Reconstructing the fast-ion velocity distribution in the DIII-D tokamak during Alfvén eigenmode activity

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Reconstructing the fast-ion velocity distribution in the DIII-D tokamak
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B. Madsen¹, M. Salewski¹, W. W. Heidbrink², L. Stagner², M. Podestà³, D. Lin²,
A. V. Garcia², P. C. Hansen⁴, J. Huang⁵ and the DIII-D team

¹ Department of Physics, Technical University of Denmark, Kgs. Lyngby, Denmark
² Department of Physics and Astronomy, University of California Irvine, Irvine, CA, USA
³ Princeton Plasma Physics Laboratory, Princeton, NJ, USA
⁴ Department of Applied Mathematics and Computer Science, Technical University of
Denmark, Kgs. Lyngby, Denmark
⁵ Institute of Plasma Physics, Chinese Academy of Sciences, Hefei, Anhui, China

Abstract

Fast-ion velocity-space tomography enables reconstruction of the 2D fast-ion velocity
distribution from a set of measurements. Here, we reconstruct the fast-ion velocity distribu-
tion for plasmas with strong and weak Alfvén eigenmode (AE) activity using the four-view
fast-ion D-alpha (FIDA) diagnostics in the DIII-D tokamak. We find that the fast-ion losses
due to strong AE activity are selective in velocity-space with a particularly strong density
decrease in the population of co-going fast-ions with energies between 40 and 70 keV.

Introduction

Resonant interactions between fast ions and Alfvén eigenmodes (AEs) can lead to enhanced
fast-ion transport [1, 2] that can potentially cause reduced heating efficiency and damage to
the reactor walls. The fast-ion D-alpha (FIDA) diagnostics provides a way to map the fast-
ion transport by measuring the radiation following Balmer-alpha transitions in neutralized fast
deuterium ions [3]. The ions are neutralized by charge-exchange with beam neutrals, and the
measurement volume is determined by the intersection between the FIDA line-of-sight and
the neutral beam. From a set of FIDA measurements, the local fast-ion distribution can be
reconstructed by fast-ion velocity-space tomography [4, 5, 6]. Here, we employ the four-view
FIDA diagnostic to reconstruct the central fast-ion velocity distribution during weak and strong
AE activity during the sequential discharges #153071 and #153072 in the DIII-D tokamak [2,
7, 8].

Velocity-space coverage of the DIII-D FIDA diagnostics

At a location in position space, the discrete FIDA spectrum $S$ relates to the local fast-ion
velocity distribution $F$ through $S = WF$, where $W$ is the transfer matrix containing the FIDA
weight functions [4, 9, 10] computed with FIDASIM [11]. Four distinct FIDA views looking onto three different neutral beams are installed in the DIII-D tokamak [2]. Fig. 1a-d show the modelled spectra for each view originating from an analytic fast-ion slowing down distribution (Fig. 2). The DIII-D FIDA diagnostics measures only one-sided spectra (unlike the FIDA systems at the ASDEX Upgrade [6], NSTX [12] and EAST [13] tokamaks), as indicated by grey-shaded areas in the figure. Fig. 1e shows how many FIDA views are sensitive to any given point in velocity space. Multiple-view coverage is required to obtain reliable reconstructions [14]. However, a large fraction of the negative-pitch velocity space is covered by only one view. Therefore, reliable reconstructions can be obtained for only positive pitches.

Figure 1: (a-d) Synthetic signals from an analytical co-going NBI slowing down distribution for each of the four FIDA views. The experimentally reliable wavelength ranges are marked in grey, whereas the red-shaded regions correspond to synthetic null-measurements [15]. (e) The number of views covering any given point in velocity space related to the grey-shaded areas in panels a-d. Here we trust only reconstructions in the area that is not scratched.

Reconstructions for positive pitches

In order to reconstruct the fast-ion distribution from measurements, the problem must be regularized [5]. This can be done by first-order Tikhonov regularization [16], where the solution is

$$ F^* = \min_F \left\| \begin{pmatrix} W \\ \lambda_1 L_1 \end{pmatrix} F - \begin{pmatrix} S \\ 0 \end{pmatrix} \right\|_2. $$

(1)

Here $L_1$ approximates the gradient with respect to $E$ and $p = v_\parallel/v$ [5], and $\lambda_1$ is the regularization strength. The solution is improved by including a non-negativity constraint [15] and null-measurements marked with red-shaded regions in Fig. 1a-d as a penalty [17]. Additionally, we introduce two new types of prior information for reconstructing for only positive pitches by assuming: (i) a known simulated signal originating from ions with negative pitches, and
(ii) isotropy for negative-pitch ions deposited in a one-dimensional energy-resolved bin. The resulting reconstructions from a noisy synthetic signal based on the analytical slowing-down distribution in Fig. 2a are shown in Fig. 2b-e. For all methods, the positive-pitch regions of the reconstructions are in good agreement with the ground truth.

Figure 2: (a) Analytic fast-ion slowing down distribution from co-going neutral beam injection and (b-d) inversions of a synthetic signal from panel a. The superscripts indicate the employed inversion methods, and the normalization factors are the same.

The effect of AE activity on the positive-pitch fast-ion population

The DIII-D discharges #153071 and #153072 exhibit weak and strong AE activity, respectively [2, 7]. Due to calibration uncertainties, we only consider relative differences caused by increased mode activity. Fig. 3 shows the pitch spectra integrated over energies for \( E > 30 \text{ keV} \) and energy spectra integrated over positive pitches for the pixel-differences between the central fast-ion velocity distributions during strong and weak AE activity. The spectra are computed both from the classically expected TRANSP/NUBEAM [18] distributions and reconstructions from measurements using the methods tested in Fig. 2. Compared to what is classically expected, the reconstructions observe a strong decrease in the local fast-ion density for energies between 40 and 70 keV and pitches greater than 0.2 caused by the increased AE activity.

Figure 3: (a) Pitch spectra integrated over energy for \( E > 30 \text{ keV} \) and (b) energy spectra integrated over pitch for \( p > 0 \) for the differences between the discharges with strong and weak AE activity from both the classical TRANSP/NUBEAM distributions and reconstructions.
Conclusion

The discharge with strong AE activity shows a strong decrease in the density of fast ions from the central plasma with pitches greater than 0.2 and energies between 40 and 70 keV compared to the discharge with weak AE activity. This is captured by all methods used to reconstruct the distribution from measurements, but not by the classical TRANSP/NUBEAM distributions, and is attributed to the increased AE activity.

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