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Estimation of High Velocities in Synthetic Aperture Imaging: II: Experimental Investigation

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Abstract—The paper describes the performance of a new pulse sequence design and estimation approach for increasing the maximum detectable velocity in synthetic aperture (SA) velocity imaging. Measurements are conducted for conventional imaging for comparing the velocity range detectable by a directional Transverse Oscillation (TO) autocorrelation estimator to a new cross-correlation estimator. For conventional focused emissions a 192-elements, 3 MHz convex array transducer is used together with the SARUS experimental scanner on a flow rig at beam-to-flow angles of 60°, 70° and 90°. Here the new estimator always yields a higher precision, and the aliasing limit is increased by a factor 3. The new SA inter-spaced scheme was investigated using Field II simulations and SARUS measurements. A 3 MHz, 128-elements phased array was employed with a 5 virtual source emissions scheme for flow estimation and 15 emissions for B-mode imaging. The scheme was interleaved three times for a positive, negative, and positive transmission, so that non-linear pulse inversion also could be made. The experiments were conducted at three angles and for 4 different pulse repetition frequencies. A peak transverse velocity of 0.51 m/s could be estimated at $f_{prf}=450$ Hz, translating to 5.6 m/s at $f_{prf}=5$ kHz showing the theoretical increase of a factor 10 predicted in the accompanying theory paper.

I. INTRODUCTION

Estimation of high velocities is of interest in many applications where diseases in the cardiovascular system are found. Velocities are increased proportionally to stenosis degree in vessels, which can result in peak velocities in excess of 2.3 m/s for stenosis of more than 70% in the carotid artery [1]. Such vessels are superficial, and the pulse repetition frequency f_{prf} can often be maintained high above e.g. 10 kHz. In cardiology incompetent valves can result in regurgitant jets with velocities up to 7 m/s, which currently have to be diagnosed using continuous wave Doppler systems. This limits the available isonation angles, and makes it difficult to select the measurement depth. A full view of the heart with

a flexible placement of the interrogation volume would be highly beneficial. This demands an $f_{prf} = 5$ kHz for imaging down to 15 cm, which in current synthetic aperture (SA) vector flow imaging (VFI) systems usually limits the maximum velocity to around 0.25 m/s for a 3 MHz probe. The aim of the paper is therefore to experimentally investigate methods with the capability of measuring velocities up to 3-7 m/s using pulse repetition frequencies of 5 to 10 kHz for SA VFI.

Methods for increasing the maximum detectable velocity were described in the accompanying paper [2]. The first method relies on a new pulsing scheme for SA VFI. In previous methods a series of emissions were produced to generate low resolution images (LRIs), and a high resolution image (HRI) was then assembled from the LRIs. Repeating the process generates a second HRI to be correlated with the first for estimating the spatial velocity distribution [3]. The correlation between measurements N emissions apart results in an effective pulse repetition frequency of $f_{eff,prf} = f_{prf}/N$, which lowers the maximum detectable velocity by a factor of N . Inter-leaving the sequences for the two or more HRIs reduces the time distance to $1/f_{prf}$ resulting in the largest maximum detectable velocity possible.

The second method employs a new directional transverse oscillation (DTO) cross-correlation vector flow estimator, which avoids the phase aliasing from the previous autocorrelation method [4], [5]. This can increase the maximum detectable velocity beyond the $\lambda_x f_{prf}/4$ aliasing limit, where λ_x is the lateral oscillation wavelength. The estimator uses data beamformed orthogonal to the ultrasound propagation direction and cross-correlates data along this directions. This makes the estimator independent of the actual value of λ_x , and no calibration has to be performed to yield close-to unbiased results.

The new estimator is investigated using flow rig measurements in Section IV, where the auto- and cross-

TABLE I
PARAMETERS USED FOR THE DUPLEX TO MEASUREMENTS.

Parameter	Value	Unit	Variable
Transducer elements	192		
Center frequency	3.0	MHz	f_0
Cycles in transmit pulse	4		
Transmit apodization	von Hann		
Element pitch	0.33	mm	P_i
Element height	13	mm	
Elevation focus	65	mm	
Convex radius	60	mm	
Transmit focus for vessel depths below 80 mm	105	mm	
Transmit focus for vessel depths beyond 80 mm	300	mm	
Pulse repetition frequency	2.5	kHz	f_{prf}
Active transmit elements	64		
Measurement sampling frequency	17.5	MHz	

correlation estimates are compared using conventional duplex ultrasound data. The new estimator is combined with the inter-leaved SA sequence in Section V for a phased array probe, and the performance found for both simulated and measured data.

II. MEASUREMENT SET-UP FOR CONVENTIONAL IMAGING

A convex array transducer with 192 elements and a center frequency of 3 MHz is used for the measurements in conjunction with the SARUS ultrasound scanner [6]. A duplex sequence is used, as described in [5]. A circulating flow rig with a 6 mm radius tube was scanned, and the volume flow recorded by a Danfoss Magnetic flow meter MAG1100 (Danfoss, Nordborg, Denmark). The volume flow was 106.4 l/h corresponding to a peak velocity in the vessel of 0.52 m/s assuming a fully developed parabolic profile. The beam-to-flow angles were 60°, 75°, and 90°, and the measurement depths varied from 24 to 156 mm. The data are identical to the ones used in [5], and all parameters are shown in Table I.

III. MEASUREMENT SET-UP FOR INTER-LEAVED SA DTO

A phased array transducer with 128 elements and a center frequency of 3 MHz is used for the measurements in conjunction with the SARUS ultrasound scanner [6]. The emission sequence contains both B-mode and flow emissions as shown in Fig. 1, and the polarity of the emissions are also inverted, so that pulse inversion imaging can be made. The pulse inversion is emitted

Fig. 1. Geometric depiction of the SA flow imaging sequence. The red circles indicate virtual source position for flow emissions and green circles are for B-mode emissions. The solid lines are the directions for the emissions, and the dashed lines indicate the acceptance angle used in the beamformation. The solid blue lines show the possible imaging region.

in a short sequence and the B-mode emissions are inter-spaced between the blocks of flow emissions as shown in Table II. The positions of the different emitted virtual sources and their directions are shown in Fig. 1.

In Table II $-v_2^{(1)}$ denotes a negative excitation waveform emission in direction 2 for the first time, and $+B_{15}^{(1)}$ denotes a positive B-mode emission in direction 15. Fifteen emissions are emitted for a full B-mode image, and five emissions are used for creating a high resolution image for velocity estimation. Five of the velocity emissions are also used in forming the B-mode image, and there are, thus, 10 blocks of emissions starting with two B-mode emissions (positive and negative for the same virtual source) followed by 3 times 5 flow emissions.

The scheme makes it possible to perform pulse inversion imaging. It also makes it possible to correlate HRIs created from $+v_{1...5}^{(1)}$ with $+v_{1...5}^{(3)}$, which only has a time difference of $2T_{prf}$. Alternatively sequences like $+v_{1...5}^{(1)} - v_{1...5}^{(2)}$ could be correlated with $+v_{1...5}^{(3)} - v_{1...5}^{(2)}$. Another possibility is to change the sign of the negative emission $-v_1^{(2)}$ to make the time difference equal to T_{prf} , which increases the maximum detectable velocity by a factor of 2.

The same circulating flow rig set-up as described in Section II were used for the measurements. All parameters are shown in Table III. Simulations using Field II [7], [8] were conducted with the exact same transducer parameters, inter-spacing, and acquisition parameters. The impulse response of the probe was used, but no

TABLE II

EMISSION SEQUENCE FOR THE SA B-MODE/FLOW SEQUENCE AND THE CORRESPONDING LOW AND HIGH RESOLUTION IMAGES. THE LEFT SIDE SHOWS THE EMISSION SEQUENCE AND THE RIGHT SIDE THE CORRESPONDING LRIS AND THEIR SUMMATION TO HRIS.

Emission 1	Emission 2	Emission 3	Low resolution images:		
$+B_1^{(1)}$	$-B_1^{(2)}$		$+LRI_{B_1}^{(1)}$	$-LRI_{B_1}^{(2)}$	
$+v_1^{(1)}$	$-v_1^{(2)}$	$+v_1^{(3)}$	$+LRI_1^{(1)}$	$-LRI_1^{(2)}$	$+LRI_1^{(3)}$
$+v_2^{(1)}$	$-v_2^{(2)}$	$+v_2^{(3)}$	$+LRI_2^{(1)}$	$-LRI_2^{(2)}$	$+LRI_2^{(3)}$
$+v_3^{(1)}$	$-v_3^{(2)}$	$+v_3^{(3)}$	$+LRI_3^{(1)}$	$-LRI_3^{(2)}$	$+LRI_3^{(3)}$
$+v_4^{(1)}$	$-v_4^{(2)}$	$+v_4^{(3)}$	$+LRI_4^{(1)}$	$-LRI_4^{(2)}$	$+LRI_4^{(3)}$
$+v_5^{(1)}$	$-v_5^{(2)}$	$+v_5^{(3)}$	$+LRI_5^{(1)}$	$-LRI_5^{(2)}$	$+LRI_5^{(3)}$
LRIs sums to:			$\sum_{i=1}^5 +LRI_i^{(1)}$ $= +HRI^{(1v)}$	$\sum_{i=1}^5 -LRI_i^{(2)}$ $= -HRI^{(2v)}$	$\sum_{i=1}^5 +LRI_i^{(3)}$ $= +HRI^{(3v)}$
$+B_2^{(1)}$	$-B_2^{(2)}$		$+LRI_{B_2}^{(1)}$	$-LRI_{B_2}^{(2)}$	
$+v_1^{(4)}$	$-v_1^{(5)}$	$+v_1^{(6)}$	$+LRI_1^{(4)}$	$-LRI_1^{(5)}$	$+LRI_1^{(6)}$
$+v_2^{(4)}$	$-v_2^{(5)}$	$+v_2^{(6)}$	$+LRI_2^{(4)}$	$-LRI_2^{(5)}$	$+LRI_2^{(6)}$
$+v_3^{(4)}$	$-v_3^{(5)}$	$+v_3^{(6)}$	$+LRI_3^{(4)}$	$-LRI_3^{(5)}$	$+LRI_3^{(6)}$
$+v_4^{(4)}$	$-v_4^{(5)}$	$+v_4^{(6)}$	$+LRI_4^{(4)}$	$-LRI_4^{(5)}$	$+LRI_4^{(6)}$
$+v_5^{(4)}$	$-v_5^{(5)}$	$+v_5^{(6)}$	$+LRI_5^{(4)}$	$-LRI_5^{(5)}$	$+LRI_5^{(6)}$
LRIs sums to:			$\sum_{i=1}^5 +LRI_i^{(4)}$ $= +HRI^{(4v)}$	$\sum_{i=1}^5 -LRI_i^{(5)}$ $= -HRI^{(5v)}$	$\sum_{i=1}^5 +LRI_i^{(6)}$ $= +HRI^{(6v)}$
$\vdots$	$\vdots$	$\vdots$			
$+B_{15}^{(1)}$	$-B_{15}^{(2)}$		$+LRI_{B_{15}}^{(1)}$	$-LRI_{B_{15}}^{(2)}$	
$+v_1^{(28)}$	$-v_1^{(29)}$	$-v_1^{(30)}$	$+LRI_1^{(28)}$	$-LRI_1^{(29)}$	$+LRI_1^{(30)}$
$\vdots$	$\vdots$	$\vdots$			
			B-mode emissions sums to:		
			$\sum_{i=1}^{10} +LRI_{B_j}^{(1)} + \sum_{i=1}^5 +LRI_i^{(1)}$ $= +HRI^{(1B)}$	$\sum_{i=1}^{10} -LRI_{B_j}^{(2)} + \sum_{i=1}^5 -LRI_i^{(2)}$ $= -HRI^{(2B)}$	

TABLE III

PARAMETERS USED FOR THE SA INTER-LEAVED MEASUREMENTS.

Parameter	Value	Unit	Variable
Transducer elements	128		
Center frequency	4.0	MHz	f_0
Transmit frequency	2.7	MHz	
Cycles in transmit pulse	2		
Transmit apodization	von Hann		
Element pitch	0.22	mm	P_t
Element height	15	mm	
Elevation focus	85	mm	
Active transmit elements	32		
Focal depth	5.74	mm	
Measurement sampling frequency	17.5	MHz	
Simulation sampling frequency	350	MHz	
Flow angles	60, 70, 90	°	
Pulse repetition frequencies	225, 450 1000, 4500	Hz	f_{prf}

non-linear effects were introduced in the entirely linear simulation.

A. Stationary echo canceling filters

Stationary echo canceling was performed by subtracting the mean signal. The mean signal was calculated across all HRIs used for velocity estimation, and it was subsequently subtracted from the individual HRIs. The same echo canceling approach was also employed for the measurements described in Section II. Other more advanced schemes based on decomposition could be used [9], [10], [11], [12], [13], [14]. This is, however, not necessary here, as all data employed are either simulated or measured in a flow rig, whether tissue motion is not present.

IV. RESULTS FOR DTO ESTIMATORS

The autocorrelation DTO estimator described in [5] was used as a reference. It uses beamformed data measured using SARUS with TO receive apodization

Fig. 2. Estimated mean velocity profiles for the DTO autocorrelation approach (top plot) and for the cross-correlation method (bottom) at a beam-to-flow angle of 90° using 32 emissions.

and focusing in a line perpendicular to the ultrasound propagation direction. This data is used in the estimator to yield the lateral velocity component. The same beamformed data is used in the new cross-correlation estimator developed in the accompanying paper [2].

A. Estimation for different angles

The performance of the cross-correlation approach compared to autocorrelation DTO can be seen in Fig. 2, where 32 emissions have been used for imaging the vessel at a depth of 64 mm and a beam-to-flow angle of 90° . The true profile is shown as the dashed green line, and the mean of all estimated profiles is shown as the blue curve $\pm$ one standard deviation (SD). Velocity estimates outside the vessel boundary vary widely, as no discriminator has been applied. The overall statistics has therefore been calculated on the central 90% of the vessel to avoid the influence from the echo canceling filter [15], [10] at the boundary and outside the vessel. For DTO the relative SD is 7.24% whereas it is 4.76% for the cross-correlation approach. The bias levels are similar, showing that the new estimator has a slight advantage compared to DTO in performance.

Estimates for a beam-to-flow angle of 60° are shown in Fig. 3. There is, thus, a significant axial component, which is compensated out in the estimator. Again the SD is slightly lower for the new method compared to DTO as shown on the top plots. It should be noted that the same autocorrelation estimator is used for both methods for the axial velocity component.

Fig. 3. Estimated mean velocity profiles for the DTO approach for the axial velocity (top), the lateral velocity using the autocorrelation approach (middle), and the cross-correlation method (bottom) at a beam-to-flow angle of 60° .

B. Increased velocity range

Fig. 4 demonstrates the larger velocity range for the cross-correlation method. The pulse repetition frequency f_{prf} was lowered by a factor of 4 by using every fourth line in the estimation process. Estimates at the vessel center are aliased for the autocorrelation estimator in the top graph, whereas the cross-correlation estimator still can find reliable estimates. The vessel was placed at 65 mm, the beam-to-flow angle was 60° , and 32 emissions were used in the estimator.

The lateral wavelength λ_x used was 2.39 mm and f_{prf} was 2.5 kHz giving an aliasing limit of $v_{z,max} = \lambda_x f_{prf} / 4 = 1.49$ m/s. Reducing f_{prf} by a factor N , thus reduces $v_{z,max}$ by this factor to 0.373 m/s. Fig. 4 shows that this limit applies for the autocorrelation estimator, but the new DTO cross-correlation estimator can still provide the correct velocity estimate.

The bias and standard deviation are heavily influenced by aliasing, and the limit for the new estimator can therefore be determined using the bias increase as shown in Fig. 5. Here the number of skipped lines N is varied resulting in a larger displacement between correlated lines. This is expressed in wavelengths of motions between emissions. The autocorrelation DTO estimator has an aliasing limit of $\lambda_x / 4$ as indicated by the dashed, yellow line. Bias and SD increases for the autocorrelation estimates after this aliasing limit, whereas the bias is relatively constant for the cross-correlation estimator up to $0.75\lambda_x$, three times higher than for the autocorrelation estimator. The SD rises slightly for increasing displacement values due to the increased probability of false detection at higher search

Fig. 4. Velocity profiles as shown on the top graphs for the autocorrelation (left) and cross-correlation (right) DTO estimators, when f_{prf} is reduced by a factor of four to 650 Hz. The mean velocity profiles $\pm$ SD are shown in the bottom graphs.

Fig. 5. Standard deviation (top) and bias (bottom) as a function of displacement between pulse emissions in wavelengths for the autocorrelation (blue) and cross-correlation (red) estimators, when the effective f_{prf} is varied.

ranges [16]. This is shown in Fig. 6, calculated as estimates which deviates more than 1/10 of the peak velocity from the correct velocity. Probabilities below 20% are sparse enough that they can be compensated for either visually or through considering adjacent estimates to correct outliers [17]. This gives a limit around $0.75\lambda_x$ corresponding to 9 m/s, if $f_{prf} = 5$ kHz. The SD is also affected by the false peaks, and this is the main reason for the increase in SD for velocities higher than 5 m/s corresponding to motions beyond $0.4\lambda_x$ per emission.

Fig. 6. Probability of false detection for the cross-correlation estimator.

Fig. 7. Performance as a function of emissions used for the estimates at a beam-to-flow angle of 60° and a vessel depth of 64 mm.

C. Performance comparison

Fig. 7 shows the performance as a function of emissions used for the estimates at a beam-to-flow angle of 60° and a vessel depth of 64 mm. The relative standard deviation is shown on the top and the bias on the bottom graph. The red color indicates estimates from the new cross-correlation approach, and the orange color is for DTO. The relative SD is in all cases lower for the cross-correlation approach demonstrating that a better precision is attained for the same number of emissions using exactly the same beamformed data.

V. RESULTS FOR INTER-LEAVED SA DTO

The sequence described in Section III has been used to acquire data from the flow rig for different values of f_{prf} and for different beam-to-flow angles. Eight high resolution images were used for the inter-leaved sequence. Echo canceling was performed by subtracting the mean value of the HRIs from the data. Three HRIs were created in one sequence by combining all the first positive emission in $+HRI^{(1v)}$, then all negative

emissions in $-HRI^{(2v)}$, and finally all of the last positive emissions in $+HRI^{(3v)}$. Correlation was then conducted between $+HRI^{(1v)}$ and $+HRI^{(2v)}$ and between $+HRI^{(2v)}$ and $+HRI^{(3v)}$. The correlation functions were averaged across all the high resolution images. The scheme maintains the shortest time difference between correlated signals, yields data that can be beamformed everywhere in the images, and is continuous in time. The data shown here spans $8 \times (3 \times 5 + 2) = 136$ emissions, which is comparable to the normal time span of a spectral velocity estimate, which often uses 128 emissions.

The first measurement uses $f_{prf} = 4500$ Hz and the second uses only a 450 Hz f_{prf} . The peak velocity in the vessel is 0.52 m/s. Fig. 8 shows the estimated velocity profiles. The standard deviation of the lateral velocity component is maintained around 2.5% in both cases although the time between correlations is increased by a factor of 10. For a normal SA flow sequence the distance between HRIs to correlate would here be 17 emissions, where it now is only one emission. The largest velocity possible to estimate in the lower sequence would then be above 0.52 m/s and scaling f_{prf} to 4500 Hz gives a maximum velocity of 5.2 m/s. Using even higher pulse repetition frequencies can further increase the maximum detectable velocity. It should also be noted that the data is continuously available in the full imaging region shown in Fig. 1. It is therefore also possible to estimate low velocities as averaging can take place over as long time as the velocity can be considered stationary. This also makes echo canceling much easier, as any echo canceling filter can be employed for the continuous data.

The velocity estimation has been conducted for different pulse repetition frequencies of 225, 450, 1000 and 4500 Hz and for beam-to-flow angles of 60° , 75° , and 90° . Directional data orthogonal to the imaging direction has been beamformed using TO, and data has also been beamformed along the flow directional to enable the use of directional cross-correlation estimation of the velocity. The results for the DTO estimates are shown in Fig. 9 for measured data. A fairly low bias and SD is maintained for both $f_{prf} = 4500$ Hz and $f_{prf} = 1000$ Hz, whereas the results deteriorates significantly for $f_{prf} = 450$ Hz at angles of 60° and 75° . Scaling f_{prf} to 5 kHz for the 1 kHz case corresponds to a peak velocity in the vessel of 2.6 m/s.

The same pattern is seen for simulated data in Fig. 10, although with a slightly lower bias and SD. The maximum detectable velocity can be increased by using directional data focused along the flow direction [18], [19], where the results from simulations are shown in Fig. 11. The SD is now significantly lower around 1% and reliable velocities can be estimated at both 75° , and

Fig. 8. Mean velocity profiles $\pm$ SD for SA VFI using a 4500 Hz f_{prf} (top) and a 450 Hz f_{prf} (bottom) at a beam-to-flow angle of 90° .

Fig. 9. Performance for SA imaging using the DTO cross-correlation estimator as a function of f_{prf} for measured data.

Fig. 10. Performance for SA imaging using the DTO cross-correlation estimator as a function of f_{prf} for simulated data.

Fig. 11. Performance for SA imaging using the directional cross-correlation estimator as a function of f_{prf} for simulated data.

90°. Scaling f_{prf} from 450 Hz to 5 kHz in this case corresponds to a peak velocity in the vessel of 5.8 m/s. At 90° the profile can still be estimated with a relative SD of 2.6% and bias of 5.7% for $f_{prf} = 225$ Hz. For a 10 kHz f_{prf} this translates to peak velocities of 23.1 m/s. For the other angles there is a significant amount of false detections, which potentially could be corrected to yield the true value by employing clustering [17] or regularization [20].

VI. DISCUSSION AND CONCLUSIONS

A new approach for vector flow estimation using transverse oscillation and lateral beamforming in combination with cross-correlation was developed and investigated. A new pre-processing step removes the influence from the axial velocity component and makes it possible to decompose the estimation into an axial and a lateral

cross-correlation estimator. The advantages are a significantly higher maximum detectable velocity and a slightly improved standard deviation on the results improving on the estimates for the same number of data. The estimate is also self-calibrating, as no estimation of the lateral oscillation wavelength λ_x is needed resulting in a low bias. Flow rig measurements using a duplex setup for a convex probe demonstrated that the relative SD is between 4.8% and 8.7% and the bias is between -3.5% and -1.2% at angles from 60° to 90° when using 32 emissions at $f_{prf} = 2.5$ kHz. The new estimator goes beyond the normal aliasing limit of $\lambda_x f_{prf}/4$ and could therefore estimate considerably higher velocities. Peak velocities up to 9 m/s could be estimated at $f_{prf} = 5$ kHz with a probability of false detection around 20%, which is visually acceptable or could be corrected by clustering estimates over time or space to correct false values. The peak detection is, thus, 3 times higher than for the autocorrelation estimator.

An inter-leaved pulse emission scheme for synthetic aperture flow imaging makes it possible to maintain the highest possible maximum detectable velocity in a SA sequence with multiple emissions. This decouples the link between the sequence length N and the maximum detectable velocity. Long sequences with good focusing properties and low side lobes can be used and at the same time maintain the highest possible maximum detectable velocity. Continuous data are obtained, which yield a low standard deviation, arbitrary echo canceling possibilities, and the ability to detect low velocities. The sequence investigated contained a flow sequence inter-leaved with a longer B-mode sequence. The flow sequence used 5 virtual source, where each source was repeated with a positive, negative, and positive polarity emission. This makes it possible to perform velocity estimation and B-mode imaging with pulse inversion for both. Two different velocity estimation schemes were implemented using directional TO and directional beamforming along the flow direction, and they were evaluated using a combination of eight high-resolution images. A satisfactory performance was obtained for both methods using a f_{prf} of 4.5 kHz and 1 kHz. The bias and SD were lower for the simulated data indicating a bit more noise in the measured data. Performance was significantly better for the directionally focused data and SD below 1.5% were obtained with a bias below 5% down to an f_{prf} of 450 Hz. Acceptable results could even be obtained at $f_{prf} = 225$ Hz at a beam-to-flow angle of 90°. This, however, assumes a known beam-to-flow angle, if significantly more beamforming is to be avoided in the directional approach to estimate the angle. A combination of DTO autocorrelation and directional estimation has been made

for plane wave imaging, and this reduces the beamforming demands significantly [21].

These results indicate a peak detectable velocity of at least 2.9 m/s for the DTO method and 5.8 m/s for the directional approach, when $f_{prf} = 5$ kHz is used. This is lower than for the duplex sequence, and this is due to the combination of 5 emissions in SA before the velocity is estimated. Compensation for the motion could also be introduced [22], and this might in certain cases improve on the estimates. The profiles detected for lower values of f_{prf} have high SDs due to many false peak detections. These can be removed using clustering [17] or regularization [20]. The peak velocity can also be increased by using $f_{prf} = 10$ kHz resulting in a maximum detectable velocity of 23.1 m/s for directional beamforming and cross-correlation at 90° . Another approach is to increase the lateral wavelength, which will also proportionally increase the maximum detectable velocity. The axial velocity is more difficult to estimate for high velocities as the wavelength is small and cannot be easily manipulated. For high velocities the transducer should therefore be oriented, so that the main velocity component is along the transverse direction to reduce the axial velocity component.

The demand enforced in cardiology for visualizing the whole heart is $f_{prf} = 5$ kHz to scan down to depths of 15 cm. Regurgitation jets can yield velocities up to 7 m/s and today have to be measured by continuous wave Doppler. This creates ambiguity with regard to depth and angle and, thus, actual value of the velocity. The methods presented here open the possibility of scanning the heart from new angles and yield the true peak velocity and direction of the flow without ambiguity at more than 60 independent frames per second.

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