps," when thrusters are fired to reorient the spacecraft and allow the reaction wheels to spin down. The times of the momentum dumps are marked in Figure 2. There were 10 such events during Sector 1 observations, occurring every 2 1/2 days.

2.2. Radial-velocity Data

 π Men has been monitored for 20 yr as part of the Anglo-Australian Planet Search, which uses the 3.9 m Anglo-Australian Telescope (AAT) and the University College London Echelle Spectrograph (UCLES; Diego et al. 1990). The long-period giant planet π Men b was discovered in this survey (Jones et al. 2002; Butler et al. 2006). A total of 77

 $^{^{45}}$ The last test in the list, the centroid test, was complicated by the fact that the star is bright enough to cause blooming in the *TESS* CCD images. The associated systematic effects were removed using the method of Günther et al. (2017).

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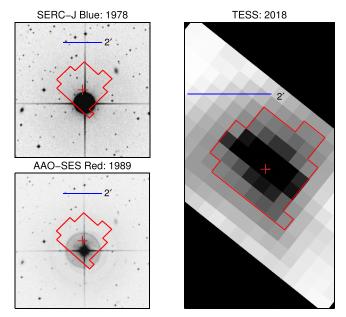


Figure 1. Images of the field surrounding π Men. Top-left panel: from the Science and Engineering Research Council J survey, obtained with a bluesensitive photographic emulsion in 1978. The red cross is the current position of π Men. Red lines mark the boundary of the *TESS* photometric aperture. Bottom-left panel: from the AAO Second Epoch Survey, obtained with a redsensitive photographic emulsion in 1989. Right panel: summed *TESS* image. North is up and East is to the left in all of the images.

radial velocities are available, obtained between 1998 and 2015, with a mean internal uncertainty of 2.13 m s⁻¹.

The star was also monitored with the High-Accuracy Radialvelocity Planet Searcher HARPS (Mayor et al. 2003) on the ESO 3.6 m telescope at La Silla Observatory in Chile. A hardware upgrade in 2015 June led to an offset in the velocity scale (Lo Curto et al. 2015). For this reason, our model allows for different constants to be added to the pre-upgrade and postupgrade data. A total of 145 radial velocities are available, obtained between 2003 December and 2016 March with irregular sampling. The mean internal uncertainty of the 128 pre-upgrade velocities is 0.78 m s⁻¹, while that of the 17 postupgrade velocities is 0.38 m s⁻¹.

The top panel of Figure 3 shows the radial-velocity data. It is easy to see the 400 m s⁻¹ variations from the giant planet. To search for evidence of the new planet, we subtracted the bestfitting single-planet model from the data and computed the Lomb–Scargle periodogram of the more precise HARPS data, shown in the middle panel. The highest peak is far above the 0.1% false alarm threshold and is located at the transit period of 6.27 days. The next highest peaks, bracketing a period of 1 day, are aliases of this signal. The phase of the 6.27 day signal is also consistent with the measured transit times.

We consider this to be a decisive confirmation of the existence of π Men c. Still, as another precaution against false positives, we checked the HARPS spectra for any indication of a second star, or spectral-line distortions associated with the 6.27 day signal. We re-analyzed the HARPS cross-correlation functions with the BLENDFITTER routine (Günther et al. 2018) and found no sign of any correlated bisector variations.

3. Determination of System Parameters

We performed a joint analysis of the two-planet system using the *TESS* transit light curve and the 222 radial velocities from

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Table 1 System Parameters for π Men

| Stellar Parameters | Value | Source |
|--|------------------------|--------------------------|
| Catalog Information | | |
| R.A. (h:m:s) | 05:37:09.89 | Gaia Data Release 2 (DR2 |
| Decl. (d:m:s) | -80:28:08.8 | Gaia DR2 |
| Epoch | 2015.5 | Gaia DR2 |
| Parallax (mas) | 54.705 ± 0.067 | Gaia DR2 |
| $\mu_{R.A.}$ (mas | 311.19 ± 0.13 | Gaia DR2 |
| yr^{-1}) | | |
| $\mu_{\text{Decl.}}$ (mas | 1048.85 ± 0.14 | Gaia DR2 |
| yr^{-1}) | | |
| Gaia DR2 ID | 4623036865373793408 | |
| HD ID | HD 39091 | |
| TIC ID | 261136679 | Stassun et al. (2018) |
| TOI ID | 144.01 | |
| Spectroscopic Proper | rties | |
| $T_{\rm eff}$ (K) | 6037 ± 45 | Ghezzi et al. (2010) |
| $\log g \ (\text{cgs})$ | 4.42 ± 0.03 | Ghezzi et al. (2010) |
| [Fe/H] (dex) | 0.08 ± 0.03 | Ghezzi et al. (2010) |
| $v \sin i (\mathrm{km} \mathrm{s}^{-1})$ | 3.14 ± 0.50 | Valenti & Fischer (2005) |
| Photometric Properti | es | |
| B (mag) | 6.25 | |
| V (mag) | 5.67 | |
| TESS (mag) | 5.1 | TIC V7 |
| Gaia (mag) | 5.491 | Gaia DR2 |
| $Gaia_r$ (mag) | 5.064 | Gaia DR2 |
| $Gaia_b (mag)$ | 5.838 | Gaia DR2 |
| J (mag) | 4.87 ± 0.27 | Two Micron All-Sky |
| | | Survey (2MASS) |
| H (mag) | 4.42 ± 0.23 | 2MASS |
| K_s (mag) | 4.241 ± 0.027 | 2MASS |
| Derived Properties | | |
| $M_{\star}(M_{\odot})$ | 1.094 ± 0.039 | This work |
| $R_{\star} \ (R_{\odot})$ | $1.10\pm0.023^{\rm a}$ | This work |
| L_{\star} (L_{\odot}) | 1.444 ± 0.02 | This work |
| Age (Gyr) | $2.98^{+1.4}_{-1.3}$ | This work |
| Distance (pc) | 18.27 ± 0.02 | Gaia DR2 |
| $\rho_{\star} (\text{g cm}^{-3})$ | 1.148 ± 0.065 | This work |

Note.

^a This is consistent with the radius of star derived from spectra energy analysis by Stassun et al. (2018).

the AAT and HARPS surveys. The orbit of planet c was assumed to be circular in the fit.⁴⁶ As noted previously, we assigned a different additive constant to each of the three radial-velocity data sets. We also required the systemic velocity to be zero and allowed for three independent values of the "jitter," a term that is added in quadrature to the internally estimated measurement uncertainty to account for systematic effects.

We assumed the star to follow a quadratic limb-darkening law and used the formulas of Mandel & Agol (2002) as implemented by Kreidberg (2015). We fixed the limbdarkening coefficients at $u_1 = 0.28$ and $u_2 = 0.27$, based on the tabulation of Claret (2017). The photometric model was computed with 0.4 minutes sampling and then averaged to 2 minutes before comparing with the data.

⁴⁶ We also tried allowing planet c to have an eccentric orbit, which resulted in an upper limit of $e_c < 0.3$ (1 σ). All of the other orbital parameters remained consistent with the results of the $e_c \equiv 0$ model, although naturally, some parameters were subject to slightly larger uncertainties.

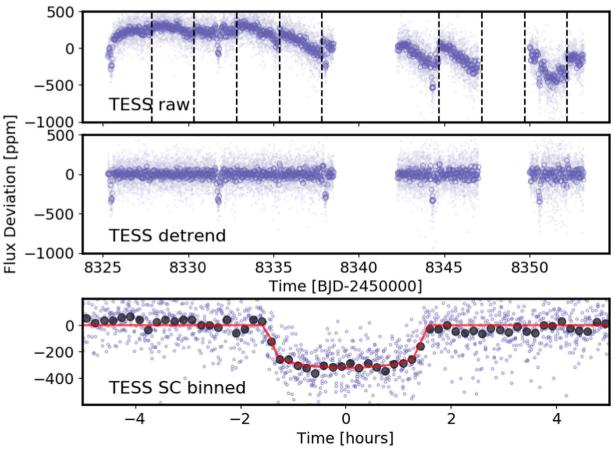


Figure 2. Raw (top panel) and corrected (middle panel) *TESS* light curves. The lighter points are based on the short cadence (SC) data with 2 minute sampling. The darker points are 30 minute averages. The dashed lines indicate the times of momentum dumps. The interruptions are from the data downlink and the pointing anomaly. The bottom panel shows the phase-folded light curve, along with the best-fitting model. The black dots represent 5 minute averages.

We also fitted for the mass and radius of the star, which were constrained by measurements of the spectroscopic parameters (Ghezzi et al. 2010) as well as the stellar mean density ρ_{\star} implicit in the combination of P, a/R_{\star} , and i (Seager & Mallén-Ornelas 2003; Winn 2010). For a given choice of mass, age, and metallicity, we relied on the Dartmouth stellar-evolutionary models (Dotter et al. 2008) to determine the corresponding radius R_{\star} , effective temperature T_{eff} , and *Gaia* absolute magnitude. The likelihood function enforced agreement with the measurements of T_{eff} , ρ_{\star} , log g, and parallax (based on the absolute and apparent *Gaia* magnitudes).

To determine the credible intervals for all of the parameters, we used the Markov Chain Monte Carlo (MCMC) method as implemented in emcee by Foreman-Mackey et al. (2013). Detrending was performed simultaneously with the transit fitting: at each step in the Markov Chain, the transit model (batman, Kreidberg 2015) was subtracted from the data and the residual light curve was detrended using a basis spline with knots spaced by 0.5 days. To avoid trying to model the discontinuities in the data related to momentum dumps, we only fitted the segment of the light curve in between momentum dumps. The results are given in Tables 1 and 2, and the best-fitting model is plotted in Figures 2 and 3. As a consistency check, we also fitted each of the five transits independently. Figure 4 shows the results, which are all consistent to within the estimated uncertainties.

 Table 2

 Parameters for the HD 39091 Planetary System

| Additional RV Parameters | RV offset | Instrument Jitter |
|---|-----------------------------------|--------------------------------------|
| AAT (m s ⁻¹) | 32.07 ± 0.86 | 6.7 ± 0.60 |
| HARPS pre-fix (m s ⁻¹) | 108.51 ± 0.40 | 2.33 ± 0.18 |
| HARPS post-fix (m s ⁻¹) | 130.60 ± 0.70 | 1.74 ± 0.33 |
| Planet Parameters | Planet b | Planet c |
| P (days) | 2093.07 ± 1.73 | 6.2679 ± 0.00046 |
| T_p (BJD) | 2445852.0 ± 3.0 | |
| T_c (BJD) | 2446087.0 ± 8.4 | $2458325.50400^{+0.0012}_{-0.00074}$ |
| $K ({\rm m}{\rm s}^{-1})$ | 192.6 ± 1.4 | $1.58^{+0.26}_{-0.28}$ |
| $\sqrt{e}\cos\omega$ | 0.6957 ± 0.0044 | |
| $\sqrt{e} \sin \omega$ | -0.392 ± 0.006 | |
| e | 0.637 ± 0.002 | 0 |
| ω | 330.61 ± 0.3 | |
| T_{14} (hr) | | 2.953 ± 0.047 |
| a/R_{\star} | | 13.38 ± 0.26 |
| R_p/R_{\star} | | $0.01703^{+0.00025}_{-0.00023}$ |
| $b \equiv a \cos i / R_{\star}$ | | $0.59^{+0.018}_{-0.020}$ |
| i_c (deg) | | $87.456_{-0.076}^{+0.085}$ |
| Derived Parameters | | |
| M_p | | $4.82^{+0.84}_{-0.86}~M_{\oplus}$ |
| $M_p \sin i$ | $10.02\pm0.15~\textit{M}_{\rm J}$ | |
| $R_p (R_{\rm c})$ | | 2.042 ± 0.050 |
| $\rho_p \text{ (g cm}^{-3}\text{)}$ | | $2.97_{-0.55}^{+0.57}$ |
| $\log g_p$ (cgs) | | $3.041\substack{+0.07\\-0.86}$ |
| a (AU) | 3.10 ± 0.02 | 0.06839 ± 0.00050 |
| $T_{\rm eq}$ (K) ^h | | $1169.8^{+2.8}_{-4.3}$ |
| $\langle F_j \rangle \ (10^9 \ {\rm erg \ s^{-1} \ cm^{-2}})$ | | $0.42_{-0.09}^{+0.04}$ |

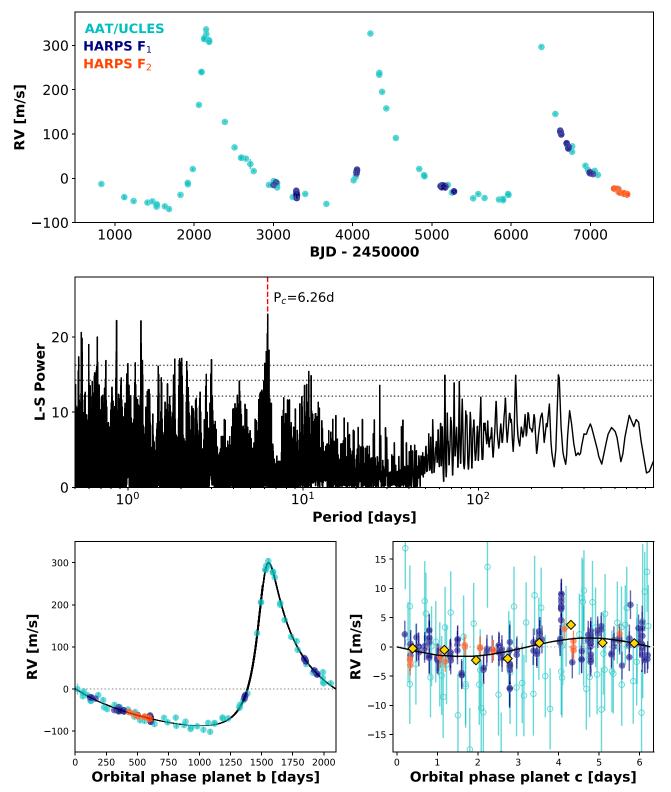


Figure 3. Top panel: relative radial velocity of π Men as measured with UCLES and HARPS (both pre- and post-upgrade). The zero points of each of the three data sets have been adjusted to coincide. Middle panel: Lomb–Scargle periodogram of the HARPS data, after subtracting the single-planet model that best fits the entire data set. The dotted lines are the power levels corresponding to false alarm probabilities of 10%, 1%, and 0.1%. Bottom-left panel: radial velocity as a function of the orbital phase of planet b, after subtracting the best-fitting model for the variation due to planet c. Bottom-right panel: similar, but for planet c. The orange point is binned in phase space.

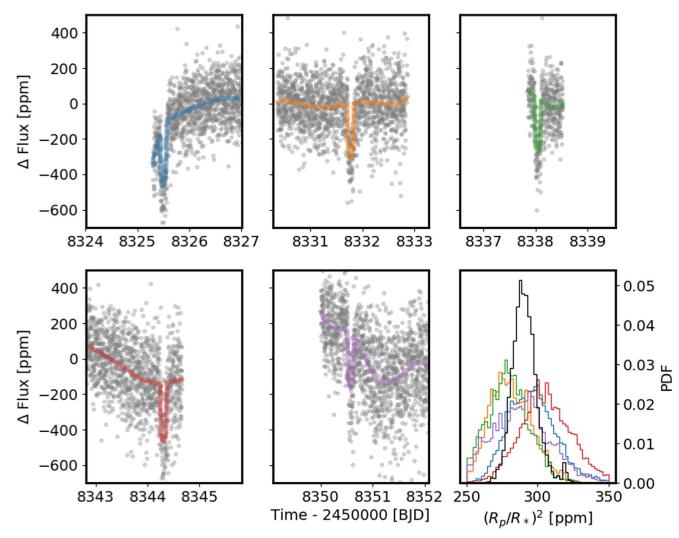


Figure 4. Plotted are the *TESS* data that surround each of the five observed transits and were obtained in between momentum dumps. Each panel shows 3 days of data and spans the same range of flux deviations. In the bottom-right panel, the colored histograms are the five posterior distributions for $(R_p/R_*)^2$, obtained from independent fits to the five transit data sets. The black histogram is the posterior based on the fit to all of the data.

4. Discussion

Among the known stars with transiting planets, π Men is the second brightest in the visual band, as illustrated in the top panel of Figure 5. *TESS* has begun to fulfill its promise to enlarge the collection of small, transiting planets orbiting bright stars. Such stars enable precise measurements of that planet's mass and radius. The bottom panel of Figure 5 shows the measured masses and radii of the known planets smaller than Neptune, overlaid with theoretical mass/radius relationships for different compositions. π Men c falls above the "pure rock" curve on the diagram, and near curves for planets composed of either pure water or rocky interiors surrounded by a lightweight 1% H/He envelope. π Men c must not have a purely rocky composition, but instead may have a rocky core surrounded by layers of volatiles, such as hydrogen/helium (see Owen & Wu 2017), or water/methane (Vanderburg et al. 2017).

With a near-infrared magnitude of K = 4.24, π Men is also one of the brightest stars available for planetary atmospheric characterization with the *James Webb Space Telescope (JWST)*. π Men c is one of the top 10 most favorable systems in the ranking scheme of Kempton et al. (2018), although this ranking scheme does not take into account the practical difficulties in achieving photon-limited observations of such a bright star. Transit spectroscopy would be difficult if the planet has an Earth-like atmospheric scale height of order 10 km, in which case the atmospheric signals would be on the order of only 1 ppm. On the other hand, given the intense stellar irradiation, there may be larger signals form an escaping atmosphere (see, e.g., Ehrenreich et al. 2015; Spake et al. 2018). Spectroscopy of occultations (secondary eclipses) is also promising. The occultation depth is predicted to be 60 ppm in the Rayleigh–Jeans limit, assuming the entire surface radiates as a blackbody at the equilibrium temperature of 1200 K.

Another interesting possibility is to measure the stellar obliquity by observing the Rossiter–McLaughlin (RM) effect. Stars with close-in giant planets show a surprising diversity of orientations (Winn & Fabrycky 2015; Triaud 2017). However, we know relatively little about the obliquities of stars with smaller planets, because the relevant signals are smaller and harder to detect. In the case of π Men c, the amplitude of the RM effect is on the order of 1 m s⁻¹, the product of the transit depth (300 ppm) and the sky-projected rotation velocity (3.1 km s⁻¹; Valenti & Fischer 2005).

The π Men system consists of a giant planet on a longperiod, highly eccentric orbit, along with a planet with an orbit

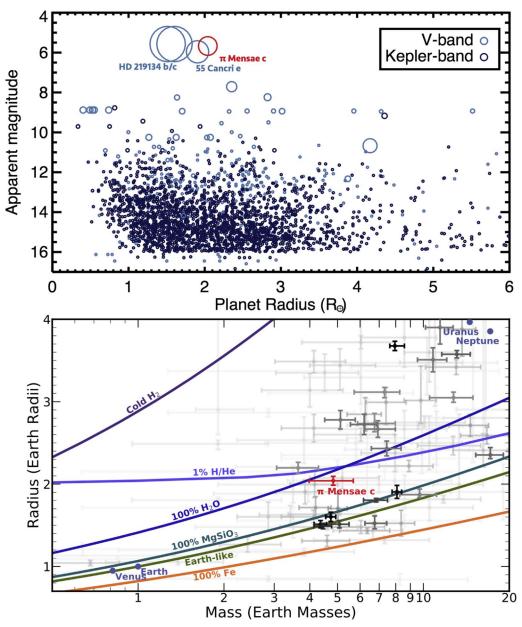


Figure 5. π Men c in the context of other known exoplanets. Top panel: apparent magnitude and planet radius for all of the known transiting planets. The V magnitude is plotted when available, and otherwise the *Kepler* magnitude is plotted. The symbol size is proportional to the angular diameter of the star. Bottom panel: mass-radius diagram for small exoplanets. Darker points represent more precise measurements. Based on data from the NASA Exoplanet Archive, accessed on 2018 September 13 (https://exoplanetarchive.ipac.caltech.edu/cgi-bin/TblView/nph-tblView?app=ExoTbls&config=planets). Model curves are: H_2 (Seager et al. 2007); 100% H_2O , 100% $M_gS_iO_3$, 100% F_e , Earth-like (Zeng et al. 2016); and 1% H/ H_e (Lopez et al. 2012).

and mass that are both smaller by two orders of magnitude. Recent follow-up studies of *Kepler* systems have suggested that they maybe intrinsically common (Bryan et al. 2018; Zhu & Wu 2018). Thus, we might find many similar cases with *TESS*, providing clues about the formation of close-orbiting planets, whether by disk migration, Lidov–Kozai oscillations, or other mechanisms.

Astrometric observations with the *Gaia* spacecraft might ultimately reveal the full 3D geometry of the system. Ranalli et al. (2018) predicted that the astrometric signal of π Men b will be detectable with a signal-to-noise ratio higher than 10 by the end of the mission. Indeed, the fit to the existing *Gaia* data exhibits an excess scatter of 295 μ'' (37 σ), perhaps a hint of planet-induced motion. Direct imaging might also be fruitful some day, although Zurlo et al. (2018) have already ruled out any companions with orbital separation 10-20 AU and an infrared contrast exceeding 10^{-6} , corresponding roughly to 30 Jupiter masses.

While some of these observations may be far off, we will not have to wait long for another opportunity to learn more about π Men. As mentioned earlier, *TESS* is scheduled to collect five additional months of data. This will allow us to refine our knowledge of planet c, search for additional transiting planets, and try to detect asteroseismic oscillations. The π Men system has already been generous to the exoplanet community, and with a little luck, the gifts will keep arriving.

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