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Polarization manipulation and management are important for 4H-silicon carbide (SiC) integrated photonics, as 4H-SiC has material-based birefringent properties. In this paper, we propose a low-birefringence polarization beam splitter (PBS) based on asymmetric directional coupler (ADC) mode converters with overall high performances. We numerically and experimentally demonstrate the ADC mode conversion based PBS on a 4H-SiC chip. The experimental results show that the device exhibits high transmittance of -0.6 dB and -1.3 dB for the transverse-electric (TE) and transversemagnetic (TM) polarized light, respectively and broad operational bandwidth over 130 nm. The polarization extinction ratio of >25 dB and >17 dB covering the whole C band for the TE and TM polarized light, respectively, and an ultra-large polarization extinction ratio of >32 dB for both polarizations at around 1560 nm are achieved. © 2022 Optica **Publishing Group**

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Silicon carbide (SiC), exhibiting superior physical properties, is 5 one of the most widely used third-generation semiconductors [1]. Its wide bandgap, as well as high saturation velocity and large breakdown 6 voltage, makes SiC an attractive candidate for power electronics^[2]. The wide bandgap property of SiC also enables a wide window of 8 optical transparency in a wavelength range from near ultraviolet to 9 mid-infrared and an eradication of two-photon absorption at telecom 10 wavelength, benefiting optical applications [3]. Meanwhile, thanks to 11 its strong second- and third-order optical nonlinearities and material-12 based color centers, SiC also attracts a lot of attentions in nonlinear 13 and quantum optics [4-6]. Therefore, SiC is regarded as a growing 14 alternative to silicon based electronic and photonic components. In re-15 cent years, the ion-cut technique and the grinding-chemical mechanical 16 polishing method facilitate high-quality 4H-SiC-on-insulator (SiCOI) 17 stacks, and the advanced nanofabrication technologies promote the 18 development of 4H-SiC integrated photonics [7–13]. 19

Polarization manipulation is very important in both photonic inte grated circuits and quantum photonic integrated circuits, which can be
 used for on-chip polarization-division multiplexing and polarization
 entanglement-enabled quantum key distribution [14–17]. Especially,
 for the 4H-SiCOI integrated platform, the polarization control is an

essential issue. Because 4H-SiC is a uniaxial crystal, optical properties, including refractive index and nonlinear refractive index, are crystal-orientation dependent. As a result, light with different polarizations, propagating in 4H-SiC waveguides, also show different properties [18, 19]. The polarization beam splitter (PBS) is one of the important devices for polarization manipulation, which is able to split or combine light beams with different polarizations. The performances of PBSs are evaluated by transmission efficiency, operational bandwidth, and polarization extinction ratio [20]. Although plenty of high-performance PBSs have been demonstrated in the silicon photonic integrated platform [21], it is still challenging to achieve in the SiC photonic integrated platform. Compared to silicon, SiC has relatively lower refractive index, which results in low birefringence of the orthogonal polarizations in the waveguides with large dimensions [7]. As a result, it is not efficient to make low-birefringence beam splitting with conventional directional couplers. A novel vertical-dual-slot structure has been proposed to achieve low-birefringence beam splitting in the SiC integrated platform with a small footprint, but the polarization extinction ratio is not satisfactory [7]. Therefore, it remains a demanding challenge to achieve PBSs with overall good performances in the SiC integrated platform.

Directional couplers and asymmetric directional couplers (ADC) have been widely used for efficient power coupling [22], mode conversion [23], as well as polarization beam splitting especially in silicon integrated platforms. The scheme of the directional coupler, assisted by a bridged waveguide, has been demonstrated to achieve high polarization extinction ratio for polarization beam splitting [24]. By using thin silicon layer to have large birefringence for the two polarizations, a higher polarization extinction ratio can be achieved [25]. An alternative way to increase the polarization extinction ratio is to apply ADC, for the purpose of enlarging the phase mismatching of the mode coupling for the two polarizations [26], and it is also valid for thick silicon layer [27]. Similar designs have also been reported in low-refractive-index materials based integrated platforms, such as silicon nitride and lithium niobate, showing great advantages for low-birefringence beam splitting [28, 29].

In this letter, we apply a pair of mode converters, based on ADC, for the low-birefringence polarization beam splitting in the SiC integrated platform. One mode converter is used to convert the fundamental TM (TM_0) mode to the first-order TM (TM_1) mode, according to the phase matching between the two modes. At the same time, it can split the orthogonal polarizations, because the phase-matching

condition is not valid for the TE modes. A cascaded identical mode 67 converter is directly connected to the first mode converter, in order 68 to convert the TM_1 mode back to the TM_0 mode. More importantly, 69 it can further eliminate the leakage of the TE mode from the first 70 mode converter, so that the polarization extinction ratio can be largely 71 increased. The device is numerically and experimentally demonstrated 72 with high transmission efficiency, broad operational bandwidth, and 73 large polarization extinction ratio in the SiC integrated platform, within 74 a small footprint. 75



Fig. 1. Schematic of the ADC mode conversion based PBS.

Figure 1 shows the schematic of the proposed ADC mode conver-76 sion based PBS. The inset shows the cross section of an ADC mode 77 converter. The device sits on a 4H-SiCOI integrated platform, with a 78 4H-SiC thin film thickness of 400 nm. The extraordinary optical axis of 79 4H-SiC is perpendicular to the wafer plane. The ADC mode converter 80 consists of a single-mode waveguide and a multi-mode waveguide, 81 with widths of W_1 and W_2 , respectively. The coupling length and the 82 gap between the two waveguides are L and G, respectively. The two 83 mode converters in the whole device have the same dimension. When 84 light is injected from Input, after transmitting a pair of ADC mode 85 converters, the TM-polarized light is coupled out to the upper waveg-86 uide at Output 1, while the TE-polarized light remains propagating in 87 the same waveguide, and is coupled out at Output 2, with an S-bend 88 89 waveguide.

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121 In order to design efficient ADC mode converters, the waveguide 90 122 widths are supposed to match the phase matching condition between the 91 123 TM₀ and TM₁ modes, which means that the effective refractive index 92 124 of the two modes are equivalent. We use Lumerical MODE Solutions 93 125 to simulate the effective refractive index in 4H-SiC waveguides, and 94 126 the anisotropic Sellmeier equation is applied to the refractive index of 95 127 4H-SiC material [18]. The effective refractive index, n_{eff} , of different 96 guided modes is simulated with varying waveguide widths at the wave-97 129 length of $\lambda = 1550$ nm, as shown in Figure 2(a). The width of the input 98 130 waveguide is selected to be $W_1 = 0.60 \,\mu\text{m}$, which only confines the TE₀ 99 and TM₀ modes. Then, the corresponding width of the multi-mode 131 100 waveguide in the ADC mode converter should be $W_2 = 1.44 \,\mu\text{m}$, to sat-101 isfy the phase matching condition, given by $n_{eff,TM_0} = n_{eff,TM_1} = 1.78$, 133 102 approximately. 103 134

After selecting the waveguide dimensions, we need to set the gap 135 104 between the two waveguides in the ADC mode converter. The power 136 105 coupling between the two waveguides in the ADC is jointly determined 137 106 by the gap and the coupling length. The gap can influence the su- 138 107 permode distribution in the coupling region, as well as the coupling 139 108 efficiency. The coupling length can be theoretically calculated through 109 140 $L = \lambda / [2(n_{even} - n_{odd})]$, where λ is the wavelength, $n_{even/odd}$ is the 141 110 effective refractive index of the even/odd supermode in the ADC mode 142 111 converter [30]. We find that as the gap is increasing, the effective 143 112 refractive index difference between the even and the odd mode is re- 144 113 114 duced. Hence, the coupling length, which is inversely proportional 145



Fig. 2. (a) Effective refractive index of TE_0 , TE_1 , TM_0 , TM_1 modes of 400 nm thick 4H-SiC waveguides as a function of the waveugide width, and the corresponding eigenmode profiles at 1550 nm. The magnetic-field profiles of (b) the even and (c) the odd TM supermodes in the coupling region.

to the effective refractive index difference, is increasing accordingly. Narrow gaps make efficient power coupling, but also require high nanofabrication resolution. Taking the trade-off between fabrication difficulty and the device compactness into consideration, we decide to choose the gap to be G = 200 nm. Figure 2(b) and 2(c) show the magnetic-field profiles of the even and the odd TM supermodes in the coupling region, respectively, with a gap of G = 200 nm. The effective refractive index of the even and the odd supermode is $n_{even} = 1.811$ and $n_{odd} = 1.763$, respectively. Thus, the coupling length is calculated to be 16 µm, when the wavelength is at $\lambda = 1550$ nm.

We also numerically simulate the coupling efficiency as a function of the coupling length, using Lumerical FDTD Solutions. The optimized coupling length is simulated to be $L = 14 \,\mu\text{m}$, considering a high transmittance of >-0.04 dB and a large polarization extinction ratio of >23 dB for both polarizations at the corresponding output, as shown in Figure 3(a) and 3(b). The polarization extinction ratio is defined by $PER_{\text{TE}} = 10\log_{10} \frac{T_{\text{TE}_{\text{Output}2}}}{T_{\text{TE}_{\text{Output}1}}}$ and $PER_{\text{TM}} = 10\log_{10} \frac{T_{\text{TM}_{\text{Output}2}}}{T_{\text{TM}_{\text{Output}2}}}$, for TE and TM polarized light, respectively, where *T* is the transmission efficiency. The simulated coupling length is slightly shorter than the calculated one. It is because the start of the S-bend waveguide is still close to the multimode waveguide, and there could be some light coupling in that region.

We then simulate the transmission spectrum with the optimized PBS, as shown in Figure 3(c) and 3(d). At Output 1, the highest transmittance of the TM polarized light is >-0.04 dB with a wide 1 dB bandwidth of 190 nm from 1450 nm to 1640 nm. Thanks to the cascaded ADC mode converters, the TE polarized light leaks weakly at Output 1, <-20 dB. At Output 2, the TE polarized light also shows very high transmittance with an ultra wide 1 dB bandwidth of >300 nm. The polarization extinction ratio is plotted in Figure 3(e). The TE and TM polarized light shows high polarization extinction ratio of



Fig. 3. Transmittance as a function of the coupling length at (a) Output 1 and (b) Output 2. Transmittance as a function of the wavelength at (c) Output 1 and (d) Output 2. (e) Polarization extinction ratio as a function of the wavelength. (f) Energy distribution of TE (upper) and (lower) TM polarized light in the PBS. It is noticed that the dimensions in x and y directions are not in scale.



Fig. 4. (a) Microscope image of the ADC mode conversion based PBS with gratings at the input and output ports, fabricated in 4H-SiCOI integrated platform. (b) Zoom-in SEM image of the SiC ADC mode conversion based PBS.

>20 dB within a bandwidth of >300 nm and 40 nm from 1530 nm to 1570 nm, respectively. The device is sensitive to the wavelength for the TM polarization, due to the dispersion induced phase mismatching. Generally, from the simulation results, the proposed ADC mode conversion based PBS shows high efficiency, high polarization extinction ratio, and satisfactory operation bandwidth covering the whole C band. The energy distribution of the TE and TM polarized light in the device is shown in Figure 3(f). It is noticed that the energy distrubution in the central waveguide is not a perfect TM₁ mode, which is due to the multimode interference with the cross-coupled TM₀ mode [26]. All the parameters of the ADC mode conversion based PBS are summarized in Table 1.

Table 1. Optimized parameters of the ADC mode conversion based PBS.

W_1 (µm)	$W_2 \ (\mu m)$	G (nm)	L (µm)
0.6	1.44	200	14

After designing the PBS, we fabricate the device on a 4H-SiCOI chip, which is prepared through the ion-cut method [7]. The devices are patterned on the chip through electron-beam lithography and inductively coupled plasma-reactive ion etching [7]. Figure 4(a) shows the microscope image of the whole fabricated SiC testing device. The input and output ports are connected with polarization-insensitive grating couplers, which can couple light with both polarizations between the device and the fiber-based characterization system [19]. Figure 4(b) shows the scanning electron microscope (SEM) image of the fabricated SiC ADC mode converter based PBS.

We build the characterization setup, shown in Figure 5(a), to test the device. A tunable continuous-wave laser source (Santec TSL-510) launches light into a polarization controller. The polarized light is coupled in and out of the device through fibers and grating couplers. The output light is detected by a power meter (Santec MPM-212). The laser and the power meter are controlled by a computer for synchronized sweep to measure the transmission spectrum. Before the characterization, a microring resonator with resonances of both polarizations is used to calibrate the input polarization, according to the different free spectral ranges [19].

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The normalized transmittance of the ADC mode conversion based PBS with both polarizations, measured from Output 1 and Output 2, is plotted in Figure 5(b) and 5(c), respectively. It is noticed that the normalized transmittance only shows the PBS performance, excluding the grating transmittance. As can be seen, the transmittance of the TM and TE polarized light at Output 1 and Output 2, respectively, is very high and flat, within the measured wavelength range between 1500 nm and 1630 nm. At Output 1, the highest transmittance of the TM polarized light is -1.3 dB at \sim 1560 nm. It can be seen there are many ripples in the spectrum, that are deemed to be generated due to the interference in the ADC, as some light is reflected by the grating couplers. In spite of the ripples, the overall transmittance is >-3.3 dB over the whole measurement range. At Output 2, the highest transmittance of the TE polarized light is -0.6 dB, and the overall transmittance is >-1.5 dB over the whole measurement range, which means that the 1 dB transmission bandwidth is beyond 130 nm. The polarization extinction ratio as a function of the wavelength is plotted in Fig.5(d). It is seen that the polarization extinction ratio for the TM polarized light as high as 35 dB is achieved at ~1560 nm, which is the phase matching wavelength. And it is >17 dB, between 1535 nm and 1595 nm, covering the whole C band and half L band. The polarization extinction ratio for the TE polarized light is also >30 dB at \sim 1560 nm, and is >20 dB within the whole measured wavelength range. Such high

polarization extinction ratio is attributed to phase mismatching of the 226 201 TE mode in the ADC, as well as the sidewall roughness induced light 202 227 scattering in the coupling region, that results in weak cross coupling. 203 228 It is seen that the simulation well predicts and guides to perform the 204 experiment, and we acquire a PBS with overall good performances, 229 205 230 except for a slight red shift of the spectrum. The phase matching 206 231 wavelength is red shifted by ~ 10 nm in the experiment, which could 207 be due to the linewidth variation and fabrication imperfections in the 208 fabrication processes. We find that a linewidth reduction of 10 nm can ²³² 209 induce a red shift of the phase matching wavelength by ~ 10 nm, from 210 233 the simulation. 211 234



Fig. 5. (a) Measurement setup schematic, CW: tunable continuouswave laser, PC: polarization controller, PM: power meter. The measured results of normalized transmittance at (b) Output 1 and (c) Output 2, and (d) polarization extinction ratio.

In conclusion, an ADC mode conversion based PBS is proposed 212 277 for the low-birefringence polarization beam splitting in SiC integrated 213 278 platforms. We numerically simulate and experimentally demonstrate 279 214 the device in a 4H-SiCOI chip, with high transmission efficiency, broad 280 215 operational bandwidth, and large polarization extinction ratio. The 281 216 measurement results show that it has high transmittance of -0.6 dB and -282 217 1.3 dB for the TE and TM polarized light, respectively. The polarization 283 218 284 extinction ratio is >32 dB at \sim 1560 nm for both polarizations, and is 219 285 >25 dB and >17 dB in the whole C band, for the TE and TM polarized 220 286 light, respectively. This work provides an effective solution for low-221 birefringence polarization beam splitting in SiC integrated platforms 222 288 with overall high performances, especially high polarization extinction 223 ratio, and paves the way for the polarization diversity applications with 224

SiC photonic integrated circuits. 225

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Data availability. Data underlying the results presented in this paper are not publicly available at this time but may be obtained from the authors upon reasonable request.

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