A Water-Based Huygens Dielectric Resonator Antenna

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ABSTRACT

With its high-permittivity in the microwave frequency range, water has the potential as an alternative, inexpensive, bio-friendly and abundant material for many microwave applications such as dielectric resonator antennas and tunable metamaterials. Huygens antennas, composed of single-element structures supporting a special combination of electric and magnetic modes, provide an alternative route for compact and directive antennas. In this work, we investigate a subwavelength water-based Huygens dielectric resonator antenna with strongly excited electric and magnetic dipoles at around 350 MHz. Our antenna leads to the directivity of 6 and a radiation front-to-back ratio of 40.3 dB. Additionally, good matching to the feed-line and the surrounding free space were achieved with a reflection coefficient of $-38.3\,\text{dB}$ and a total efficiency of 57.8%. The antenna was fabricated and characterized with excellent agreement between the measurements and the numerical results. Furthermore, several means of tuning the antenna were tested numerically and experimentally. We believe that the proposed water-based Huygens dielectric resonator antenna may serve as an easy-to-fabricate and cheap alternative for the VHF and low end of the UHF bands.

INDEX TERMS

Dielectric resonator antenna, reconfigurable antennas, water.

I. INTRODUCTION

OVer several decades, dielectric resonator antennas (DRAs) have been of great interest as compact alternatives to traditional dipole and waveguide antennas [1]. Many designs have been proposed demonstrating different attributes such as low-profile, large bandwidth, frequency-tunability and phase-steering capabilities. Generally, DRAs rely on electric and magnetic modes supported in expensive high-permittivity ceramic structures. Alternatively, water-filled cavities exploiting water’s relatively high permittivity in the microwave frequency range can be used instead [2]. Water holds great potential as an abundant, inexpensive and bio-friendly material for tunable antennas [3]–[5], as well as inclusions in metamaterials (MMs) and metasurfaces (MSs) [6]–[11].

Various liquid-based antennas have been demonstrated with some relying on water’s high real and low imaginary parts of permittivity, whereas other use conductive liquids [3]–[5], [12]. In general, many of these antennas operate in the VHF and UHF bands due to the decreasing (increasing) real part of water’s permittivity at higher frequencies. The working principle of water-based DRAs relies on resonant electric (TM) and magnetic (TE) modes in water-filled cavities. These modes are the natural eigen-modes belonging to the dielectric cavity, and they can be excited by sources placed outside and inside the cavity. In general, a single cavity has multiple modes spread-out in narrow bands of frequencies across the spectrum. In [5], we demonstrated an electrically small water-based hemispherical DRA with such characteristics. The water provided an excellent way to place an adjustable metallic monopole antenna inside the cavity without the issues of air gaps between the dielectric and the monopole antenna. Furthermore, the water-based DRA could be tuned by changing the water-temperature and/or the water-volume in the cavity.
In antenna configurations, the dipole modes are the most broadly used, which by their nature have low directivity. Higher-order modes can also be exploited for high-directivity antennas, but such higher-order modes require larger dielectric structures, have very small bandwidths and are more sensitive to losses [1]. Another way to increase the directivity in a compact single-element configuration is to combine magnetic and electric modes [13]–[15], and such antennas are often referred to as Huygens antennas. Furthermore, when the fundamental dipole modes are combined in dielectric structures, Kerker’s condition may be satisfied [16]–[17], which is also of great interest in MM's and MS's. Additionally, dipoles and higher-order modes can also be combined to increase the directivity even further [17]–[18]. Kerker’s condition is satisfied in many types of dielectric structures, including dielectric spheres and cylinders [19]–[20]. However, the dipole modes are often weakly excited resulting in very low scattering. Designing structures exhibiting strongly excited magnetic and electric modes with the right integration of their magnitudes and phases is not a simple task. Combining an electric and magnetic dipole will increase the directivity in one direction to twice of that of a single dipole, while in the opposite direction, the directivity is greatly minimized. Several Huygens antenna designs have already been demonstrated [13]–[14], [18], but thus far no water-based Huygens DRAs have been proposed.

The purpose of the present work is to investigate a simple subwavelength water-based Huygens DRA. The antenna consists of a short monopole antenna fed against a large ground conducting plane and encapsulated by a rectangular cuboid-shaped cavity filled with distilled water. The DRA is designed with overlapping fundamental electric and magnetic dipole modes at around 350 MHz. The final design has a directivity of 6 and a front-to-back ratio (FBR) of 40.3 dB. Furthermore, our antenna is well matched to both the free-space wavelength (wave number). A prototype of the UHF bands.

Throughout this article, the time-factor \(\exp(j\omega t)\), with \(\omega\) being the angular frequency and \(t\) being the time, is assumed and suppressed.

**II. ANTENNA CONFIGURATION AND PARAMETERS**

A sketch of the water-based Huygens DRA is shown in Fig. 1. It consists of a rectangular cuboid-shaped Rohacell 51 HF cavity attached to a ground conducting plane and filled with distilled water. The Rohacell 51 HF materials has a measured relative permittivity of 1.075 [9]. A Cartesian coordinate system is introduced as shown in Fig. 1 with the associated spherical angular coordinates \((\theta, \phi)\). A short monopole is inserted into the cavity as shown in Fig. 1. The feed line is a coaxial transmission line with the inner and outer diameter of 1.28 mm and 4.1 mm, respectively, and a dielectric material with the relative permittivity of 2.1. The characteristic impedance of the feed line is 48.15 \(\Omega\). The diameter and the length of the monopole is 1.6 mm and \(l_m\), respectively. The monopole is displaced by \(x = -x_m\) from the center of the Huygens DRA to excite the magnetic mode. The Huygens DRA is designed for 350 MHz operation with a water temperature of 20°C. The final design with the highest total efficiency \((\eta_{tot})\) and FBR were determined through a parametric study. In a previous work, we investigated a water-based hemispherical DRA, which exhibited both magnetic and electric dipole resonances, however, with the latter being at higher frequencies [5]. A cubic DRA has similar behavior, and to shift the electric resonance to lower frequencies, we simply extend the dielectric along the y- and z-axis until we obtain the wanted properties of the antenna. The cavity size and the values of \(x_m\) and \(l_m\) are listed in Table 1 for the final antenna design.

The \(xz\)- and \(xy\)-planes correspond to the E- and H-planes of the antenna, respectively. The Debye model is used to describe the permittivity of distilled water inside the cavity [2], and at 350 MHz and 20°C, the relative permittivity is \(\varepsilon_r, \text{water} = 80.2 - j1.6\). A model of the Huygens DRA is built in COMSOL Multiphysics 5.4 used for all numerical calculations [21]. The COMSOL model consists of the Huygens DRA enclosed by a hemisphere of free space and with a Perfect Electric Conductor (PEC) plane as the ground. The Huygens DRA is simulated without the Rohacell 51 HF material as its low permittivity minimally affects the antenna response. An outer layer of PML is attached to the free-space hemisphere and a matched port is placed on the bottom of the feed line. A symmetry plane was placed on the bottom of the model size and a convergence test was conducted leading to total of 1,283,616 Number of Degrees of Freedom solved for in the final model.

We use different well-known antenna parameters in this article, such as total efficiency (including the mismatch loss)
TABLE 1. Simulated water-based Huygens antenna parameters.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Cavity size [mm²]</td>
<td>116.8×52.2×132.9</td>
</tr>
<tr>
<td>(l_m) [mm]</td>
<td>27.2</td>
</tr>
<tr>
<td>(x_m) [mm]</td>
<td>24</td>
</tr>
<tr>
<td>(f) [MHz]</td>
<td>351.6</td>
</tr>
<tr>
<td>(G_{\text{realized, max}}) [dBi]</td>
<td>5.42</td>
</tr>
<tr>
<td>(D_{\text{max}})</td>
<td>6</td>
</tr>
<tr>
<td>(\eta_{\text{tot}})</td>
<td>57.8 %</td>
</tr>
<tr>
<td>(\eta_{\text{rad}})</td>
<td>57.8 %</td>
</tr>
<tr>
<td>(</td>
<td>S_{11}</td>
</tr>
<tr>
<td>(Z_A) [Ω]</td>
<td>47.27 + j0.77</td>
</tr>
<tr>
<td>(\text{FBW}_{10dB})</td>
<td>4.0 %</td>
</tr>
<tr>
<td>(\text{FBR}_{\text{max}}) [dB]</td>
<td>40.3</td>
</tr>
<tr>
<td>(k_0a) [rad]</td>
<td>1.08</td>
</tr>
<tr>
<td>(\varepsilon_{r,\text{water}}) (350 MHz, 20 °C)</td>
<td>80.2 − j1.6</td>
</tr>
</tbody>
</table>

\(\eta_{\text{tot}}\), radiation efficiency \(\eta_{\text{rad}}\), reflection coefficient \(S_{11}\) etc., and we refer to [5] or [22] for more information about these parameters. Throughout the paper we use the front-to-back ratio (FBR) to evaluate how well the radiated power is concentrated in one direction. The FBR is defined as

\[
\text{FBR} = \frac{D_{+x}}{D_{-x}},
\]

where \(D_{\pm x}\) is the directivity in the \(\pm x\)-direction (in angular coordinates: \((\theta, \phi)_{+x} = (90°, 0°)\) and \((\theta, \phi)_{-x} = (90°, 180°)\)).

III. NUMERICAL AND EXPERIMENTAL RESULTS
A. ANTENNA CHARACTERIZATION

The final results are summarized in Table 1. We present \(\eta_{\text{tot}}\) and \(\eta_{\text{rad}}\) as functions of frequency in Fig. 2(a). The magnitude of \(S_{11}\) as function of frequency is shown in Fig. 2(b). Around 350 MHz the efficiencies are high, and \(S_{11}\) is low. At 351.6 MHz, where FBR is maximum (see Fig. 4(a)), the \(S_{11}\) drops to \(-38.3\) dB indicating excellent matching of the antenna to the feed-line. The antenna input impedance \(Z_A\) is calculated using \(S_{11}\), and its real and imaginary parts are shown in Fig. 2(c) as functions of frequency. At 351.6 MHz, \(Z_A = 47.27 + j0.77\)Ω. \(\eta_{\text{tot}}\) is modest and is 57.8 % as the water absorbs 42.2 % of the input power. The absorbed power was calculated in COMSOL by integrating the power loss density over the volume of water.

The peak in \(\eta_{\text{tot}}\) comes from the modes excited in the water-filled cavity. From Fig. 2, we would conclude that a single resonant mode is excited around 350 MHz. However, inspecting the near fields of the antenna in both the \(xz\)- and \(xy\)-plane, which are shown in Fig. 3, we observe two dipoles being excited. The displacement of the monopole along the \(x\)-axis inside the water cavity renders an asymmetric field distribution, and in the \(xz\)-plane in Fig. 3(a), we see that the electric field (arrows) forms a half-loop with an enhanced magnetic field intensity (colors) near the ground plane inside the water. This is identified as a magnetic dipole mode (\(\text{TE}_{111}\)). If we turn our attention to the \(xy\)-plane in Fig. 3(b), we find an alternated field distribution: the magnetic field (arrows) forms a loop with an enhanced electric field intensity (colors), i.e., an electric dipole (\(\text{TM}_{011}\)). In our previous work on water-based antennas, such strong dipole modes were also excited, but at different frequencies [5]. In the present configuration, the dipoles are excited at the same frequency. The outcome of the double dipole excitation can also be observed in the far-field, see Fig. 4(a), where the directivity pattern has the cardioid shape. The directivity in the \(-x\)-direction is greatly...
reduced, whereas the maximum in the +x-direction is still maintained. Such behavior is highlighted in Fig. 4(b) with the FBR shown as a function of frequency. At 351.6 MHz, the FBR has the remarkable peak value of 40.3 dB. The maximum directivity $D_{\text{max}} = D_{+x}$ as a function of frequency is included in Fig. 4(b). At 351.6 MHz, $D_{\text{max}} = 6$, which is expected for ground-supported Huygens antennas [15]: a single dipole antenna has a directivity of 1.5, and with two dipoles and a ground plane, the directivity is increased by a factor of 4. Furthermore, the maximum realized gain is $G_{\text{realized, max}} = 5.42 \, \text{dBi}$, which is 0.65 dB higher than that of a matched and lossless monopole antenna.

The antenna was fabricated at the local workshop, and a photograph is shown in Fig. 5(a). The antenna was made by hollowing out the cavity in a Rohacell 51 HF block, which was then glued on to a circular aluminum plate of 1 m in diameter. Holes for insertion of the monopole antenna and distilled water were drilled in the aluminum plate. A sketch of the measurement setup is shown in Fig. 5(b). The antenna was mounted on a tripod with a rotating joint such that the antenna could be rotated $\theta'$ as shown in Fig. 5(b). In this way, the radiation pattern was recorded using a simple dipole antenna designed for 350 MHz operation and positioned 3.3 m away. Accordingly to the definition of the far-field distance, if we use the ground plane as the largest dimension ($d = 1 \, \text{m}$ is the size of the ground plane), we get $d_{\text{ff}} > 2d^2/\lambda_0 \approx 2.3 \, \text{m}$ and thus fulfill the far-field condition. For measurement angles $\theta' = [90^\circ, 0^\circ]$, the antenna was rotated $180^\circ$ in the xy-plane. The reflection coefficients ($S_{11}$) and transmission coefficients ($S_{21}(\theta')$) were measured with an Anritsu MS2024B vector network analyzer calibrated for 50 Ohm matching. The measured $S_{11}$ is shown in Fig. 2(b) and the measured FBR is included in Fig. 4(b), which is calculated as $S_{21}(\theta' = 90^\circ)/S_{21}(\theta' = -90^\circ)$. The measured FBR is maximal at 354.7 MHz with a measured $S_{11}$ of $-30.9 \, \text{dB}$ showing nice agreement with the numerical predictions. The small mismatch of 3.1 MHz in frequency between experimental and numerical results can be explained by a small difference in temperature and/or water filling in the cavity. The normalized measured $S_{21}(\theta')$ in the xz-plane is included in Fig. 4(a), and is in good agreement with the numerical results.

By default, the antenna bandwidth is characterized by its 3-dB fractional matched voltage standing wave ratio bandwidth $\text{FBW}_{3\,\text{dB}}$ [22]. However, due to the second resonance around 390 MHz, this is quite large (19.7 %). Therefore, it is more convenient in this case to use the 10-dB fractional matched voltage standing wave ratio bandwidth $\text{FBW}_{10\,\text{dB}}$, which is 4.0 %. According to [23], the quality factor $Q$ is given by $Q = 2/3\text{FBW}_{10\,\text{dB}}$, which in the case of our Huygens DRA gives around 17. In the absence of water, one is left with the monopole antenna alone, which resonates
at a higher frequency (around 2.8 GHz) with an associated quality factor of around 7 [22]. From Fig. 4(b), we see that the high FBR is narrow-banded. Obviously, this is coming from the tight overlap of the resonant dipole modes. The FBR\textsubscript{10dB} bandwidth is 6.5 %.

**B. TUNABILITIES AND SENSITIVITIES**

Water-based antennas have multiple tuning parameters, and in this work, we have tested variations of the position of the monopole, water content in the cavity and temperature of water.

In total three holes were drilled for the insertion of the monopole, and the measured S\textsubscript{11} for each of them is shown in Fig. 6(a). For the alternative monopole positions, x\textsubscript{m} = 0 and x\textsubscript{m} = 12 mm, we find several minima spread out over their spectra, each representing different mode excitations. None of these modes produce reflection coefficients below −20 dB, and their radiation patterns are different from that shown in Fig. 4(a). In Fig. 7(a), we have included the simulated realized gain G\textsubscript{realized} as a function of \(\theta\) in the \(xz\)-plane as the monopole antenna position is shifted along the \(x\)-axis. The inset shows the angle of each minima (null) in G\textsubscript{realized}, \(\theta\)\textsubscript{min}, as a function of \(x\)\textsubscript{m}. At \(x\)\textsubscript{m} = 0, the radiation pattern is symmetric with the null being at 0°. As \(x\)\textsubscript{m} increases, the null shifts towards the negative angles, and at \(x\)\textsubscript{m} = 24 mm, the null reaches −90°. These results demonstrate the beam-steering capabilities of the antenna with a simple re-positioning of the monopole antenna. Similar study was conducted in [3] showing same tendency. Additionally, further beam-steering should be possible by also displacing the monopole antenna in the \(y\)-direction.

Another way to tune the water-based antennas is to simply remove/add water (see, e.g., [4]–[5]). Presently, we stepwise remove 2 % of the water from the cavity, until 70 % water is left, measuring S\textsubscript{11} after each water extraction. The result is shown in Fig. 6(b) and an aggravation of S\textsubscript{11} is observed as the water content in the cavity is reduced (please note that the measurements in Fig. 6(a) and Fig. 6(b) were carried out at days with different temperatures causing the minor differences in the measured S\textsubscript{11}). A similar study has previously been conducted for another water-based DRA, where a linear blue-shift was observed [5]. However, we do not observe similar tendency for the present design. As the water content is reduced, the radiation in the \(xz\)-plane becomes less directive and more omnidirectional as shown in Fig. 7(b) by G\textsubscript{realized} as a function of \(\theta\) in the \(xz\)-plane. The antenna has the response shown in Fig. 6(b) and Fig. 7(b) as long as it is kept in level. If any deviations in angle theta in positioning of the AUT happens, then water no more can be accurately described as a cuboid dielectric. This is quite complicated behavior as the resonance frequency and radiation pattern will change for different theta angles. This can be a subject of further analysis. Seen from another perspective, this presents a great way to, e.g., completely level the antenna or to check if the cavity is 100 % filled. In addition,
the antenna should be sensitive to vibrations. We have not made such comprehensive studies this work.

Water’s complex permittivity is temperature-dependent, and it decreases with increasing temperature. This property can also be exploited to tune the response of water-based antennas as well as other types of water-based devices [4]. Presently, we simulate minor variations of water’s temperature, and in Fig. 8, $S_{11}$, FBR and $\eta_{\text{tot}}$ are shown as functions of frequency. As the temperature increases from 10°C to 30°C, the response is blue-shifted by approximately 16 MHz. Furthermore, both $\eta_{\text{tot}}$ and FBR improve due to the lower losses, whereas the matching, indicated by the minimum of $S_{11}$, is optimum around 20 °C. Still, the minimum of $S_{11}$ stays below −20 dB at all simulated temperatures.

The results in Fig. 6, 7 and 8 demonstrate the antenna’s tunabilities, but it can also be viewed as its sensitivities to fabrication variations and defects as well as changes in the local environment. The antenna has other tunabilities/sensitivities as, e.g., the impurities in the water, which we did not test in this work.

C. FURTHER DISCUSSIONS

The antenna’s electrical size is $k_0 a \approx 1.08 (\lambda_0/6)$, and therefore, not categorized as being electrically small, which requires $k_0 a \leq 0.5$ for ground-supported antennas, see, e.g., [5] and [24]. Still, the antenna is several times smaller than many other DRAs [1], [4], and the modest efficiency is higher than, e.g., a directive subwavelength DRA composed of an expensive low-loss ceramic material [18]. In our case, we have a simple and cheap antenna. Furthermore, the antenna is tunable and self-matched to a 50 Ohm transmission line, and thus, no matching components are needed. The directivity is twice that of traditional single-unit DRAs in which only magnetic or electric modes are excited [1]. Due to the losses in water and the compact design, the efficiency is modest compared to other DRAs.

In Fig. 2, we find a second mode being excited around 385 MHz. The mode is weakly excited, which explains the low efficiency. Inspecting the field distribution and radiation pattern (not shown in the manuscript), we conclude that it is another magnetic mode.

IV. SUMMARY AND CONCLUSION

A simple water-based Huygens DRA consisting of a short monopole antenna fed against a large conducting ground plane and encapsulated by a rectangular cuboid-shaped cavity filled with water was investigated. With strongly excited magnetic and electric dipoles, an excellent front-to-back ratio of 40.3 dB was achieved as well as good matching to the feed-line and the surrounding space with a reflection coefficient of −38.3 dB and a total efficiency of 57.8 %. Furthermore, the achieved maximum directivity was 6 in accordance with the expectations. The antenna was fabricated and characterized with the measurements exhibiting...
excellent agreement with the numerical results. We demonstrated several ways of tuning the Huygens DRA’s response with both numerical and experimental studies.

With its simple and versatile, as well as extremely cheap and bio-friendly design, the Huygens DRA may serve as an alternative antenna for microwave communication systems operating in the VHF and low end of the UHF bands. Additionally, the Rohacell material is not essential for the functionality of the antenna and can therefore be replaced by cheaper and more bio-friendly alternatives to even further minimize the ecological footprint of the device.

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REFERENCES