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ABSTRACT
The specific far-infrared spectral signatures associated with highly localized large-amplitude out-of-plane librational motion of water molecules have recently been demonstrated to provide sensitive spectroscopic probes for the micro-solvation of organic molecules [Mihrin et al., Phys. Chem. Chem. Phys. 21(4), 1717 (2019)]. The present work employs this direct far-infrared spectroscopic approach to investigate the non-covalent intermolecular forces involved in the micro-solvation of a selection of seven ether molecules with systematically varied alkyl substituents: dimethyl ether, diethyl ether, diisopropyl ether, ethyl methyl ether, t-butyl methyl ether, and t-butyl ethyl ether. The ranking of the observed out-of-plane water librational band signatures for this selected series of ether–water complexes embedded in inert neon matrices at 4 K reveals information about the interplay of directional intermolecular hydrogen bond motifs and non-directional and long-range dispersion interactions for the micro-solvated structures. These far-infrared observables differentiate minor subtle effects introduced by specific alkyl substituents and serve as rigorous experimental benchmarks for modern quantum chemical methodologies of various levels of scalability, which often fail to accurately predict the structural variations and corresponding vibrational signatures of the closely related systems. The accurate interaction energies of the series of ether–water complexes have been predicted by the domain based local pair natural orbital coupled cluster theory with single-, double-, and perturbative triple excitations, followed by a local energy decomposition analysis of the energy components. In some cases, the secondary dispersion forces are in direct competition with the primary intermolecular hydrogen bonds as witnessed by the specific out-of-plane librational signatures.

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I. INTRODUCTION

The large-amplitude out-of-plane hindered rotational motion of individual water molecules introduced upon complexation has proven enormously important for the unique hydrogen bond network rearrangement dynamics in liquid water2–4 and even in the smallest of the water clusters.5–7 The associated vibrational fundamental transitions occurring in the experimentally challenging far-infrared spectral region is the most prominent spectroscopic manifestation of hydrogen-bonded water molecules. The distinct far-infrared spectroscopic observables associated with this highly localized out-of-plane librational motion of individual water molecules have recently been shown to provide very sensitive spectroscopic probes for the hydrogen bond acceptor capabilities of importance for micro-solvation processes of alcohol molecules8,9 and several other classes of organic molecules.10 The intermolecular out-of-plane librational motion of the first solvating water molecule causes a significant change of dipole moment and gives rise to a strong characteristic far-infrared band correlated quantitatively with the complexation energy. The conventional relationship between the complexation energies and spectral redshifts of the intramolecular OH stretching bands relative to the free water molecule observed in the mid-infrared region is less well defined for hydrated systems owing to the more effective decoupling of the two OH stretching modes of water as the intermolecular hydrogen bonds become stronger.
In the present work, this far-infrared cluster spectroscopy approach supported by theoretical interaction energy decomposition methodologies has been employed to explore the effect of specific alkyl substituent compositions on the intermolecular forces involved in the micro-solvation processes for a variety of aliphatic ethers. The +I inductive effect is one of the fundamental forces in which molecular surroundings influence the hydrogen bond donor–acceptor capability of a molecule, which, in turn, govern macroscopic phenomena as solvation processes.16 In the case of non-linear molecular complexes, the dispersion interactions with bulky alkyl substituents are furthermore expected to compete with the intermolecular hydrogen bonds and affect the overall stabilities of the molecular complexes.13,14 Quantum chemical methodologies, in general, struggle to provide quantitative descriptions of this interplay between different classes of non-covalent interactions, which are all critically important in both the energy and materials and life sciences. Rigorous experimental benchmarking of such intermolecular energy balances is thus crucial both for the validation and future development of modern theoretical methods.

Aliphatic ethers are excellent model systems for the classical intermolecular O–H⋯O hydrogen bonding type. These compounds have the R–O⋯R′ motifs, resulting in higher sensitivity to the inductive effect compared to, e.g., ketones. The selected sample set expands upon previously acquired experimental datasets on various alcohol–water complexes and is still within reach for accurate quantum chemical approaches. As the structurally related ether–water complexes have two alkyl groups substituted to the O atoms of the hydrogen bond acceptor motifs instead of one, the joint investigation of symmetric and asymmetric ether molecules with the same number of electrons, e.g., diethyl ether and methyl propyl ether, may shed further light on the competition between the different intermolecular forces. The experimental datasets on ether–water mixtures are of value themselves, as many of the ethers addressed in the present investigation are widely used in industry15 and are useful in many applications, such as organic synthesis,14 extraction of natural compounds, and oil production.15

II. METHODOLOGIES
A. Experimental details

The diethyl (DEE), diisopropyl (DIPE), tert-butyl methyl (TBME), tert-butyl ethyl (TBBE), and propyl methyl (PME) ethers of HPLC grade (>99.9%, inhibitor-free) were purchased from Sigma-Aldrich. Several milliliters of the liquid were sampled into a glass vial fitted with a polytetrafluoroethylene cold CsI stopcock containing activated molecular sieves 4 Å. The samples were all degassed under vacuum at about −50°C to remove the dissolved air and minor traces of CO₂.

The dimethyl ether (DME) gas (99.9%, Sigma-Aldrich) was transferred from a lecture bottle into a glass volume and further distilled from traces of air by a series of LN₂ freeze–pump–thaw cycles. For sample preparation, DME was nitrogen-cooled and transferred at the sublimation pressure. Ethyl methyl ether (EME) was prepared in house according to the Williamson method16,17 from ethyl bromide and a sodium methoxide solution in methanol, with the product vacuum distilled in a LN₂ trap. The water samples were prepared by placing a few milliliters of Milli-Q grade water in a glass vial fitted with a vacuum valve and degassing under vacuum.

The setup for far-infrared matrix isolation spectroscopy consists of an immersion helium cryostat in conjunction with a Bruker IFS 120 high-resolution FTIR spectrometer. The cryostat contains a rotatable vacuum shroud with positions for a CsI window and a gate valve for the inlet system. The inert gas matrix was obtained by simultaneous deposition of neon gas (99.999%, Air Liquide) with the sample vapor through separate inlet tubes onto a gold plated oxygen-free high thermal conductivity copper mirror at 3 K. The experimental sensitivity was increased by means of a specially designed dual-pass optical arrangement guiding the focused probe beam onto and back from the cold copper mirror of the cryostat and has previously been employed for investigations of weakly bound molecular complexes.18,20 The neon was supplied via a MKS mass flow controller at 8 sccm flow rate, passing a LN₂-cooled spiral, depositing ∼20 mmol of host gas per hour. The sample gas was supplied through a small-flow metering valve from a 5 l glass volume. For all the matrix isolation experiments, the volume was filled with ∼2 mbar of the respective ether sample and ∼0.5 mbar of water vapor unless stated otherwise. The setup provided spectra in the mid-infrared range (650−4000 cm⁻¹, spectral resolution of 1.0 cm⁻¹) employing the combination of a water-cooled Globar light source, Ge on KBr beam splitter, and a semi-conductor HgCdTe detector, as well as the far-infrared range (150−650 cm⁻¹) employing the source together with a 6 μm multilayer Mylar beam splitter and a liquid helium cooled Si-bolometer. A complementary setup was used to acquire mid-infrared spectra down to 400 cm⁻¹ for some of the systems without the consumption of liquid helium. This setup is based around a closed-cycle Advanced Research Systems (ARS) Cryo DE-204 4 K cryostat with a cold CsI window for measurements in the transmission mode. The inlet system designed for this cryostat is analogous to the system described above. The mid-infrared spectra of doped neon matrices obtained with this setup were recorded by a Bruker Vertex 80 V spectrometer, equipped with a broadband MCT detector, Ge on KBr beam splitter, and an air-cooled globalar source.

Following the acquisition of the initial spectra of the neon matrices (denoted pre-annealing spectra), the matrices were annealed by raising the temperature to 8.5 K for 10 min using a resistive heater and a Si-diode temperature sensor attached to the sample holder. As the cryogenic neon host medium is heated gently, the light water molecules diffuse in this softer environment and enable the micro-solvation of ether molecules in neighboring matrix cages. At the end of annealing, the matrices were all allowed to cool before the post-annealing sample spectra were recorded. The background spectra were collected for the evacuated cryostats with the clean sample mirror/window in the same optical arrangements at room temperature.

B. Computational details

The quantum chemical calculations were performed using the ORCA (v.4.1).21 Potential energy minima structures and corresponding harmonic vibrational frequencies of the ether–water molecular complexes and their constituting fragments were calculated using the MP22 method with an aug-cc-pVQZ (AVQZ)3 basis set. The resolution of the identity approximation was used for
the MP2 and SCF integrals together with the numerical “chain-of-spheres” approximation and appropriate fitting basis sets for both the energy and gradients. Partial atomic charges were calculated with the CHELPG method using COSMO radii.

The electronic dissociation energies \( (D_e) \) of the complexes were obtained by applying the domain based local pair natural orbital-coupled cluster theory with single-, double-, and perturbative triple excitations \( [\text{DLPNO-CCSD(T)}] \) for single point calculations, with the aug-cc-pV5Z basis set and the RI-JK approximation, on the respective MP2 geometries. The cutoffs were set using the TightPNO keyword.\(^\text{15,16} \) The zero-point energy corrected dissociation energies \( (D_0) \) of the binary complexes were calculated as the sum of the DLPNO-CCSD(T) electronic dissociation energies and the respective harmonic AZPE values from the MP2 calculations. The interaction energies of the complexes were analyzed with the DLPNO-based local energy decomposition (LED) scheme. This procedure is described in detail in the literature.\(^\text{20,21} \)

### III. EXPERIMENTAL RESULTS

Subsections III A–III G contain detailed analyses of the experimental far-infrared spectra collected for the selection of ether–water systems. The complete list of band assignments for the investigated systems are given in Table I together with the corresponding calculated structural parameters and harmonic vibrational frequencies.

#### A. Diethyl ether (DEE)–water

Figure 1 shows the far-infrared absorption spectra (200–650 cm\(^{-1}\)) collected for cryogenic neon matrices doped with pure water \((\text{H}_2\text{O}:\text{Ne}) = (1:600)\), pure diethyl ether \((\text{DEE}:\text{Ne}) = (1:600)\), and mixtures of water and diethyl ether \((\text{H}_2\text{O}:\text{DEE}:\text{Ne}) = (1:4:600)\) at 4 K. The far-infrared spectrum collected for DEE embedded in neon alone reveals a series of distinct bands at 202, 245, 252, 315, 378, 436, 441, and 502 cm\(^{-1}\). These bands have all previously been assigned to fundamental transitions associated with intramolecular skeleton bending motion \((202\text{ cm}^{-1})\), large-amplitude \( \text{C}–\text{CH}_3 \) torsional motion \((245\text{ and }252\text{ cm}^{-1})\), skeleton deformation motion \((378\text{ and }502\text{ cm}^{-1})\), anti-phase \( \text{CCO} \) bending motion \((436\text{ cm}^{-1})\), and \( \text{COC} \) bending motion \((441\text{ cm}^{-1})\) for two energetically close \( \text{trans}–\text{trans} \) (denoted \( \text{DEE}_g \)) and \( \text{trans}–\text{gauche} \) (denoted \( \text{DEE}_t \)) conformations of DEE. The detailed vibrational assignments, including the torsional–vibrational anharmonic coupling of modes for these DEE conformations, are beyond the scope of the present contribution and we refer to the scientific literature for additional information.\(^\text{27,28} \)

The far-infrared spectrum of pure water embedded in neon alone reveals several bands previously assigned to the water dimer and water trimer. The chosen \( \text{H}_2\text{O}:\text{Ne} \) mixing ratio mostly favors dimer formation and, to a lesser extent, trimer formation. The band observed at 309 cm\(^{-1}\) is associated with the donor in-plane libration mode \( (\nu) \) of the water dimer, whereas the band at 522 cm\(^{-1}\) is associated with the highly localized out-of-plane libration mode of the water dimer.\(^\text{39} \)

The series of bands observed from 450 to 350 cm\(^{-1}\) and the band at 280 cm\(^{-1}\) in the reference spectrum of pure water in neon have all previously been assigned to vibrational fundamentals associated with librational motion of the cyclic water trimer.\(^\text{40,41} \) These bands vanish completely in the spectrum collected when both DEE and water are co-deposited in neon due to the effective quenching of water trimer formation with 4:1 excess of DEE in the matrix relative to water.

### Table I. Summary of the computed relevant structural parameters at the MP2/AOQZ level of theory: the intermolecular hydrogen bond distances \( (r_{O..O}) \) in Å, the \( \angle \text{O}.\text{H}.\text{O} \) angles (in degrees), the CHELPG partial atomic charges \( (\delta) \) on the O atom of the acceptor moiety, and the harmonically predicted \( (\omega_{\text{oop}}) \) and the experimentally assigned \( (\nu_{\text{oop}},\text{exp}) \) out-of-plane libration bands (in cm\(^{-1}\)) for the global potential energy minima of the investigated hydrogen bonded-complexes of water with aliphatic ethers and aliphatic alcohols with the water dimer reference system.

<table>
<thead>
<tr>
<th>Systems</th>
<th>( \delta_{\text{O}} )</th>
<th>( r_{\text{O..O}} )</th>
<th>( \angle \text{O}.\text{H}.\text{O} )</th>
<th>( \omega_{\text{oop}} )</th>
<th>( \nu_{\text{oop},\text{exp}} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \text{H}_2\text{O} )</td>
<td>0.73</td>
<td>1.943</td>
<td>171.7</td>
<td>628.0</td>
<td>522.4</td>
</tr>
<tr>
<td>( \text{DME} )</td>
<td>0.19</td>
<td>1.867</td>
<td>164.8</td>
<td>653.2</td>
<td>546.6</td>
</tr>
<tr>
<td>( \text{EME} )</td>
<td>0.30</td>
<td>1.864</td>
<td>165.4</td>
<td>658.0</td>
<td>549.1</td>
</tr>
<tr>
<td>( \text{PME}_g )</td>
<td>0.24</td>
<td>1.862</td>
<td>169.9</td>
<td>656.4</td>
<td>547.3</td>
</tr>
<tr>
<td>( \text{DEE} )</td>
<td>0.41</td>
<td>1.866</td>
<td>168.6</td>
<td>669.2</td>
<td>560.1</td>
</tr>
<tr>
<td>( \text{DPE} )</td>
<td>0.59</td>
<td>1.858</td>
<td>168.8</td>
<td>667.8</td>
<td>559.6</td>
</tr>
<tr>
<td>( \text{TBME} )</td>
<td>0.45</td>
<td>1.848</td>
<td>167.4</td>
<td>685.4</td>
<td>567.0</td>
</tr>
<tr>
<td>( \text{TBE} )</td>
<td>0.54</td>
<td>1.854</td>
<td>165.2</td>
<td>693.0</td>
<td>574.5</td>
</tr>
<tr>
<td>( \text{MeOH} )</td>
<td>0.62</td>
<td>1.893</td>
<td>164.7</td>
<td>634.4</td>
<td>527.5</td>
</tr>
<tr>
<td>( \text{EOH}_g )</td>
<td>0.68</td>
<td>1.888</td>
<td>161.6</td>
<td>642.6</td>
<td>551.6</td>
</tr>
<tr>
<td>( \text{EOH}_t )</td>
<td>0.68</td>
<td>1.896</td>
<td>165.6</td>
<td>651.5</td>
<td>545.4</td>
</tr>
<tr>
<td>( \text{i-PHOH} )</td>
<td>0.73</td>
<td>1.896</td>
<td>164.8</td>
<td>656.5</td>
<td>554.3</td>
</tr>
<tr>
<td>( \text{t-BuOH} )</td>
<td>0.77</td>
<td>1.885</td>
<td>165.7</td>
<td>662.7</td>
<td>556.4</td>
</tr>
</tbody>
</table>
For the neon matrices doped simultaneously with water and DEE, several new bands appear in the far-infrared spectrum in close vicinity to the donor OH librational bands of the water dimer. Two overlapped but still distinct bands are observed at 559.7 and 542.7 cm\(^{-1}\) and thereby with significantly higher band origins than the large-amplitude intermolecular out-of-plane librational band of the water dimer (522 cm\(^{-1}\)). At the same time, a new broad band with a less apparent subband structure is observed around 310 cm\(^{-1}\) and thereby strongly overlapped both with the transitions assigned to the large-amplitude intermolecular in-phase \(c\)-type libration of the water dimer (309 cm\(^{-1}\)) and the intramolecular CCO bending mode of DEE at 315 cm\(^{-1}\). In addition, a fourth new band is observed at 259 cm\(^{-1}\) and slightly blue-shifted relative to the free C–CH\(_3\) torsional band of DEE\(_{\text{tr}}\) assigned at 252 cm\(^{-1}\). The new bands all gain significant intensity upon annealing up to 9 K and are for that reason assigned to mixed molecular complexes of DEE and water. The annealing procedure promotes the diffusion of the light water molecules in the soft neon environment above 7 K, favoring the formation of complexes.\(^{39}\) The two distinct bands at 559.7 and 542.7 cm\(^{-1}\) are straightforwardly assigned to transitions associated with the highly localized out-of-plane librational modes of the 1:1 complex of DEE and water. The improved hydrogen bond acceptor capability of DEE strengthens the intermolecular hydrogen bond between DEE and water relative to the hydrogen bond in the water dimer. The out-of-plane librational band origins are expected to follow the trend reported in our previous far-infrared investigations concerned with hydrogen-bonded alcohol–water complexes.\(^{38}\) The observation of two distinct out-of-plane librational bands clearly points at two different conformations of the DEE-H\(_2\)O complex with slightly different hydrogen bond orientations. This is also supported by the annealing procedure indicating that the less stable conformation of the complex with band origin at 542.7 cm\(^{-1}\) is formed slower than the more stable conformation with the band origin at 559.7 cm\(^{-1}\), compared to the pre-annealing ratio between the two forms. The new broad band around 310 cm\(^{-1}\) and the band at 324 cm\(^{-1}\) must then be assigned to the in-phase \(c\)-type librational transitions expected for the two conformations of the DEE-H\(_2\)O complex. The final new band observed at 259 cm\(^{-1}\) may then be assigned to transitions associated with the slightly perturbed C–CH\(_3\) torsional modes of the DEE-H\(_2\)O complex picking up infrared intensity from the dipole moment change of the hydrogen-bonded water molecule.

The annealing procedure further distinguishes features associated with the \textit{trans–gauche} isomer of DEE. The positions of the bands (502 and 376 cm\(^{-1}\)) are in agreement with the literature, but the relative abundance of the conformers is significantly shifted compared to estimates at ambient temperature,\(^{38,40,41}\) and annealing of the matrix further induces relaxation of the remaining \textit{trans–gauche} DEE molecules into the \textit{trans–trans} conformation. This experimental evidence indicates that the two observed conformations of the DEE complex with water differ in the way of how the water molecule is attached to DEE in the \textit{trans–trans} conformation, rather than due to different conformations of the DEE fragment itself. The potential energy minima structures of the different conformations will be investigated computationally in Sec. IV A.

**B. Diisopropyl ether (DIPE)–water**

Figure 2 shows the far-infrared absorption spectra (200–650 cm\(^{-1}\)) collected for cryogenic neon matrices at 4 K doped with pure water (H\(_2\)O:Ne) = (1:600), pure diisopropyl ether (DIPE:Ne) = (1:600), and a mixture of water and diisopropyl ether with a mixing ratio of (H\(_2\)O:DIPE:Ne) = (1:4:600).

The band associated with the low-frequency COC bending mode is observed at 202 cm\(^{-1}\), and the two bands associated with the in-phase and out-of-phase CCC bending modes are located at 301 and 390 cm\(^{-1}\) for the DIPE monomer, respectively. The bands at 407 and 428 cm\(^{-1}\) correspondingly are assigned to the COC twisting and wagging modes of DIPE, the weak 488 cm\(^{-1}\) band is assigned to the COC bending mode, and the band at 510 cm\(^{-1}\) is assigned to the CCCCO umbrella mode of DIPE. The band at 515 cm\(^{-1}\) should likely be assigned to the umbrella mode of a less stable DIPE conformation. These observations in the solid neon environment are in qualitative accordance with previous investigations of both thin films,\(^{31}\) the liquid phase\(^{41}\) and the gas phase.\(^{45}\)

The relevant \(\nu_{\text{oop}}\) band for the respective DIPE-H\(_2\)O complex is clearly observed at 559.6 cm\(^{-1}\), nearly at the same position as has been found for the similar DEE-H\(_2\)O complex; however, only one single slightly asymmetric broad band is observed for the DIPE-H\(_2\)O system. At 320.7 cm\(^{-1}\), an intense band with a shoulder around 313 cm\(^{-1}\) is observed. These bands are assigned to the in-plane librational mode \(\nu_p\) of two different conformations of the DIPE-H\(_2\)O complex. Evidently, the structural differences between the two conformations have a stronger influence on this \(c\)-type water librational mode (the hindered overall rotation of the water molecule) than the highly localized \(\nu_{\text{oop}}\) mode where only the donor H atom is engaged. Upon annealing, the spectral subtraction between the post- and pre-annealing spectra shows a substantial growth of intensity of the 320.7 cm\(^{-1}\) absorption, supporting the fact that this transition should be assigned to the in-plane librational mode of the most stable DIPE-H\(_2\)O conformation.
C. Ethyl methyl ether (EME)-water

Figure 3 shows the far-infrared spectra for the neon matrices doped with pure water (H₂O:Ne = 1:600) and a mixture of ethyl methyl ether with water having a mixing ratio of (EME:H₂O:Ne = 3:1:600). The bands observed in the spectra at 213 and 256 cm⁻¹ should be assigned to the two O–CH₃ torsional modes, the bands at 296 and 471 cm⁻¹ should be assigned to the in-plane and out-of-plane hybrid CCO and COC bending modes of EME, respectively.¹⁶,⁴⁶ The band at 511 cm⁻¹ has not previously been observed in the gas phase. Kitagawa et al.⁴⁶ estimated the gauche conformation of EME to have the bending band at 510 cm⁻¹. This band exhibits a reversed annealing effect (less band intensity in the post-annealing spectrum) in agreement with other monomer bands in the spectrum, suggesting that it might potentially be due to a combination of the 213 and 296 cm⁻¹ bands.

The remaining observed bands are assigned to water clusters and the mixed EME·H₂O complex. The broader structured band with an origin at ~313 cm⁻¹ is assigned to the in-plane librational mode νip, and finally the band with maximum at 549.1 cm⁻¹ is assigned to the target out-of-plane librational mode νoop of the complex. As expected from the previously established trend among homologous ether–water complexes, the band position is lower than found for the DEE·H₂O complex, indicating a lower hydrogen bond acceptor capability of the EME molecule relative to DEE. In contrast to the DEE·H₂O system, only a single conformation of this complex was detected for this system.

D. Propyl methyl ether (PME)-water

Figure 4 shows the infrared spectra for neon matrices doped with pure water (H₂O:Ne = 1:600), a mixture of propyl methyl ether with water having a mixing ratio of (PME:H₂O:Ne = 4:1:600), and a spectrum of pure PME (PME:Ne = 1:200). The observed bands with maxima at 509 and 516 cm⁻¹ are associated with the COC bending/CH₃ rocking motion of the trans–gauche conformation. The bands at 434 and 412 cm⁻¹ band are assigned to the COC and CCC bending modes of the trans–trans conformation.⁴⁶ The bands at 356 and 301 cm⁻¹ are associated with twisting modes of the trans–gauche conformation. The band at 233 cm⁻¹ is associated with CH₃–O torsions of both conformers, and the weak absorption features at 208 and 197 cm⁻¹ are associated with CH₁–C torsional (trans–gauche) and skeletal deformation motions (trans–trans), respectively.

The out-of-plane librational band νoop of the most stable conformation of the PME·H₂O complex is assigned at 547.3 cm⁻¹. This very distinct band has a reproducible shoulder around 537.2 cm⁻¹ in the pre-annealing spectrum, which is significantly reduced in intensity in the post-annealing spectrum. The difference spectrum clearly reveals the relaxation of the less stable conformation (intensity decrease of the 537.2 cm⁻¹ band) into the global potential energy minimum (intensity increase of the 547.3 cm⁻¹ band). This trend is supported by the observations in the region of the in-plane librational motion νip. A distinct band at 306 cm⁻¹ shows a strong positive annealing effect, matching this band with the 547.3 cm⁻¹ band of the global minimum conformation, whereas a band around 287 cm⁻¹ decreases in intensity upon annealing, indicative of the local potential energy minima. As all the detected out-of-plane librational bands νoop are located above 500 cm⁻¹, the remaining of the selected ether–water systems are solely presented in the spectral region down to 400 cm⁻¹.
E. Tert-butyl methyl ether (TBME)–water

Figure 5 shows the infrared spectra (400–800 cm\(^{-1}\)) for the neon matrices doped with a mixture of tert-butyl methyl ether with water (TBME\(\cdot\)H\(_2\)O:Ne = 4:1:600), pure water (H\(_2\)O:Ne = 1:600), and pure TBME (H\(_2\)O:Ne = 1:150). There is scarcity of information about the vibrational spectrum of TBME below 700 cm\(^{-1}\); however, our theoretical harmonic predictions as described in Sec. IV A are in excellent agreement with our matrix experiment in the 300–700 cm\(^{-1}\) spectral range. The TBME monomer bands located at 414, 459, and 508 cm\(^{-1}\) should be assigned to the tert-butyl substituent’s umbrella mode, COC skeletal deformation, and the COC bending motion, respectively. The band at 726 cm\(^{-1}\) should be assigned to the tert-butyl breathing motion.\(^{49}\) The far-infrared bands at 371 and 346 cm\(^{-1}\) are assigned to the COC/CCC bending motion and the CCCC asymmetric bending motion of the tert-butyl group, respectively. The two broad bands related to mixed complexes are present in the range of the spectrum where the out-of-plane librational band \(\nu_{\text{oop}}\) is expected for the respective binary system. These two bands have been reproduced in experiments with several mixing ratios (8:1:1200 to 1:1:300), but the bands exhibit different relative annealing responses. These intensity changes with annealing can be correlated with the two respective bands in the OH-stretching range of the spectrum (not shown), located at 3485 and 3505 cm\(^{-1}\). Additional experimental work would be required for an unambiguous confirmation of the proposed assignment, but based on the similar findings for the TBME\(\cdot\)H\(_2\)O complex, we believe that this band should likewise be attributed to a ternary TBEE\(\cdot\)H\(_2\)O complex. The assumption that a water-rich ternary TBEE\(\cdot\)H\(_2\)O system is formed in the matrix above an annealing threshold, is somewhat reinforced by the observation that the number density of water dimers seems to go down simultaneously. The band origin for the out-of-plane librational band \(\nu_{\text{oop}}\) of the binary TBME\(\cdot\)H\(_2\)O complex is estimated to be 574.5 cm\(^{-1}\).

C. Dimethyl ether (DME)–water revisited

In our recent work,\(^{8}\) the out-of-plane librational band of the DME\(\cdot\)H\(_2\)O complex in the neon matrix environment was reassigned at 555.9 cm\(^{-1}\). This system is revisited again in light of the present extensive complementary experimental results for the series of homologous ether–water complexes. The DME\(\cdot\)H\(_2\)O system was initially investigated in several noble gas matrices, and...
the far-infrared assignments of the out-of-plane librational mode of this complex were proposed at 524.0, 555.9, and 547.4 cm\(^{-1}\) in doped neon, argon, and krypton matrices, respectively.\(^{31}\) The substantial shift from argon to neon matrices reported is not consistent with typically associated matrix spectral shifts and these observations do not follow the host material trend reported for the corresponding out-of-plane librational mode of the water dimer. In addition, this 524.0 cm\(^{-1}\) band assignment to the out-of-plane librational mode of the most stable conformation of the DME-H\(_2\)O complex seems too close to the assignment for the water dimer at 522.4 cm\(^{-1}\) considering the substantial difference in interaction energies between the two complexes.\(^{15}\) Our recent experimental findings showed that, apart from the very broad 524 cm\(^{-1}\) band, a comparatively weak absorption could be found at 555.9 cm\(^{-1}\). However, the present experimental assignments of the out-of-plane librational bands for the homologous ether–water complexes, namely, DME-H\(_2\)O at 549.1 cm\(^{-1}\) and PME-H\(_2\)O at 547.3 cm\(^{-1}\), indicate that this very weak 555.9 cm\(^{-1}\) absorption also should not be assigned to the out-of-plane librational transition of DME-H\(_2\)O.

In our revisit of this system, we have performed a new series of experiments with a wide range of DME-H\(_2\)O:Ne mixing ratios ranging from 1:1:600 to 24:1:1200. The far-infrared spectra of these independent matrices are shown in Fig. 7 both before and after annealing together with the reference spectra of DME and water embedded in neon alone. The rather weak 555.9 cm\(^{-1}\) absorption is again reproduced in the spectra collected for both mixtures, as well as the broad band at 524 cm\(^{-1}\). The annealing of both matrices, however, independently reveal the substantial growth of a third quite distinct band at 546.6 cm\(^{-1}\), while the aforementioned absorption features show little to no growth. Remarkably, this new 546.6 cm\(^{-1}\) band is only seen as a weak shoulder in the pre-annealing spectra and only slightly more evident when DME is abundant in large excess. At the same time, we observe in the range of the in-plane librational mode that, previously assigned for this complex at 293 cm\(^{-1}\), a strong and broad asymmetric band indeed appears at 293 cm\(^{-1}\). On annealing of the matrix, this band, however, shows only minor growth. At the same time, a new distinct band emerges at 314 cm\(^{-1}\) as a shoulder in spectra with higher DME-H\(_2\)O mixing ratios and stands out clearly in the experiment with huge DME excess. This new band shows a pronounced annealing effect very similar to the behavior of the 546.6 cm\(^{-1}\) band. In addition, the 314 cm\(^{-1}\) band position agrees better with theoretical predictions for the in-plane librational transition of the global minimum for the DME-H\(_2\)O complex as will be evident in Sec. IV A. Based on these new strong experimental annealing evidences, we then believe that all the reported observations can be explained by the presence of two different energetically close conformations of the DME-H\(_2\)O complex.

In the investigation on the similar DME-methanol complex by Han and Kim,\(^{32}\) the authors reported experimental evidences suggesting the presence of two different conformations. The difference between the two ways the hydrogen bond donor molecule can attach DME can be viewed as the relative orientation of the unbound side of the donor fragment. The two calculated potential energy minima structures in Fig. 8 represent our best attempts to capture these two conformations for the DME-H\(_2\)O complex. While both the conformational wells in the DME-methanol complex probably are well defined, owing to the additional dispersion interactions from the methyl group, the energy barrier between the two similar conformations of the DME-H\(_2\)O complex is expected to be smaller. The implication is that the local minimum cannot be reliably captured computationally at higher levels of theory. Nevertheless, at the DFT-level, we have still predicted the difference in the out-of-plane librational band origin for the two conformations to be about 20 cm\(^{-1}\) in the harmonic approximation close to the behavior of the 546.6 cm\(^{-1}\) band. In addition, the 314 cm\(^{-1}\) band position agrees better with theoretical predictions for the in-plane librational transition of the global minimum for the DME-H\(_2\)O complex as will be evident in Sec. IV A. Based on these new strong experimental annealing evidences, we then believe that all the reported observations can be explained by the presence of two different energetically close conformations of the DME-H\(_2\)O complex.
to the experimental findings. Based on these considerations, we assign the newly observed bands at 546.6 and 314.0 cm$^{-1}$ as the out-of-plane and in-plane librational modes, respectively, for the global minimum conformation (conformation A), while the previous assignments by Engdah and Nelander could be attributed to the less stable conformation of the complex (conformation B).

IV. SUPPORTING THEORETICAL RESULTS AND DISCUSSIONS
A. Intermolecular potential energy minima and structural parameters

To aid the interpretation of the present experimental datasets, the present results have been supported by $ab$ initio calculations of the potential energy minima structures and the corresponding harmonic vibrational spectra for all systems at the MP2/AVQZ level of theory. For some of the systems, only one conformation was found at this computational level and all of these systems were found to be clearly hydrogen-bonded complexes with water serving as the hydrogen bond donor. The relevant structural parameters, such as intermolecular hydrogen bond lengths and O–H⋯O angles, were calculated as well as the partial charges located on the O atoms of the hydrogen bond accepting ether molecules. The computational results are presented in Table I together with the corresponding values for the previously investigated hydrogen-bonded alcohol–water complexes.

Multiple potential energy minimum conformations were found for some of the ether–water systems. For the mono-hydrated complexes of the linear DEE, EME, and DME molecules, two conformers associated with a different rotation of the water molecule along the hydrogen bond were optimized (see Fig. 9). The conformations of type B are less stable and have lower calculated out-of-plane librational band positions than the respective type A conformations. For the DEE$\cdot$H$_2$O system, the harmonically predicted out-of-plane librational band origin of the conformation B is 18 cm$^{-1}$ lower than for conformation A, explaining the two observed bands in the corresponding matrix experiment. This isomerism is less likely for the asymmetric linear ether molecules as the interaction with the corresponding matrix experiment. This isomerism is less likely for conformation A, explaining the two observed bands in the spectral range.

The predicted structural parameters suggest that, in those cases where the bond angle deviates from the ideal linear 180° configuration, strong secondary interactions presumably of dispersive nature are acting in competition with the primary hydrogen bond. The hydrogen bond angles of the DME$\cdot$H$_2$O and EME$\cdot$H$_2$O complexes are noticeably lower than calculated for the water dimer. The long-chained ethers, such as PME, pose an additional challenge for an accurate comparison to theory as the fully stretched conformation of the skeleton is not necessarily the most abundant as discussed above. Another observable effect is the twist of the water molecule in the PME$\cdot$H$_2$O complex (Fig. 9). In contrast to the complex of a symmetric ether, e.g., the DME$\cdot$H$_2$O system, the water molecule is turned with the oxygen closer to the larger substituent due to dispersive interaction of the water molecule with the alkyl substituent in competition with the hydrogen bond. The twist is also present in the calculated structure of the asymmetric EME$\cdot$H$_2$O system but to a much smaller extent.

The inductive effect of the substituents falls off rapidly with bond count and the difference is almost negligible between the ethyl and propyl substituents. The partial charge alone does not appear to provide sufficient information for the explanation of the experimental observations. It is clear that the high negative charge and the large bond angle found in the water dimer are indicative of the dominant electrostatic component of this interaction. The theoretical vibrational frequencies generally correlate quite well with the experimental results with a scaling factor to the harmonic predictions of ~84%. However, the differences between closely positioned band origins of the DME$\cdot$H$_2$O, EME$\cdot$H$_2$O, and PME$\cdot$H$_2$O complexes are overestimated, suggesting that an accurate treatment of electronic correlation plays an important role in a quantitative differentiation of the acceptor capability of this series of molecules.

B. Local energy decomposition of complexation energies

To investigate the interplay between the different classes of intermolecular forces more deeply, the interaction energies of the
systems were analyzed employing the local energy decomposition (LED) method, exploiting the local pair nature of the underlying DLPNO-CCSD model. The method has previously been tested on a set of diverse systems and has been applied for the analysis of the intermolecular forces of the water dimer.\textsuperscript{35}

The net dissociation energy of the complex (\(D_0\)) is decomposed into the geometry preparation energy \(\Delta E_{\text{geo}}\), the dispersive [pure dispersion \(\Delta E_{\text{disp}}^{\text{Disp}}\), \(\Delta E_{\text{disp}}^{\text{Disp}}\) combined with the (T) component\textsuperscript{35,36}] and the non-dispersive (electrostatic preparation \(\Delta E_{\text{prep}}\), electrostatic and exchange \(\Delta E_{\text{el-ex}}\)) terms, and non-dispersive contributions from the CCSD pairs \(\Delta E_{\text{res}}^{\text{res}}\) contributions. A detailed description can be found in the supplementary material and in Refs. 35 and 57. The final fragmentation of the complete zero-point corrected dissociation energy (\(D_0\)) for a given ether–water complex is

\[
D_0 = \Delta E_{\text{geo}} + \Delta E_{\text{el-prep}}^{\text{HF}} + \Delta E_{\text{el-ex}}^{\text{HF}} + \Delta E_{\text{res}}^{\text{res}} + \Delta E_{\text{disp}}^{\text{Disp}} + \Delta E_{\text{disp}}^{\text{Disp}} + \Delta E_{\text{disp}}^{\text{Disp}}(T) + \Delta ZPE. \tag{1}
\]

The results of the analysis are summarized in Table II. As can be expected, the water dimer complexation energy is heavily dominated by the non-dispersive component. The interaction energies of the mixed ether-complexes all show dramatically higher relative contributions from the dispersion component, increasing from ~30% for the water dimer to ~40% for the DME-H\(_2\)O complex and growing further as expected with the alkyl substituent sizes.

In the DME-H\(_2\)O–EME-H\(_2\)O–PME-H\(_2\)O sequence, the total dissociation energy increases, converging with the alkyl substituent length. The non-dispersive component, however, can be seen to decrease, while the dispersion contribution increases. This suggests that the dispersion interaction is in competition with the electrostatic component. From the DME-H\(_2\)O to EME-H\(_2\)O systems, the drop of the non-dispersive energy is small compared to the increase of the dispersion, whereas going to the PME-H\(_2\)O system, the non-dispersive component decreases significantly. This is likely due to the fact that in the PME-H\(_2\)O complex, the dispersive forces act from a greater angle, straining the hydrogen bond. On the other hand, in the EME-H\(_2\)O complex and specifically in the DME-H\(_2\)O complex, the mean vector of dispersive force is more aligned with the hydrogen, so the reduction in the net non-dispersive component comes primarily from the increase in electrostatic repulsion. The latter is calculated to be identical for the EME-H\(_2\)O and PME-H\(_2\)O complexes but drastically lower for the DME-H\(_2\)O system. As the out-of-plane librational band position does not follow the dissociation energy for these complexes, the aforementioned effect can be seen as the one of the causes for this deviation.

In line with the conclusion, the symmetric DEE-H\(_2\)O complex shows an increase of both the net N/D interaction and the electrostatic preparation component over the PME-H\(_2\)O system. The librational band position for this complex is higher by 11 cm\(^{-1}\) than found for EME-H\(_2\)O, despite the fact that the dispersion ratio is increased to 50% over the 45% in the latter.

Both conformations of the DIPE-H\(_2\)O complex show an increase of the N/D and dispersive energy components, owing to the presence of two large branched substituents. However, the experimentally observed librational band position is nearly identical to that of the DEE-H\(_2\)O complex, also in agreement with the calculations. It is very likely that the specific surroundings of the water molecule is affecting the band position, as the \(\Delta ZPE\) for this complex is measurably higher than predicted for the DEE-H\(_2\)O.

The TBME-H\(_2\)O complex shows the highest net non-dispersive interaction component of all the investigated ether–water complexes, almost on par with the value for the water dimer. However, a

### Table II

The results of the computational analysis of the dissociation energies for the binary complexes of water with ethers, aliphatic alcohols, and the water dimer using the LED method at the DLPNO-CCSD(T)/AVZ level (units of \(kJ\) mol\(^{-1}\)). This table includes the geometric preparation energy \(\Delta E_{\text{geo}}\), the electronic preparation (EP), the HF electrostatic and exchange terms (e\(st+X\)), the sum of the HF interaction terms (Int\(_{\text{HF}}\)), the CCSD strong and weak pairs (SP\(_{\text{res}}\) and WP), the pure dispersive contribution (Disp.) and the triples term (T). The total dispersive (\(\text{Total}_{\text{Disp}}\)) and the non-dispersive (\(\text{Total}_{\text{ND}}\)) terms are indicated in bold. The total change of zero-point energies \(\Delta ZPE\) upon complexation are calculated at the MP2/AVQZ level in the harmonic approximation.

<table>
<thead>
<tr>
<th>System</th>
<th>(\Delta E_{\text{geo}})</th>
<th>EP</th>
<th>e(st+X)</th>
<th>Int(_{\text{HF}})</th>
<th>SP(_{\text{res}})</th>
<th>WP</th>
<th>Total(_{\text{ND}})</th>
<th>Disp</th>
<th>Total(_{\text{Disp}})</th>
<th>(D_0)</th>
<th>(\Delta ZPE)</th>
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<tr>
<td>H(_2)O</td>
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<td>-15.2</td>
<td>0.3</td>
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<td>-14.4</td>
<td>-0.3</td>
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<td>-1.5</td>
<td>-10.4</td>
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<tr>
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<td>142.7</td>
<td>-156.5</td>
<td>-13.8</td>
<td>-0.7</td>
<td>-0.09</td>
<td>-14.6</td>
<td>-10.3</td>
<td>-1.8</td>
<td>-12.1</td>
<td>-25.9</td>
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<tr>
<td>PME(_x)</td>
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<td>-155.7</td>
<td>-14.0</td>
<td>-0.8</td>
<td>-0.10</td>
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<td>-11.4</td>
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<td>-0.10</td>
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<td>-15.6</td>
<td>-29.4</td>
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<tr>
<td>TBME</td>
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<td>162.5</td>
<td>-176.5</td>
<td>-14.0</td>
<td>-0.7</td>
<td>-0.12</td>
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<td>EtOH(_x)</td>
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<td>0.7</td>
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<td>-1.6</td>
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<td>EtOH(_t)</td>
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<td>-144.9</td>
<td>-15.0</td>
<td>0.6</td>
<td>-0.10</td>
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<td>-1.6</td>
<td>-11.3</td>
<td>-25.1</td>
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<td>-161.6</td>
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<td>-0.13</td>
<td>-14.6</td>
<td>-12.0</td>
<td>-1.9</td>
<td>-13.8</td>
<td>-27.7</td>
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</tbody>
</table>
large portion comes from the CCSD charge-polarization component (SP$_{ctt}$). The dispersion component of the energy is somewhat below that of the DIPE system, as well as the total $D_0$. With the increase of the alkyl chain length in the TBEE-H$_2$O complex, most of the changes resemble the cases of the linear ethers. There is a reduction in the net N/D component, associated primarily with the rapid growth of the electrostatic repulsion outweighed by the increase in dispersive forces, while all other sub-components increase in magnitude as well. These differences manifest in a measurably higher librational band position for the TBEE-H$_2$O complex.

The set of previously studied alcohol–water complexes falls well into the newly acquired results, with a decent correlation between the theoretical dissociation energies of the complexes and their respective observed out-of-plane librational band positions. Each of the alcohol molecules forms a somewhat weaker complex with water than their respective ether analog, e.g., $r$-BuOH-H$_2$O vs TBME-H$_2$O and EtOH$_2$-H$_2$O vs EME-H$_2$O. The primary sources of these differences, according to the LED results, stem from the higher dispersion interaction with the introduced methyl group and from a large gain in the charge-transfer component. The present quantum chemical calculations correlate very well with the experimental observations for the out-of-plane librational bands of this series of water complexes, unlike in the comparison of the EME–PME pair. The calculated O–H stretching band origins show that the PME-H$_2$O complex has a larger redshift relative to the EME-H$_2$O complex. This suggests that the out-of-plane librational band position is more sensitive to the electrostatics and the N/D components of the electronic correlation and, therefore, appears to provide a somewhat better probe of the localized intermolecular hydrogen bond itself, rather than the cumulative interaction energy, at least when fine differences in the structure and dispersion energy are present. This conclusion seems to be supported by the out-of-plane librational signatures for the TBME-H$_2$O and DIPE-H$_2$O complexes where the observed band origin for the latter system is significantly smaller despite the overall higher dissociation energy. On the contrary, the local energy decomposition approach shows that the TBME-H$_2$O interaction contains the larger non-dispersive energy contribution. For cases where the bonding is highly similar and the energy differences are high, e.g., in the series of symmetrical ethers, a high degree of linearity with respect to the theoretical $D_0$-values is observed. This is likely one of the sources of the remarkably linear relationship between observed out-of-plane HF librational band origins and the square root of the respective experimental binding energies ($D_0$ vs $\nu_{oo}$) reported previously for a series of linear hydrogen-bonded HF complexes by Klemperer et al.\textsuperscript{88} The absolute band centers for the out-of-plane HF librational transitions were shown to be more reliable measures of the intermolecular hydrogen bond energies $D_0$ than the traditional complexation redshifts of the intramolecular HF stretching bands relative to the free HF molecule.

V. CONCLUSIONS

The interplay of non-covalent intermolecular forces involved in the micro-solvation of seven different aliphatic ether molecules with systematically varied alkyl substituents has been explored by far-infrared neon matrix isolation spectroscopy in combination with explorative quantum chemical modeling. The primary focus experimentally has been to recover the very informative band origins for the highly anharmonic intermolecular out-of-plane librational transitions of this series of hydrogen-bonded ether–water complexes to complement previous far-infrared investigations of more simple alcohol–water complexes. The out-of-plane librational motion of the first solvating water molecule involves a significant change of the water dipole moment and gives rise to strong characteristic far-infrared bands in the range from 546 to 574 cm$^{-1}$, which are correlated quantitatively with the intermolecular ether–water complexation energies with our currently best feasible theoretical estimates, excluding anharmonic corrections, ranging from 17.0 (DME-H$_2$O) to 21.3 kJ mol$^{-1}$ (DIPE-H$_2$O). For the energetically close complexes, a fine interplay between secondary dispersive interactions and the primary intermolecular hydrogen bonds has been revealed, which is poorly predicted by the commonly used scalable electronic structure methodologies. For the larger asymmetric systems, the directional intermolecular hydrogen bonds are in direct competition with non-directional dispersion forces comprising almost half of the total interaction energy $D_0$. The local energy decomposition of the interaction energies shows that the net non-dispersive component decreases with the size of the substituents caused by the increased electrostatic repulsion of the fragments. The present contribution suggests that the absolute out-of-plane librational band position is more closely linked to the electrostatics and the N/D components of the electronic correlation. This observable appears to provide a somewhat better probe of the local and directional intermolecular hydrogen bonds rather than the less directional cumulative intermolecular binding motifs. The present dataset provides rigorous experimental benchmark points for further improvement of high-level ab initio electronic structure methods in the field of non-covalent interactions. The conventional RI-MP2 approach with the AVQZ basis set shows good consistency for predictions of the relative alkyl substituent effect when calibrated with experimental data for closely related systems, which allows us to remove the uncertainty associated with anharmonicity. The present work has been limited to linear and branched alkyl substituents that are all resulting in an $+l$ inductive effect and thereby increased hydrogen bond acceptor capability of the O atoms compared to the smallest DME-H$_2$O system and the water dimer. The next step in further investigations of the aforementioned correlations would be to study analogous ether–water systems containing more complex substituents, e.g., alkenyl and aryl fragments, tuning the acceptor fragments and shifting the balance of electrostatic and dispersive forces, as well as expanding the range of the set toward lower interaction energies.

SUPPLEMENTARY MATERIAL

The supplementary material contains a description of the local energy decomposition (LED) procedure, the XYZ coordinates of the optimized structures, the graphic representations of the conformations of the DIPE and DEE complexes with water, and the relevant ORCA output files containing the harmonic vibrational frequencies.

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AUTHOR DECLARATIONS

Conflict of Interest

The authors have no conflicts to disclose.

DATA AVAILABILITY

The data that support the findings of this study are available from the corresponding author upon reasonable request.

REFERENCES


