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Anomalous dispersion of longitudinal optical phonons in oxygen-doped \( \text{La}_{2-x}\text{Sr}_x\text{CuO}_{4+d} \)

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Inelastic neutron scattering has been used to study the in-plane Cu-O bond-stretching mode in oxygen-doped \( \text{La}_{1.94}\text{Sr}_{0.06}\text{CuO}_{4.035} \) \( (T_c = 38 \text{ K}) \) and \( \text{La}_{2.03}\text{CuO}_{4.05} \) \( (T_c = 43 \text{ K}) \). Similar to results from optimally doped \( \text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4 \) \( (T_c = 35 \text{ K}) \), we observe anomalous features in the dispersion of this half-breathing mode in the form of a softening halfway through the Brillouin zone. Considering the differences in electronic structure and local environment between the oxygen- and strontium-doped compounds with similar \( T_c \), we rule out a connection between the phonon anomaly and structural instabilities related to the specific dopant type. We interpret the phonon anomaly as a signature of correlated charge fluctuations ubiquitous in optimally doped superconductors.

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A widely accepted description of the low-temperature electronic state of underdoped cuprate superconductors is the “stripe” picture, where, due to hole doping, the antiferromagnetic (AFM) ground state of the parent compound is segregated by channels of charge resulting in magnetic antiphase boundaries [1]. Experimentally, the magnetic part of the stripes shows up in neutron-scattering experiments as a modulated AFM structure (static magnetic stripes) and excitations emerging with similar wave vectors, in the low-energy regime (dynamic magnetic stripes).

While experimental evidence of magnetic stripes has been extensively documented (see, e.g., [2] for a review), the charge component of the stripes is more elusive. Static charge stripes only show up in superconducting (SC) samples close to the \( x = \frac{1}{8} \) anomaly \([1,3–6]\) and evidence for dynamic charge stripes has been reported for \( \text{La}_{2-x}\text{Sr}_x\text{CuO}_4 \) \([7,8]\), \( \text{YBa}_2\text{Cu}_3\text{O}_{7-\delta} \) \([9]\), and \( \text{La}_{2-x}\text{Sr}_x\text{NiO}_4 \) \([10]\). Understanding the relationship between these signatures of stripe formation as a function of doping would be a crucial leap forward in our understanding of the cuprates.

Recently, it was discovered that the dispersion of the Cu-O bond-stretching longitudinal-optical (LO) phonon in SC \( \text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4 \) (LSCO15, \( T_c = 38 \text{ K} \)) \([11]\) and \( \text{La}_{1.80}\text{Sr}_{0.20}\text{CuO}_4 \) \( (T_c \approx 35 \text{ K}) \) \([12]\) displays a strong anomalous softening interpreted as a coupling to a novel charge collective mode \([12]\). Furthermore, merely a weak signature of the anomaly is visible in the phonon linewidth of \( \text{La}_{1.93}\text{Sr}_{0.07}\text{CuO}_4 \) \( (T_c \approx 15 \text{ K}) \) and \( \text{La}_{1.73}\text{Sr}_{0.23}\text{CuO}_4 \) \( (T_c \approx 15 \text{ K}) \), suggesting that the strength of the anomaly tracks the doping dependence of \( T_c \) \([12]\). Similar phonon anomalies have been observed in \( \text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_{8+\delta} \) \([13]\), \( \text{La}_{1.87}\text{Ba}_{0.125}\text{CuO}_4 \) (LBCO) \([14]\), \( \text{La}_{1.48}\text{Nd}_{0.5}\text{Sr}_{0.12}\text{CuO}_4 \) (LNSCO) \([11]\), and \( \text{YBa}_2\text{Cu}_3\text{O}_{6.6} \) \([15]\) hinting at a ubiquitous feature of cuprate superconductors.

In order to further investigate the robustness of the phonon anomaly in the \( 0.1 < x < 0.2 \) doping range, we examine two compounds derived from \( \text{La}_2\text{CuO}_4 \) by two distinct doping procedures giving rise to unique magnetic, structural, and superconducting properties.

Hole doping of the parent compound \( \text{La}_2\text{CuO}_4 \) can be performed by introducing two distinct dopant species. Replacing \( \text{La}^{3+} \) for \( \text{Sr}^{2+} \) yields \( \text{La}_{2-x}\text{Sr}_x\text{CuO}_4 \) (LSCO) with “quenched doping” \([16]\), meaning that Sr has a fixed random distribution on La sites after crystal growth. On the other hand, an “annealed doping” \([16]\) can be obtained by introducing excess, mobile oxygen anions into the lattice by electrochemical methods \([17]\), obtaining \( \text{La}_{2-x}\text{Sr}_x\text{CuO}_4 \) (LCO+O). At low Sr concentrations \((x \leq 0.14)\) it is possible to combine Sr and O dopants, resulting in “codoped” \( \text{La}_{2-x}\text{Sr}_x\text{CuO}_{4+d} \) (LSCO+O) \([18]\). Figure 1(a) depicts the crystallographic sites of O/Sr dopant ions \([19–21]\).

Quenched \( \text{Sr}^{2+} \) doping creates a superconductor where \( T_c \) varies continuously with doping, forming the so-called superconducting dome for \( 0.05 \leq x \leq 0.25 \). Meanwhile the mobile O dopants only seem to allow for certain superconducting phases to emerge \((T_c \approx 15, 30, 40 \text{ K})\) due to oxygen content \([22]\), pressure \([23]\), or thermal treatment \([24]\).

At sufficient O doping, LSCO+O is superconducting at ambient pressure below \( T_c \approx 40 \text{ K} \), regardless of Sr content \([18]\). While the underlying mechanism for the connection between annealed disorder and superconductivity is far from settled, the distribution of O dopants is clearly distinct from the distribution of \( \text{Sr}^{2+} \) dopants. Contrary to the quenched \( \text{Sr}^{2+} \) dopants, annealed O dopants in LSCO+O are responsible for a number of structural ordering phenomena such as staging \([25,26]\) and fractal-like distributions of ordered superstructure patches \([24,27]\).

Differences can also be observed in the magnetic structure of the two compounds. In LSCO+O, the transition to static
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The extracted dispersion from the zero-field data is shown in Fig. 3 along with a normal sinusoidal dispersion, $\omega = \alpha \cos(2\pi q) + \beta$, inferred from phonon calculations on LSCO using density functional theory (see Fig. S5 in the SM [34] and Ref. [38]). We fit the cosine function to points near the zone center [$q = (0, 0, 0)$] and edge [$q = (\frac{1}{2}, \frac{1}{2}, 0)$] to obtain the dashed curves of Fig. 3(a). To quantify the magnitude of the anomaly, we define the “anomaly signal” as the difference between the normal dispersion and the measured data [gray shaded area in Fig. 3(a)]. Figure 3(b) shows our anomaly signal for LCO+O and LSCO6+O along with previous results from optimally doped LSCO15 and insulating, stripe-ordered LNSCO [11]. We emphasize the presence of similar anomaly signals on an absolute scale across all studied samples. The linewidth broadening $\gamma$ follows the softening of $\omega_q$ similar to what was observed in LSCO15 and LNSCO [11] (see Fig. S2 in the SM [34]). Finally, our data can be described by a linear combination of stripe-ordered and optimally superconducting anomaly signals (see Fig. S7 in the SM [34]), consistent with the observation of phase separation in the (super)oxygenated compounds [18].

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FIG. 2. Reduced data at selected wave vectors of the form $Q = (h, h, 0)$ for both LSCO6+O and LCO+O at $T = 5$ K and $H = 0$ T. Data at $Q = (5, 5, 0)$ and $Q = (4.85, 4.85, 0)$ was scaled by a factor of $\frac{1}{2}$ for clarity due to an increase of intensity from the phonon form factor. Data at different $h$ are offset for clarity. Solid lines are fits to a DHO line shape (see text).

To begin the discussion of our results, we remark that softening and/or broadening of phonon modes is generally a signature of an incipient structural or electronic instability. Typical examples include structural phase transitions, $q$-dependence of the electron-phonon matrix element, Fermi surface nesting, and electronic correlations [39]. In order to determine the origin of a given phonon anomaly, it is therefore important to carefully exhaust alternatives before making statements about the connection to novel phases such as dynamic charge stripes [39].

The phonon anomaly appears in the vicinity of the wave vector $q_{co} = (\frac{1}{4}, \frac{1}{4}, 0)$, consistent with charge-stripe ordering as illustrated in Fig. 1. Measurements of LNSCO [40], LSCO15 [11], and LBCO [11] have shown a suppression of the anomaly as one moves away from the bond-stretching direction [40], supporting a one-dimensional stripelike picture. Additionally, the phonon anomaly in LSCO15 and LBCO has almost no temperature dependence apart from a slightly sharper peak shape when heating from 10 to 300 K [11,14].

These phenomena rule out anharmonicity and structural inhomogeneity as mechanisms for the phonon anomaly in these systems.

A combination of inelastic x-ray and angle-resolved photoemission spectroscopy (ARPES) measurements on overdoped LSCO ($x = 0.2$ and $x = 0.3$) have shown that the phonon anomaly wave vector is inconsistent with Fermi surface nesting [12,41], contradicting the idea of a phonon softening due to a Kohn anomaly. A different, possibly $q$-dependent, electron-phonon coupling could still be responsible for the phonon anomaly. Such an effect would renormalize the electronic quasiparticle dispersion (the so-called “ARPES kink” [42]) at energies similar to the phonon softening. The ARPES kink has been observed in LSCO $x = 0.2$ and $x = 0.3$, but since only LSCO $x = 0.2$ shows anomalous phonons, the two phenomena appear to not be connected [12].

Thus, all previous studies are unable to explain the phonon anomaly through conventional means and any coupling to stripe order is likely dynamic. One possible scenario is a coupling of the Cu-O bond-stretching phonon with steeply dispersing charge fluctuations. Kaneshita et al. performed calculations based on the Hubbard model of this scenario, predicting anomalous phonon dispersions due to both transverse (meandering) and longitudinal (compression) coherent stripe fluctuations [28] (see Fig. 1 for a sketch of the transverse mode). We emphasize that the observed phonon anomaly reported here (see Fig. 3) and in LBCO/LSCO [11,14] is remarkably similar to the calculation by Kaneshita et al. (see Fig. 5 in Ref. [28]).

Despite differences in the magnetic excitation spectra as recorded by neutron scattering (including low- and zero-energy transfers), the three materials LCO+O, LSCO6+O,
and LSCO15 have remarkably similar in-plane Cu-O bond-stretching dispersions (Fig. 3). Furthermore, static charge order at zero field has been observed in a different sample of LCO+O [43] and in LNSCO [3] but so far not in LSCO6+O nor in LSCO15. These observations together rule out a unique, direct connection between static stripes (spin and charge) and the phonon anomaly. In order to further confirm this point, we performed scans of LSCO6+O at selected wave vectors in a $H = 10\ T$ magnetic field which is known to induce a considerable volume of striplike magnetic order in this particular sample [33]. While static charge order has not been observed in LSCO6+O, measurements on LSCO ($x = 0.12$) have shown that static charge and spin stripes respond identically to magnetic fields [4].

Figure 4 contains data at two wave vectors with and without an applied magnetic field of 10 T, clearly showing the absence of any detectable field effect on the in-plane Cu-O bond-stretching phonon at $\mathbf{q} = (\frac{1}{4}, \frac{1}{4}, 0)$. These measurements were performed simultaneously with measurements of the low-energy magnetic fluctuations [44], confirming a significant increase in the magnetic spectral weight towards lower energies consistent with the appearance of field-induced stripe order [33]. Thus, the appearance of static magnetic stripe order does not affect the phonon anomaly in LSCO6+O. A similar insensitivity of the phonon anomaly to an applied magnetic field has been observed in underdoped ($T_c = 66\ K$) YBa$_2$Cu$_3$O$_{6.6}$ [45].

We have shown that the phonon anomaly is a robust feature in optimally doped as well as stripe-ordered cuprates which is independent of the structural details related to the doping process. Since it is equally well formed in stripe-ordered and optimally doped systems, where the latter show no static magnetic order, the anomaly is surprisingly insensitive to low-energy magnetic characteristics. This is further confirmed by the absence of a magnetic field effect in LSCO6+O which introduces static magnetic stripe order.

The phonon anomaly is strongest in the doping region around optimal $T_c$ (0.125 $\leq n_0 \leq 0.20$) (LSCO15, LNSCO [11], LBCO [14], LSCO6+O, LCO+O), regardless of the presence of static charge order (LNSCO [1], LBCO [46]), suppression of bulk superconductivity (LNSCO [3]) or dopant disorder (LCO+O, LSCO6+O). In addition, the phonon anomaly is unaffected by magnetic fields (LSCO6+O, YBa$_2$Cu$_3$O$_{6.6}$ [45]) and temperature (LBCO, LSCO15 [11]).

Thus it appears to be an intrinsic, robust signature of doped cuprates near optimal doping. We point out that the phonon anomaly region coincides with the dome of dynamic charge fluctuations recently reported in Ref. [9]. This points to fluctuating charge stripes as the fundamental degrees of freedom in the cuprates and thereby lends credibility to the “electronic liquid crystal” model [47].

In conclusion, we have measured the in-plane Cu-O bond-stretching phonon in LSCO6+O and LCO+O and provided evidence for significant anomalous behavior. Since one sample (LCO+O) exhibits charge order [43] while the other (LSCO6+O) does not [33], and since the samples also have different magnetic spectra with distinct field dependencies [33,48] we conclude that the phonon anomaly has no direct, trivial relationship to either magnetic or charge static order. In addition, the unique structural characteristics of oxygen-doped samples rule out a connection between the specific dopant species and the phonon anomaly. We proceed to conclude that the phonon anomaly is a signature of charge stripe fluctuations, which is a common characteristic of the cuprate family and appears to be a prerequisite to optimal superconductivity in these systems.

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