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Improved Vector Velocity Estimation using Directional Transverse Oscillation

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Abstract—A method for estimating vector velocities using transverse oscillation (TO) combined with directional beamforming is presented. Directional Transverse Oscillation (DTO) is self-calibrating, which increase the estimation accuracy and finds the lateral oscillation period automatically. A normal focused field is emitted and the received signals are beamformed in the lateral direction transverse to the ultrasound beam. A lateral oscillation is obtained by having a receive apodization waveform with two separate peaks. The IQ data are obtained by making a Hilbert transform of the directional signal, and a modified TO estimator can be used to find both the lateral and axial velocity. The approach is self-calibrating as the lateral oscillation period directly is estimated from the directional signal through a Fourier transform. The approach was implemented on the SARUS scanner using a BK Medical 8820e transducer with a focal point at 105.6 mm (F#=5) for Vector Flow Imaging (VFI). A 6 mm radius tube in a circulating flow rig was scanned and the parabolic volume flow of 112.7 l/h (peak velocity 0.55 m/s) measured by a Danfoss Magnetic flow meter for reference. Velocity estimates for DTO are found for 32 emissions at a 90 degrees beam-to-flow angle at a vessel depth of 30 mm. The standard deviation (SD) drops from 9.14% for TO to 5.4%, when using DTO. The bias is -5.05% and the angle is found within +/- 3.93 degrees. At 70 mm a relative SD of 7% is obtained, the bias is -1.74%, and the angle is found within +/- 2.6 degrees showing a low bias across depths.

I. INTRODUCTION

When investigating human hemodynamics, it is important to find both the velocity magnitude and direction. The velocity vector in a plane can be estimated using the transverse oscillation (TO) method described in [1], [2], [3], [4]. A similar approach has also been suggested by Anderson [5] and studied by Sumi [6] and Liebgott et al. [7], [8], [9].

In TO vector flow imaging (VFI) two beams are formed during receive processing. The two beams have to be phased shifted a quarter of the lateral wavelength. This wavelength depends on the emit focus, receive apodization function, and the interrogation depth. The TO wavelength is predicted from [1], [10]:

$$\lambda = \frac{2\lambda d}{P_d} = \frac{2\lambda d}{N_d P_i}$$

where $\lambda$ is the normal axial wavelength, $d$ is the depth, and $P_d$ is the distance between the two peaks in the apodization function. The transducer pitch is $P_i$ and the number of elements between the peaks is $N_d$. The lateral wavelength depends on depth and therefore has to be calculated for every depth to ensure an unbiased and accurate result. The equation is also only valid in the far-field or at the focus, and for a pulsed field this can introduce a significant bias and optimization has to be employed. This complicates the implementation of the approach [11].

A new method for beamforming transverse to the ultrasound direction is therefore suggested. The beams can be used for velocity estimation, and the lateral oscillation period or wavelength can also be estimated from this data making the method self-calibrating without the need for pre-calculation of the lateral wavelength.

II. DIRECTIONAL TRANSVERSE OSCILLATION APPROACH

The suggestion is to beamform a received signal in the lateral direction as shown in Fig. 1. A normal focused field is emitted, and the signals are received on all transducer elements. A beam $x(n,i)$ is then focused at the depth of interest in a direction transverse to the ultrasound propagation direction. Here $n$ is the sample number along the directional line, and $i$ is the pulse emission number. Bonnefous [12] suggested focusing a transverse beam with normal focusing and focusing along the flow direction was suggested by [13], [14], [15]. The key idea here is, however, to introduce a lateral oscillation and not have to focus along the flow.

A lateral oscillation is obtained, if a receive apodization waveform is employed as shown in Fig. 2 with two peaks separated by a distance. Two peaks in the apodization can...
A complex signal is formed from the beamformed directional signal and its Hilbert transform as
\[ r_{sq}(n,i) = x(n,i) + jy(n,i). \] (2)

The received element signals from the transducer is then Hilbert transformed in the temporal direction, and a new directional beamformed signal formed at the same depth for these data. This gives the signal \( r_{sqh}(n,i) \). Two new signals are then formed from:
\[
\begin{align*}
    r_1(n,i) &= r_{sqh}(n,i) + jr_{sqh}(n,i) \\
    r_2(n,i) &= r_{sqh}(n,i) - jr_{sqh}(n,i).
\end{align*}
\]

A simple model for the received signal is then [16], [3]:
\[
\begin{align*}
    r_1(n,i) &= k \cdot \exp \left( j2\pi \left( \frac{2v_c}{c} f_0 T_{prf} - f_0 \left( t - \frac{2d}{c} \right) \right) \\
    &\quad + \left( \frac{v_x}{\lambda} T_{prf} - n\Delta x/\lambda_x \right) \right) \\
\end{align*}
\] (3)
when assuming monochromatic signals. Here \( v_c \) is the axial velocity, \( v_x \) lateral velocity, \( c \) speed of sound, \( f_0 \) transducer frequency, \( \Delta x \) is the sampling interval along the lateral signal, \( k \) is a constant, and \( T_{prf} \) is the time between pulse emissions. The interrogation depth is \( d \) and the two frequencies received from the axial and lateral motions are given by:
\[
    f_p = \frac{2v_c}{c} f_0 = \frac{2v_c}{\lambda}, \quad f_x = \frac{v_x}{\lambda_x}. \] (4)

The changes in phase as a function of emission number for the two signals \( r_1(n,i) \) and \( r_2(n,i) \) are [3]:
\[
\begin{align*}
    d\Theta_1 &= 2\pi T_{prf} (f_x + f_p) \\
    d\Theta_2 &= 2\pi T_{prf} (f_x - f_p).
\end{align*}
\] (5)

Adding the two phase changes gives
\[
\begin{align*}
    d\Theta_1 + d\Theta_2 &= 2\pi 2T_{prf} f_x = 4\pi T_{prf} \frac{v_x}{\lambda_x} \quad (6)
\end{align*}
\]
and subtracting them gives
\[
\begin{align*}
    d\Theta_1 - d\Theta_2 &= 2\pi 2T_{prf} f_p = 4\pi T_{prf} \frac{2v_c}{\lambda} \quad (7)
\end{align*}
\]

The transverse velocity can, thus, be found directly from:
\[
    v_x = \frac{(d\Theta_1 + d\Theta_2) \lambda}{2\pi 2T_{prf}} \quad (8)
\]
and the axial velocity from
\[
    v_c = \frac{(d\Theta_1 - d\Theta_2) \lambda}{2\pi 4T_{prf}}. \quad (9)
\]

For a complex signal the phase change is determined by [16], [17]
\[
\begin{align*}
    d\hat{\Theta} &= \arctan \left( \frac{\sum_{i=0}^{N-1} y(i)x(i-1) - y(i-1)x(i)}{\sum_{i=0}^{N-1} x(i)x(i-1) + y(i)y(i-1)} \right) \\
    &= \arctan \left( \frac{3 \{R(1)\}}{\Re \{R(1)\}} \right), \quad (10)
\end{align*}
\]
using the estimated complex autocorrelation of the signal
\[
\hat{R}(m) = \frac{1}{N-m} \sum_{i=0}^{N-m} r^*(i)r(i+m), \quad (11)
\]
where \( \Im \{ R(1) \} \) denotes the imaginary part of the complex autocorrelation and \( \Re \{ R(1) \} \) the real part both at a lag of 1.

For the directional signals the autocorrelation function estimates are:

\[
\hat{R}_1(1) = \frac{1}{NN_s} \sum_{i=0}^{N-2N_s-1} \sum_{n=0}^{N-1} r_1^*(n,i)r_1(n,i+1),
\]

and

\[
\hat{R}_2(1) = \frac{1}{NN_s} \sum_{i=0}^{N-2N_s-1} \sum_{n=0}^{N-1} r_2^*(n,i)r_2(n,i+1).
\]

The autocorrelation estimates are, thus, averaged over the number of emissions \( N \) and the number of samples in the directional lines \( N_x \). This reduces the noise and improves on estimation accuracy.

The velocity estimators for the two velocity components are then:

\[
v_x = \frac{\lambda_x}{2\pi 2f_{prf}} \arctan \left( \frac{\Im \{ R_1(1) \} \Re \{ R_2(1) \} + \Im \{ R_2(1) \} \Re \{ R_1(1) \}}{\Re \{ R_1(1) \} \Re \{ R_2(1) \} - \Im \{ R_1(1) \} \Im \{ R_2(1) \}} \right),
\]

and

\[
v_z = \frac{\lambda}{2\pi 4f_{prf}} \arctan \left( \frac{\Im \{ R_1(1) \} \Re \{ R_2(1) \} - \Im \{ R_2(1) \} \Re \{ R_1(1) \}}{\Re \{ R_1(1) \} \Re \{ R_2(1) \} + \Im \{ R_1(1) \} \Im \{ R_2(1) \}} \right),
\]

when using the derivation from [3] to combine the two phase shift estimates into one equation. Here \( \Im \) denotes imaginary part and \( \Re \) real part.

The lateral wavelength needed in the estimator can be calculated by:

\[
1/\lambda_x = f_x = \frac{\sum_{m=-N/2}^{N/2} m R_{sq}(m,i)}{\sum_{m=-N/2}^{N/2} R_{sq}(m,i)}
\]

where \( R_{sq}(m,i) \) is the Fourier transform of the complex directional signal \( r_{sq}(n,i) = x(n,i) + jy(n,i) \) along the sample direction \( n \). This makes the approach self-calibrating as in [18] for axial velocity estimation. The lateral wavelength \( \lambda_x \) can be calculated for the different depths and the estimator automatically yields an unbiased estimate of the transverse velocity. This frequency estimate can also be improved by averaging the estimator over all emissions as

\[
f_x = \frac{1}{N} \sum_{i=0}^{N-1} \frac{N/2 \sum_{m=-N/2}^{N/2} m R_{sq}(m,i)}{N/2 \sum_{m=-N/2}^{N/2} R_{sq}(m,i)}.
\]

The beamforming can be performed for a normal focused emission, for synthetic aperture flow imaging, and for plane wave imaging. The transmit and receive apodization function can be changed as a function of depth to obtain the highest possible \( f_x \). In synthetic aperture and plane wave imaging the combined transmit apodization function can also be manipulated to further increase the lateral oscillation frequency.

Fig. 4. Dynamic apodization profiles as a function of depth for the receive processing.

IV. PERFORMANCE

The velocity estimation approach has been implemented on the SARUS experimental scanner [19]. A BK Medical 8820e convex array transducer was employed and vector flow imaging (VFI) in a single direction was interleaved with a B-mode image. An active aperture of 64 elements was used during transmit for both sequences. The focal point was at 42 mm (F# = 2) for the B-mode and 105.6 mm (F# = 5) for VFI. The transducer has 192 elements with \( \lambda \) pitch and the B-mode image consisted of 129 lines. A 6 mm radius tube in a circulating flow rig was scanned and the volume flow was measured by a Danfoss Magnetic flow meter for reference. The volume flow was 112.7 l/h corresponding to a peak velocity in the vessel of 0.55 m/s. The pulse repetition frequency was 4 kHz and the beam-to-flow angle is 90 degrees.

The beamforming employed a double Hamming apodization on the receiving aperture and the F-number in receive was 2. The receiving aperture was expanded as a function of depth to maintain a constant F-number and thereby a constant \( \lambda_x \). The resulting apodization profiles are shown in Fig. 4.

An example of velocity estimation for the traditional TO approach is shown in Fig. 5 and the corresponding estimates for DTO is shown in Fig. 6 when using 32 emissions. The standard deviation on the estimates drops from 9.14% to 5.4% when using the new approach.

The self-calibrating feature of the approach is shown in Fig. 7, where data from a vessel at 70 mm is used. The same beamforming scheme is used and a relative standard deviation of 7% is obtained with a bias of -1.7% and the angle is found within 2.6 degrees.

V. DISCUSSION AND CONCLUSION

A new method for making TO vector flow estimation self-calibrating has been suggested. The techniques find the lateral oscillation period from the actual measured data and uses this in the velocity estimation. This makes it possible to have any apodization waveform without having to predict the lateral oscillation period, and this maintains a low bias of the estimates for different depths. The approach also makes it possible to average over the lateral signal to decrease variance. Results from a circulating flow rig showed a decrease of
standard deviation at 30 mm from 9.14% to 5.4% and the bias could be maintained below 5% for an interrogation depth from 30 to 70 mm.

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