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Thermodynamic Modeling of the Sr-Co-Fe-O System

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Abstract

This paper reviews and assesses phase equilibria and thermodynamic properties of phases in the Sr-Co-Fe-O system, with a focus on oxides, especially the SrCo 1−x Fe x O3−δ perovskite. In our work, the SrCo 1−x Fe x O3−δ perovskite was modeled with a three-sublattice model, where the three sublattices correspond to the A, B and oxygen sites in an ABO3 perovskite, respectively. A number of other important ternary oxide phases in Sr-Co-O and Sr-Co-Fe-O were also considered. Available thermodynamic and phase diagram data were carefully assessed. A thermodynamic description of Sr-Co-O was derived using the CALPHAD approach and was further extrapolated to that of Sr-Co-Fe-O. The thermodynamic database of Sr-Co-Fe-O established in this work allows for calculating phase diagrams, thermodynamic properties, cation distribution and defect chemistry properties, and therefore enables material composition optimization for various applications, including solid oxide fuel cells and oxygen membranes.

Key words: strontium cobaltite, phase diagram, perovskite.

1. Introduction

Strontium cobaltites offer good electronic and ionic conductivities that enable their application in solid oxide fuel cells (SOFCs) and oxygen permeation membranes [1]. Sr2−xFe2−xCo x O7−δ and SrCo 1−xFe x O3−δ in Sr-Co-Fe-O have also received special attention due to their magnetic properties, notably the magnetoresistance (MR) effect [2, 3]. Both the magnetic and the electrical properties of these oxides depend strongly on composition and preparation methods [4]. Though extensive studies have been carried out on the effect of composition and sample preparation on the structure and physical properties of the oxides in Sr-Co-Fe-O, the phase relationship and thermodynamic properties have not been well established.

In this research, thermodynamic databases of Sr-Co-O and Sr-Co-Fe-O were developed as part of our project to develop a thermodynamic database of La-Sr-Co-Fe-O for SOFC applications [5–7]. The literature data available on Sr-Co-O and Sr-Co-Fe-O were carefully reviewed. A thermodynamic description of Sr-Co-O was derived using the CALculation of PHAse Diagrams (CALPHAD) approach and was combined with a description of Sr-Fe-O developed by Povoden-Karadeniz et al. [8] and further extrapolated to Sr-Co-Fe-O. Various phase diagrams and thermodynamic properties were calculated and then compared with experimental data reported in the literature. The potential applications of the new database were examined.
2. Literature review

The Sr-Co-Fe-O system contains the following binary oxides: SrO, SrO₂, CoO, Co₃O₄, FeO, Fe₂O₃ and Fe₃O₄. Detailed information on the binary oxides can be found in previous modeling work [9–11]. There are five ternary oxides in Sr-Fe-O: SrFeO₃₋δ, Sr₂FeO₄₋δ, Sr₃FeO₇₋δ, Sr₄Fe₂O₁₀₋δ, Sr₄Fe₃O₁₃₋δ [8]. Two ternary oxides have been reported in Co-Fe-O: halite (CoₓFe₁₋ₓO) and spinel (CoₓFe₀·₃₋ₓO₄) [12]. The reported ternary and quaternary solid oxide phases for Sr-Co-O and Sr-Co-Fe-O were reviewed in this research. Below is a short summary of results from the relevant literature.

2.1. Sr-Co-O

Previous studies of phase equilibria in Sr-Co-O have focused on compositions related to the perovskite phase. The stability of other oxide phases, such as the Ruddlesden-Popper phase (Sr₃Co₂O₇₋δ), has not been fully studied. Takeda et al. [13] investigated phase relations in Sr-Co-O. They studied phase composition as a function of temperature in various atmospheres using X-ray diffraction (XRD), thermogravimetry (TG), and differential thermal analysis (DTA). They reported the stability ranges of cubic perovskite, brownmillerite and structure undefined low temperature phases. Rodriguez et al. [14] studied phase transition in Sr₂Co₂O₅ using neutron powder diffraction in an atmosphere of PO₂ ≈ 10⁻⁴ bar with continuously changing temperature. They were unable to correctly identify the low temperature phases [15]. Vashook et al. [16, 17] explored phase transformation temperatures for strontium cobaltite (SrCoO₃) in a PO₂ range of 50 – 400 Pa using XRD and TG. The stability regions for cubic perovskite, brownmillerite and low temperature phases (Sr₆Co₅O₁₅) at low PO₂ were determined. They described the transformation from cubic perovskite to brownmillerite as second-order and the one from brownmillerite to Sr₆Co₅O₁₅ and Co₃O₄ as first-order.

2.1.1. SrCoO₃₋δ
SrCoO₃₋δ is a fully disordered oxygen-deficient cubic perovskite phase with Pm3m symmetry. The stoichiometric SrCoO₃ can be synthesized only under high oxygen partial pressure (>2000 bar) [13, 18]. SrCoO₃₋δ has been extensively studied due to its attractive magnetic [18, 19] and electric properties [19] and its oxygen mobility [20]. Taguchi et al. [18] investigated the effect of oxygen deficiency on the magnetic properties of SrCoO₃₋δ (0 < δ < 0.5). They found a strong dependence of the Curie temperature on the oxygen deficiency. They determined the Co⁴⁺ content at low temperatures (< 623 K) and high oxygen pressures (50 – 2600 bar) assuming the presence of Co⁴⁺ and Co³⁺ only. Nakatsuka et al. [20] prepared a single crystal SrCoO₃₋δ at high temperature in an O₂ gas flow and reported the average Co valency as +3.28 using XRD. Co³⁺ and Co⁴⁺ ions co-exist in the oxygen-deficient SrCoO₃₋δ cubic perovskite. Takeda et al. [13] investigated the stability of SrCoO₃₋δ at a few chosen temperatures and atmospheres. Rodriguez et al. [14] reported the transformation temperature from SrCoO₃₋δ to Sr₂Co₂O₅ as 1113 K at PO₂ ≈ 10⁻⁴ bar. Vashook et al. [16] determined the stability range of SrCoO₃₋δ at various temperatures using DTA and XRD. They [17, 21] further investigated the oxygen non-stoichiometry and electrical conductivity of SrCoO₃₋δ in a temperature range of 1223 – 1323 K and an oxygen partial pressure range of 10⁻⁵ – 10⁻² bar by solid electrolyte coulometry and resistivity measurements. They reported a large decrease in resistivity accompanying with the transition from Sr₆Co₅O₁₅ to SrCoO₃₋δ. Le Toquin et al. [22] investigated crystal structure evolution from SrCoO₂·₅ to SrCoO₃ in situ by neutron diffraction as well as by X-ray absorption fine structure (XAFS) spectroscopy. From the neutron diffraction experiments, the structure of two intermediate phases, SrCoO₂·₇₅ and SrCoO₂·₈₂₀·₀₇, were characterized. Calle-Vallejo et al. [23] reported the Gibbs energy of formation for SrCoO₃ as −163 kJ/mol at 298 K obtained from first-principle calculations.

2.1.2. Sr₂Co₂O₅
Sr₂Co₂O₅ has a brownmillerite-type structure, which is a close relative of the perovskite structure with ordering of vacancies. Grenier et al. [24] investigated the crystal structure of Sr₂Co₂O₅ at 1183 – 1373 K and found that, on cooling, it undergoes an exothermic reaction and transforms into a hexagonal phase at 1183 K. Takeda et al. [13] reported that Sr₂Co₂O₅ is stable in a narrow temperature range at around 1073 – 1273 K in N₂. With increasing temperature, it transforms into the SrCoO₃₋δ perovskite phase at about 1373 K in N₂. They also found ordered oxygen vacancies in Sr₂Co₂O₅₋δ (δ = 0.04 – 0.16). No Fe solubility in Sr₂Co₂O₅ was reported. Saal [25] determined the Gibbs energy of Sr₂Co₂O₅ obtained via first-principle calculations.

2.1.3. Sr₆Co₅O₁₅

Sr₆Co₅O₁₅ has previously been described as a low-temperature rhombohedral/hexagonal form of Sr₂Co₂O₅ [24, 26, 27]. Takeda et al. [13] propose that it is cobalt-deficient 2H-SrCo₁₋ₓOₓ (x ≈ 0.1). Harrison et al. [15] were the first to confirm the existence of Sr₆Co₅O₁₅ with a stoichiometric composition and Sr₆Co₅O₁₅ is a 2H-hexagonal-type phase. When high temperature brownmillerite Sr₂Co₂O₅ cools in air, it decomposes into a two-phase mixture of Sr₆Co₅O₁₅ and CoOₓ [15]. No Fe solubility in Sr₆Co₅O₁₅ is reported in the literature. Its crystallographic, magnetic, and electronic properties have been extensively studied by Sun et al. and Iwasaki et al. [28, 29]. No experimental study was carried out on the thermodynamic properties of Sr₆Co₅O₁₅. Saal et al. [30] reported on its heat capacity and entropy up to 1300 K from first-principle calculations using the Debye-Grüneisen model and derived the Gibbs energy function for Sr₆Co₅O₁₅.

A few other strontium cobaltites with a hexagonal perovskite-like structure have been reported, with a Co content of 42.8 – 44.5 cation% (cat.%), including Sr₂₄Co₁₉O₅₇, Sr₁₄Co₁₁O₃₃, Sr₄Co₃O₉ and Sr₅Co₄O₁₂ [31-33]. The existence of Sr₂₄Co₁₉O₅₇ and Sr₁₄Co₁₁O₃₃ has been denied by Li et al. and Aksenova et al. [32, 34]. The reported Sr₄Co₃O₉ and Sr₅Co₄O₁₂ have XRD patterns and stability ranges similar to those of Sr₆Co₅O₁₅ [33]. The existence of Sr₄Co₃O₉ and Sr₅Co₄O₁₂ is doubtable and further verified in this work.

2.1.4. Sr₃Co₂O₇₋δ

Sr₃Co₂O₇₋δ is an oxygen-deficient Ruddlesden-Popper type phase. Its crystal structure was characterized by Dann and Weller [35] using X-ray and neutron diffraction and was reported as orthorhombic-type. They indicated that this phase can be obtained at high temperatures (1273 – 1473 K) and ambient pressure. At low temperatures and/or high P₀₂, Sr₃Co₂O₇₋δ decomposes into a mixture of Sr₆Co₅O₁₅, SrCoO₃₋δ and SrO. The oxygen deficiency (δ in Sr₃Co₂O₇₋δ) was reported to be around 1 ± 0.2, corresponding to an average Co oxidation state of around +3 ± 0.2. The crystal structure and magnetic property of this phase have been reported in various studies [36–38]. It is possible to synthesize other Ruddlesden-Popper type phases (Srₙ₊₁CoₙO₃ₙ₊₁, n≤4) under 6 GPa at 1273 – 1773 K [39]. However, they are not stable under ambient conditions.

2.2. Sr-Co-Fe-O

The phase diagram data for the Sr-Co-Fe-O system are rather limited. Aksenova et al. [34] studied the phase equilibria in the Sr-Co-Fe-O system at 1373 K in air using 68 samples with various compositions. They determined the stability region and crystal structure of solid solutions in the Sr-Co-Fe-O system using XRD and further constructed a 1373 K isothermal section of SrO-CoO-Fe₂O₃ in air. Fossdal et al. [40] studied the Sr₄FeₓCo₁₋ₓO₁₃₋δ phase in the range of 1048–1493 K using XRD and DTA and provided a Fe rich part of the isothermal section of SrO-CoO-Fe₂O₃ at 1173 K in air.

2.2.1. SrCo₁₋ₓFeₓO₃₋δ
SrFeO$_3$-$\delta$ has a tetragonal perovskite type structure. By substituting Fe with Co, the crystal structure changes from tetragonal to cubic [34, 41]. Aksenova et al. [34] determined the stability range of SrCo$_{1-x}$Fe$_x$O$_3$-$\delta$ by XRD and reported that the perovskite phase is stable within the range of 0 $\leq x \leq 0.7$ at 1373 K in air. The oxygen non-stoichiometry and phase stability of SrCo$_{1-x}$Fe$_x$O$_3$-$\delta$ have been studied by a number of groups [42–46], focusing on the composition SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$. Liu et al. [43] studied the temperature range of 823 – 1263 K and $P_{O_2}$ = 0.015 – 1 bar using TG and XRD, while Grunbaum et al. [44] chose a temperature interval of 823 – 1223 K and $P_{O_2}$ = 10$^{-5}$ – 1 bar using TG and high temperature XRD (HT-XRD). Moreover, McIntosh et al. [45] investigated the stability of SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$ at $P_{O_2} = 10^{-5}$ – 1 bar using in situ neutron diffraction. Depending on temperature and oxygen partial pressure, SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$ is either a single phase perovskite or a two-phase mixture of perovskite and brownmillerite. The cubic SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$ perovskite single phase is stable at $T > 1073$ K and $P_{O_2} = 10^{-5}$ – 1 bar and its stability range extends to lower temperatures with increasing oxygen partial pressure. However, the stability range determined for SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$ from these studies is unreliable, because the samples used were most likely not equilibrated due to short (< 24 h) annealing time. Moreover, they were unable to distinguish the hexagonal Sr$_6$Co$_{15}$O$_{33}$ phase from the SrCo$_{1-x}$Fe$_x$O$_{3-\delta}$ perovskite and Sr$_2$Co$_2$O$_5$ brownmillerite phases. McIntosh et al. [45, 46] measured the oxygen non-stoichiometry of SrCo$_{1-x}$Fe$_x$O$_{3-\delta}$ using TG and DTA, but the results among these two studies were inconsistent.

Mitchell et al. [47] investigated the stability of SrCo$_{0.8}$Fe$_{0.2}$O$_{3-\delta}$ at 1173 K and $P_{O_2} = 10^{-0.67}$ – 10$^{-13.3}$ bar, and found that it is stable for the entire $P_{O_2}$ range. Vashuk et al. [48] studied the oxygen non-stoichiometry and defect chemistry of SrCo$_{0.25}$Fe$_{0.75}$O$_{3-\delta}$ using solid-electrolyte coulometry, TG, and conductivity measurements at 573 – 1273 K and $P_{O_2} = 10^{4}$ – 1 bar. SrCo$_{0.25}$Fe$_{0.75}$O$_{3-\delta}$ was shown to be p-type throughout the entire experimental oxygen partial pressure and temperature range. Perovskite to brownmillerite phase transformation was reported.

2.2.2. Solid solution Sr$_3$Fe$_{2-x}$Co$_x$O$_{7-\delta}$

The electronic, magnetic and magneto-resistance properties of Sr$_3$Fe$_{2-x}$Co$_x$O$_{7-\delta}$ were studied by Veith et al. [2]. They reported that single-phase Sr$_3$Fe$_{2-x}$Co$_x$O$_{7-\delta}$ exists in a composition range of 0.25 $\leq x \leq 1.75$ at 1273 K under flowing O$_2$. Aksenova et al. [34] reported single phase Sr$_3$Fe$_{2-x}$Co$_x$O$_{7-\delta}$ at 1373 K in air (0 $\leq x \leq 0.4$). At x > 0.4, they detected formation of Sr$_3$Co$_{2}$O$_{7-\delta}$.

2.2.3. Solid solution Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$

The Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ solid solution has an orthorhombic structure (space group $Iba2$) [49, 50]. The phase stability of Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ and its electrical conductivity and oxygen permeability have been investigated by a number of groups. A certain amount of Co can be dissolved in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ without forming a secondary phase. Beyond the solubility limit, the Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ single phase is replaced by a three-phase mixture of Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$, SrCo$_{1-x}$Fe$_x$O$_{3-\delta}$ and Co$_x$Fe$_{1-x}$O or Co$_x$Fe$_{3-x}$O$_4$ [50–52]. Co in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ affects not only the phase stability but also the oxygen permeability and electrical conductivity. Both oxygen permeability and electrical conductivity increase with increasing x in either the single phase region or the three-phase region [51–53]. Various solubility limits of Co in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ have been reported. Xia et al. [50] reported a maximum Co solubility corresponding to x = 1.8 at 1273 K. The solubility decreases with increasing temperature in the range 1273 – 1373 K. Armstrong et al. [53] reported a Co solubility limit of x = 1.5 in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ at 1273 K. Deng et al. [52] determined x = 1.5 in samples slowly cooled from1473 K. They reported an average valence state of Co and Fe as around 2.9 ~ 3.0. Moreover, values of x = 1.4, ~1.6, 1.6 were reported by Kim et al. [54] in samples sintered at 1423 K, by Ma et al. [51] in samples sintered at 1473 K, and by Aksenova et al. [34] in samples sintered at 1373 K, respectively. Fossdal et al. [40] systemically measured the Co solubility in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ at several temperatures. They concluded that Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ is stable only in a narrow temperature range of 1048–1493 K. The highest Co solubility was obtained in samples sintered at 1173 K, and
the maximum Co-content was close to Sr$_4$Fe$_4$Co$_2$O$_{13}$. At higher or lower sintering temperatures, Co solubility decreases. This scattering of experimental Co solubility limits is probably due to the slow reaction kinetics for the transition between the single phase Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$ and the three-phase mixture of Sr$_4$Fe$_{6-x}$Co$_x$O$_{13-\delta}$, SrFe$_{1-x}$Co$_x$O$_{3-\delta}$ and Co$_x$Fe$_{1-x}$O or Co$_x$Fe$_{3-x}$O$_4$ during heating or cooling [40].

2.2.4. Other ternary oxides

For the other ternary oxides, no solubility of a third cation has been reported, e.g. no solubility of Sr in Co$_x$Fe$_{1-x}$O$_{1+\delta}$ or in Co$_x$Fe$_{3-x}$O$_4$ and no Co solubility in SrFe$_{12}$O$_{19-\delta}$.

3. Thermodynamic modeling

The new thermodynamic description of Sr-Co-O presented here provides descriptions of all observed equilibrium phases. The model description of the perovskite phase solves a problem of too high entropy at high temperatures assessed by Saal [25]. Our database is based on thermodynamic descriptions of the sub-systems Sr-O (by Risold et al. [9]) and Co-O (by Chen et al. [10]) and extended to Sr-Co-Fe-O by adopting recent descriptions of Sr-Fe-O by Povoden-Karadeniz [8] and of Co-Fe-O by Zhang and Chen [12]. Compound Energy Formalism (CEF) [55], which is widely used in CALPHAD assessments, was employed for modeling the Gibbs energy for all the phases in Sr-Co-Fe-O. The lattice stability for pure elements was adopted from Dinsdale [56]. The models used for the oxide phases in Sr-Co-O and Sr-Co-Fe-O are described in detail in the following sections. Table 1 provides meanings of the physical variables included in the model descriptions.

<table>
<thead>
<tr>
<th>$T$</th>
<th>Absolute temperature, K</th>
<th>$\omega A$</th>
<th>Gibbs energy of phase θ with constituent A occupying the first sublattice and B the second sublattice</th>
</tr>
</thead>
<tbody>
<tr>
<td>$R$</td>
<td>Gas constant, 8.31451 J mol$^{-1}$ K$^{-1}$</td>
<td>$y_i$</td>
<td>Site fraction of constituent $i$</td>
</tr>
<tr>
<td>$G_m^\theta$</td>
<td>Gibbs energy of phase θ per mole component</td>
<td>$H_A^{SER}$</td>
<td>Stable element reference (SER): reference state for pure element A at 298.15 K and 1 atm.</td>
</tr>
<tr>
<td>$E_G^\theta$</td>
<td>Excess Gibbs energy of phase θ per mole component. The excess Gibbs energy accounts for contributions due to non-ideal interaction between the components.</td>
<td>$L_{m,i,j}$</td>
<td>Interaction parameter for the θ phase, accounting for the non-ideal interaction between the constituents $i_m$ and $i_j$ on the first sublattice with the constituent $j$ on the second sublattice.</td>
</tr>
</tbody>
</table>

3.1. Liquid

The description for the liquid phase was ideally extrapolated from the ionic two–sublattice modeling [57, 58] of the subsystems (Sr-O [9], Co-O [10] and Fe-O [11]). The ionic liquid model assumes that cations occupy one sublattice and anions occupy the other. Hypothetical vacancies are introduced on the anion sublattice in order to extend the description to a liquid with only cations (i.e. a metallic liquid). Further detail on the model can be found in the paper by Hillert et al. [57] and Sundman [58]. In the current work, the liquid phase in Sr-Co-Fe-O is described as:

$$(\text{Co}^{2+}, \text{Co}^{3+}, \text{Fe}^{2+}, \text{Fe}^{3+}, \text{Sr}^{2+})_p(\text{O}^{2-}, \text{Vaq}^{-})_q$$

where Vaq$^{-}$ represents hypothetical vacancies with a valency corresponding to the average for the cations on the other sublattice,
The Gibbs energy of the liquid phase is expressed as:

\[ G^L_m = q \sum_{i} y_i y_{Va} G^L_{i, Va} + \sum_{i} y_i y_{O^{2-}} G^L_{i, O^{2-}} + pRT \sum_{j} y_j \ln y_j + qRT \sum_{j} y_j \ln y_j + E G^L_m \]

(3)

where \( i \) represents the constituents in the first sublattice, \( j \) represents the constituents in the second sublattice.

The excess Gibbs energy \( E G^L_m \) is formulated as:

\[ E G^L_m = \sum_{i_m} \sum_{i_n} y_{i_m} y_{i_n} (y_{O^{2-}} L^L_{i_m,i_n,O^{2-}} + q y_{Va} y_{O^{2-}} L^L_{i_m,i_n,3O}) \]

(4)

where \( i_m \) and \( i_n \) represent the constituents in the first sublattice. In the above expressions, colons were used to separate species on different sublattices and commas to separate species on the same sublattice. No ternary or quaternary parameters were optimized for the liquid phase due to lack of experimental data, so the calculated liquidus for Sr-Co-O and Sr-Co-Fe-O should be treated with caution.

3.2. \( \text{Sr}_2\text{Co}_2\text{O}_5 \) and \( \text{Sr}_6\text{Co}_5\text{O}_{15} \)

This research treated \( \text{Sr}_2\text{Co}_2\text{O}_5 \) and \( \text{Sr}_6\text{Co}_5\text{O}_{15} \) as stoichiometric compounds. Their Gibbs energy functions were derived from the first-principle calculation results of Saal et al. [25, 30], with the enthalpy and entropy terms further optimized using phase diagram data.

3.3. \( \text{Sr}_3\text{Co}_2\text{O}_7 \)

\( \text{Sr}_3\text{Co}_2\text{O}_7 \) was modeled with a 5-sublattice model reflecting its actual crystal structure. The model reads as follows:

\[ (\text{Sr}^{2+})_1(\text{Sr}^{2+})_2(\text{Co}^{3+})_2(\text{O}^{2-})_6(\text{O}^{2-})_4 \]

where \( \text{Va} \) represents vacancy (in this case oxygen ion vacancy). According to the Compound energy model (for modeling solid phases), vacancies with zero net charge are introduced on either the cation or oxygen ion sublattice or both according to the actual crystal structure information (i.e. sublattice occupation). The Gibbs energy for the end member \( (\text{Sr}^{2+})_1(\text{Sr}^{2+})_2(\text{Co}^{3+})_2(\text{O}^{2-})_6(\text{O}^{2-})_4 \) is formulated as

\[ a G_{\text{Sr}^{2+} \text{Sr}^{2+} \text{Co}^{3+} \text{Co}^{3+} \text{O}^{2-} \text{O}^{2-}} = 3 a G_{\text{Sr}^{2+}} + 2 a G_{\text{Co}^{3+}} + a G_{\text{O}^{2-}} + A_0 + B_0 * T \]

(5)

and for the other one \( (\text{Sr}^{2+})_1(\text{Sr}^{2+})_2(\text{Co}^{3+})_2(\text{Va})_6(\text{O}^{2-})_4 \) as

\[ a G_{\text{Sr}^{2+} \text{Sr}^{2+} \text{Co}^{3+} \text{Va} \text{O}^{2-}} = 3 a G_{\text{Sr}^{2+}} + 2 a G_{\text{Co}^{3+}} + 2 a G_{\text{O}^{2-}} \]

(6)

where \( a G_{\text{Sr}^{2+}} \), \( a G_{\text{Co}^{3+}} \) and \( a G_{\text{O}^{2-}} \) were taken from Risold et al. [9], Chen et al. [10] and Dinsdale [56], respectively, and \( A_0 \), \( B_0 \) values were optimized using phase diagram data.

3.4. Perovskite \( \text{SrCo}_{1-x}\text{Fe}_x\text{O}_{3-\delta} \)
A three-sublattice model was used for the perovskite phase, with the first sublattice (A site) for Sr cation (Sr\(^{2+}\)) and vacancy, the second (B site) for Co and Fe cations (Co\(^{2+}\), Co\(^{3+}\), Co\(^{4+}\), Fe\(^{2+}\), Fe\(^{3+}\) and Fe\(^{4+}\)) and vacancy, and the third (O site) for oxide ion and oxygen vacancy. Just as in the Sr-Fe-O system [8], cation and oxygen vacancies (Va) were introduced to the A/B and O site, respectively, in order to model oxygen non-stoichiometry and to maintain charge neutrality. The model reads as follows:

(Sr\(^{2+}\),Va\(_{i}\))(Co\(^{2+}\),Co\(^{3+}\), Co\(^{4+}\), Fe\(^{2+}\),Fe\(^{3+}\), Fe\(^{4+}\),Va\(_{i}\))(O\(^{2−}\), Va\(_{i}\))\(_{3}\)

The Gibbs energy of the perovskite phase is given by the following expression:

\[
G_{\text{perovskite}} = \sum_i \sum_j \sum_k y_i y_j y_k G_{i,j,k}^{\text{perovskite}} + RT \sum_i y_i \ln y_i + RT \sum_j y_j \ln y_j + 3RT \sum_k y_k \ln y_k + E_{\text{perovskite}}
\]

where \(i, j, k\) represent the constituents in the first, second and third sublattices, respectively.

According to the current model, 28 end-members (\(G_{i,j,k}^{\text{perovskite}}\) terms) need to be assigned a Gibbs energy value. In order to be consistent with the perovskite model used in Sr-Fe-O [8] and La-Co-O [59], the same Gibbs energy terms must be applied to all the common end-members. The thermodynamic descriptions of 16 common end-members were taken from Sr-Fe-O [8] while the other six were taken from La-Co-O [59]. The remaining six end-members belong to the Sr-Co-O system and were determined during this research. The strategy to obtain the Gibbs energy terms for these six end-members was to choose appropriate neutral end-members or their combinations as model parameters, which could be optimized using experimental data. As in the previous modeling of the perovskite phase [8, 60–62], a number of the most important end-members or their combinations were chosen as model parameters and these are listed below:

Stoichiometric SrCoO\(_3\) (Sr\(^{2+}\))\(_{1}\)(Co\(^{4+}\))\(_{1}\)(O\(^{2−}\))\(_{3}\). Its Gibbs energy function is given by:

\[
\alpha G_{\text{SrCOO}_3} = \alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{4+} \cdot \text{O}_3^{-}} = \alpha G_{\text{Sr}^{2+}} + \alpha G_{\text{Co}^{4+}} + 0.5 \alpha G_{\text{O}_3^{-}} + A_1 + B_1 * T
\]

where \(A_1, B_1\) are parameters optimized to reproduce thermodynamic and phase diagram data.

Reduced SrCoO\(_{3−\delta}\) (Sr\(^{2+}\))\(_{1}\)(Co\(^{3+}\))\(_{1}\)(O\(^{2−5/6}\), Va\(_{1/6}\))\(_{3}\). Its Gibbs energy function is given by:

\[
G_{\text{a}} = \frac{5}{6} \alpha G_{\text{SrCOO}_3} + \frac{1}{6} \alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{3+} \cdot \text{O}_{5/6}^{-}} + 3RT \left( \frac{5}{6} \ln \frac{5}{6} + \frac{1}{6} \ln \frac{1}{6} \right) = \alpha G_{\text{Sr}^{2+}} + \alpha G_{\text{Co}^{3+}} + 0.25 \alpha G_{\text{O}_{5/6}^{-}} + A_1 + B_1 * T
\]

and further reduced SrCoO\(_{3−\delta}\) (Sr\(^{2+}\))\(_{1}\)(Co\(^{3+}\))\(_{1}\)(O\(^{2−2/3}\), Va\(_{1/3}\))\(_{3}\)

\[
G_{\text{m}} = \frac{2}{3} \alpha G_{\text{SrCOO}_3} + \frac{1}{3} \alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{3+} \cdot \text{Va}_1} + 3RT \left( \frac{2}{3} \ln \frac{2}{3} + \frac{1}{3} \ln \frac{1}{3} \right) = \alpha G_{\text{Sr}^{2+}} + \alpha G_{\text{Co}^{3+}} + A_1 + B_1 * T
\]

where \(A_2, B_2, A_3\) and \(B_3\) are parameters to be optimized.

All the other end-members were correlated by the following reciprocal relations:

\[
\alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{3+} \cdot \text{O}_2^{-}} - \alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{3+} \cdot \text{Va}_1} - \alpha G_{\text{Sr}^{2+} \cdot \text{Co}^{3+} \cdot \text{O}_2^{-}} = \Delta G_1
\]
Where $\Delta G_x (x = 1, 2, 3)$ is the energy of the above the reciprocal reactions (Eqs. 11-13) and is called reciprocal energy. It determines the morphology of Gibbs energy surface of the phase, with a flat surface when $\Delta G = 0$ and a curved one when $\Delta G > 0$ (indicating a tendency for demixing). Further detail on reciprocal terms can be found in Hillert [55] and Povoden-Karadeniz et al. [63]. In this research, the reciprocal energy $\Delta G_x = 0 (x = 1, 2, 3)$ was chosen, following the same treatment by Povoden-Karadeniz [63], as no tendency of demixing was reported for the nonstoichiometric oxide solid solutions that were treated in this study. The Gibbs energy functions for the 6 end-members can be derived by solving Equations 8 to 13.

The excess Gibbs energy $E_{\text{perovskite}}^m$ is formulated as:

$$E_{\text{perovskite}}^m = \sum_{i_k} \sum_{i_l} \sum_{j_m} L_{i_k j_m} y_i y_j y_m (O_{i_k j_m} - L_{i_k j_m})$$

where $i_k, i_l$ represent the constituents in the first sublattice (A site), and $j_m, j_n$ represent the constituents in the second sublattice (B site).

3.5. $\text{Sr}_2\text{Fe}_{2-x}\text{Co}_x\text{O}_{7-\delta}$ and $\text{Sr}_4\text{Fe}_{6-x}\text{Co}_x\text{O}_{13-\delta}$

To model the Co solubility in these two phases, Co cation was introduced to the Fe site. As the average Co valence state reported is around +3, only Co$^{3+}$ was included to keep the model simple. The Gibbs energy functions for $\text{Sr}_2\text{Fe}_{2-x}\text{Co}_x\text{O}_{7-\delta}$ and $\text{Sr}_4\text{Fe}_{6-x}\text{Co}_x\text{O}_{13-\delta}$ can be found in Table 2.

3.6. Other ternary oxides phases

Two more ternary oxides in Sr-Co-O were included in the present work: $\text{Sr}_3\text{Co}_3\text{O}_9$ and $\text{Sr}_5\text{Co}_4\text{O}_{12}$ were modeled as stoichiometric compounds. In Sr-Fe-O, the description for $\text{Sr}_4\text{Fe}_3\text{O}_{10}$ and $\text{SrFe}_2\text{O}_{19}$ was taken from Sr-Fe-O [11] without modification, as no Co solubility in these two compounds was reported in the literature.

Table 2 Parameters obtained for the ternary oxide phases in the Sr-Co-O system and the quaternary solid solution phases in the Sr-Co-Fe-O system

<table>
<thead>
<tr>
<th>Phase</th>
<th>Parameters</th>
</tr>
</thead>
<tbody>
<tr>
<td>Perovskite$^b$</td>
<td>(Sr$^{2+}$, Va)$_1$(Co$^{3+}$,Co$^{4+}$,Co$^{5+}$,Fe$^{3+}$,Fe$^{4+}$,Va)$_1$(O$^{2-}$, Va)$_3$</td>
</tr>
<tr>
<td>$\gamma G_{\text{perovskite}}$</td>
<td>$-H_{\text{Sr}}^\text{SER} - H_{\text{Co}}^\text{SER} - 3H_{\text{O}}^{\text{SER}} = +\text{GS}_2\text{OV} + \text{GHSEROO} + 15.8759T$</td>
</tr>
<tr>
<td>$\gamma G_{\text{perovskite}}$</td>
<td>$-H_{\text{Sr}}^\text{SER} - H_{\text{Co}}^\text{SER} - 3H_{\text{O}}^{\text{SER}} = +0.5\text{GHSEROO} + \text{GS}_3\text{OV} + 11.2379T$</td>
</tr>
<tr>
<td>$\gamma G_{\text{perovskite}}$</td>
<td>$-H_{\text{Sr}}^\text{SER} - H_{\text{Co}}^\text{SER} - 3H_{\text{O}}^{\text{SER}} = +\text{G}_4\text{O}$</td>
</tr>
<tr>
<td>$\gamma G_{\text{perovskite}}$</td>
<td>$H_{\text{Sr}}^\text{SER} - H_{\text{Co}}^\text{SER} = +\text{GS}_2\text{OV} - 2\text{GHSEROO} + 15.8759T$</td>
</tr>
<tr>
<td>$\gamma G_{\text{perovskite}}$</td>
<td>$H_{\text{Sr}}^\text{SER} - H_{\text{Co}}^\text{SER} = +\text{GS}_3\text{OV} - 2.5\text{GHSEROO} + 11.2379T$</td>
</tr>
</tbody>
</table>
Sr$_2$Co$_2$O$_5$  
\[ (\text{Sr}^{2+})_2(\text{Co}^{4+})_1(\text{Co}^{3+})_1(\text{O}^{2-})_5 \]

$G_{\text{Sr}_2\text{Co}_2\text{O}_5} = -38661.9 \text{kJ mol}^{-1}$  

Sr$_3$Co$_2$O$_7$  
\[ (\text{Sr}^{2+})_3(\text{Co}^{3+})_2(\text{O}^{2-})_2(\text{Va}^{2+})_1(\text{O}^{2-})_1 \]

$G_{\text{Sr}_3\text{Co}_2\text{O}_7} = -29398.0 \text{kJ mol}^{-1}$  

Sr$_3$Fe$_{2-x}$Co$_x$O$_{7-8}$  
\[ (\text{Sr}^{2+})_3(\text{Fe}^{3+})_2(\text{Co}^{3+})_x(\text{Fe}^{4+})_2(\text{O}^{2-})_2(\text{Va}^{2+})_1(\text{O}^{2-})_1 \]

$G_{\text{Sr}_3\text{Fe}_{2-x}\text{Co}_x\text{O}_{7-8}} = -165000 \text{kJ mol}^{-1}$  

Sr$_4$Fe$_{6-x}$Co$_x$O$_{13}$  
\[ (\text{Sr}^{2+})_4(\text{Fe}^{3+})_6(\text{Co}^{3+})_x(\text{Fe}^{2+})_2(\text{Fe}^{4+})_2(\text{O}^{2-})_2(\text{Va}^{2+})_2(\text{O}^{2-})_2 \]

$G_{\text{Sr}_4\text{Fe}_{6-x}\text{Co}_x\text{O}_{13}} = -100000 \text{kJ mol}^{-1}$  

Sr$_6$Co$_5$O$_{15}$  
\[ (\text{Sr}^{2+})_6(\text{Co}^{4+})_5(\text{Co}^{3+})_1(\text{O}^{2-})_5 \]

$G_{\text{Sr}_6\text{Co}_5\text{O}_{15}} = -5899515.9 \text{kJ mol}^{-1}$

Functions

$GS2OV = GSROSOL + GCOOS + 28889.74 - 15.20777T$

$GS3OV = GSROSOL + GCOOS + 0.5GHSEROO - 20754.002 + 10.997T$

$GS4O = GSROSOL + GCOOS + GHSEROO - 86550.2 + 75.64357T$
\[
GCOOS = -252530 + 270.075T - 47.825T \ln(T) - 0.005112T^2 + 225008T^{-1}
\]

\[
GSROSL = -607870 + 268.9T - 47.56T \ln(T) - 0.00307T^2 + 190000T^{-1}
\]

\[
GSR2CO2O5 = -1864898.9 + 1216.424 - 213.734T \ln(T) - 0.023799T^2
\]
\[+ 1635410.1T^{-1}\]

\[
GSR3FE2O7 = 2GSRPRV + GSROSOL
\]

\[
GSRPRV = GSROSOL + 0.5GFE2O3 - 44701 - 8.73T
\]

\[
GHSERCO = +310.241 + 133.36601T - 25.0861T \ln(T)
\]
\[- 0.002654739T^2 - 1.7348 \times 10^{-7}T^3 + 72527T^{-1}\]
\[298.15 < T < 1768\]
\[\leq 17197.666 + 253.28374T - 40.5T \ln(T) + 9.3488 \times 10^{30}T^{-9}\]
\[1768 < T < 6000\]

\[
GHSEROO = -3480.87226 - 25.5028601T - 11.1355068T \ln(T)
\]
\[- 0.005098873T^2 + 6.6184604 \times 10^{-7}T^3 - 38364.8742T^{-1}\]
\[298.15 < T < 1000\]
\[- 6568.76015 + 12.6600017T - 16.8138015T \ln(T)\]
\[- 5.9579637 \times 10^{-4}T^2 + 6.78056555 \times 10^{-9}T^3 + 262904.778T^{-1}\]
\[1000 < T < 3300\]
\[- 13986.728 + 31.259625T - 18.9536T \ln(T) - 4.25243 \times 10^{-4}T^2\]
\[+ 1.0721 \times 10^{-8}T^3 + 4383200T^{-1}\]
\[3300 < T < 6000\]

\[a\] All parameters are in SI units: J, mol, K and Pa.

\[b\] The Gibbs energy terms for the end-members belonging to Sr-Fe-O can be found in Reference [8].

3.7. Optimization

Table 2 lists a set of optimized thermodynamic parameters obtained in this research for the ternary oxide phases in Sr-Co-O and for the quaternary solid solution phases in Sr-Co-Fe-O. The evaluation of the model parameters was carried out using recurrent runs of the PARROT program [64] in the Thermo-Calc software, which works by minimizing the square sum of the differences between experimental values and computed ones. In the optimization, each piece of experimental information is given a certain weight by assessing the quality of experiments. The weights were adjusted during the assessment until most of the experimental data were accounted for within the uncertainty limits stated.

In this research, for the perovskite phase, \( A_1 \) in Eq.8 was optimized with enthalpy of formation from Calle-Vallejo et al. [23], while \( B_1 \) was obtained by assuming that the entropy of \( \text{SrCoO}_3 \) at 298 K is similar to that of \( \text{Sr}_2\text{Co}_2\text{O}_5 \), and making further adjustments using phase diagram data. \( A_2 \) and \( B_2 \) in Eq.9, \( A_3 \) and \( B_3 \) in Eq. 10, and
one interaction parameter $l_{\text{perovskite}}^{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}}$ were optimized using phase diagram data and oxygen non-stoichiometry data.

As mentioned in Section 3.3, parameters for Sr$_2$Co$_2$O$_5$ and Sr$_6$Co$_5$O$_{15}$ were further adjusted using phase diagram data from Takeda et al. and Vashook et al. [13, 16], and, with lower weight in the optimization, enthalpies of formation from first-principles analysis [23]. While tremendous progress on accuracy of first-principle results has been achieved for metallic phases, the same quality does not apply to complex oxide phases due to many free parameters associated with competition between various interactions on the atomistic level. Thus, we expect higher uncertainty of first principle calculations compared to experiments. Nonetheless, deviation of assessed thermodynamic data from first principles results [23] was minimal. For Sr$_3$Co$_2$O$_7$, $A_0$ and $B_0$ in Eq.5 were adjusted such as to make the phase stable at 1373 K in air in accordance with the findings from Aksenova et al. and Cherepanov et al. [34, 65].

For the Sr-Co-Fe-O system, three interaction parameters $l_{\text{perovskite}}^{\text{Sr}^{2+}:\text{Co}^{3+}:\text{Fe}^{3+}:\text{O}^{2-}}$ for the perovskite phase were optimized using phase diagram and oxygen non-stoichiometry data. In addition, interaction parameters for Sr$_3$Fe$_2$Co$_x$O$_{7-\delta}$ ($l_{\text{Sr}^{2+}:\text{Sr}^{3+}:\text{O}^{2-}:\text{Fe}^{3+}:\text{Co}^{3+}:\text{O}^{2-}}^{\text{Sr}^{2+}:\text{Sr}^{3+}:\text{Fe}^{3+}:\text{Co}^{3+}:\text{O}^{2-}}$) and Gibbs energy terms for the end-members of Sr$_x$Fe$_{2-x}$Co$_2$O$_{13}$ ($G_{\text{Sr}^{2+}:\text{Fe}^{3+}:\text{Co}^{3+}:\text{Fe}^{3+}:\text{O}^{2-}}^{\text{Sr}^{2+}:\text{Fe}^{3+}:\text{Co}^{3+}:\text{Fe}^{3+}:\text{O}^{2-}}$) were optimized to account for the Co solubility in Sr$_3$Fe$_2$O$_7$ and Sr$_4$Fe$_6$O$_{13}$, respectively. The selection of interaction parameters was based on availability of experimental data. The parameters having direct control or the largest influence on the experimental data were selected for optimization. In general, over-optimization should be prevented, i.e. use as less interaction parameters as possible.

4. Results and discussion

4.1. The Sr-Co-O system

The enthalpy of formation, entropy, and Gibbs energy of formation calculated in this research for Sr$_6$Co$_5$O$_{15}$, Sr$_2$Co$_2$O$_5$ and SrCoO$_{3-\delta}$ at 298 K are listed in Table 3, together with the values from the literature. The CALPHAD results deviate in general from those obtained via first principle calculations. For further refinement of the Gibbs energy functions, complementary thermodynamic measurements would be advantageous.

Table 3 Enthalpy of formation, entropy, and Gibbs energy of formation for Sr$_6$Co$_5$O$_{15}$, Sr$_2$Co$_2$O$_5$ and SrCoO$_{3-\delta}$ at 298 K

<table>
<thead>
<tr>
<th>Phase</th>
<th>Results</th>
<th>Method</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>Sr$_6$Co$<em>5$O$</em>{15}$</td>
<td>$\Delta U_{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}}^{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}} (298 K) = -206.899$ kJ/mol</td>
<td>First principle</td>
<td>[25, 30]</td>
</tr>
<tr>
<td></td>
<td>$\Delta U_{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}}^{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}} (298 K) = -206.899$ kJ/mol</td>
<td>CALPHAD</td>
<td>[25]</td>
</tr>
<tr>
<td></td>
<td>$\Delta U_{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}}^{\text{Sr}^{2+}:\text{Co}^{3+}:\text{O}^{2-}} (298 K) = -218.437$ kJ/mol</td>
<td>CALPHAD</td>
<td>This work</td>
</tr>
</tbody>
</table>
The calculated phase diagrams for SrO$_x$-CoO$_y$ from this work are shown in Figs. 1–3. Fig. 1 presents the calculated SrO$_x$-CoO$_y$ phase diagram in air. Phase relations are in particular delicate inside the composition range of $x(\text{Co})/(x(\text{Co}) + x(\text{Sr})) = 0.4 - 0.5$ and the temperature range of $T = 1100 - 1500$ K. In this relatively small region, 5 two-phase equilibria are predicted by the modeling. In addition to Sr$_6$Co$_5$O$_{15}$, Wong-Ng et al. [33] reported the existence of Sr$_4$Co$_3$O$_9$ and Sr$_5$Co$_4$O$_{12}$ at 1123 K in air. According to the current modeling, however, Sr$_4$Co$_3$O$_9$ and Sr$_5$Co$_4$O$_{12}$ phases cannot co-exist with Sr$_6$Co$_5$O$_{15}$ at 1123 K in air. These two phases were therefore excluded in the final database. The calculated stable temperature range for Sr$_3$Co$_2$O$_7$-$\delta$ in air is 1246 – 1610 K. Fig. 2 shows the calculated isothermal PO$_2$-composition (PO$_2$ = 10$^{-14}$ – 1 bar) phase diagrams at temperatures 1073 – 1373 K. The stable phases at different conditions can be read directly from the figures.
Fig. 1. Calculated phase diagram of $\text{SrO}_x\text{-CoO}_y$ in air based on the parameters obtained in this research.
Fig. 2. Calculated isothermal $P_\text{O}_2$-composition phase diagrams at different temperatures based on the parameters obtained in this research. (a) 1373 K, (b) 1273 K, (c) 1173 K, (d) 1073 K. FCC is the metallic Co phase.

SrCoO$_3$-$\delta$ is one of the most important oxides in Sr-Co-O and its stability depends strongly on temperature and oxygen partial pressure. Fig. 3 shows the calculated stability phase diagram for a composition of $x(\text{Co})/(x(\text{Co})+x(\text{Sr})) = 0.5$ in comparison with experimental results. The single phase or two-phase region is labeled with phase names and the line between two neighboring regions represents a univariant reaction. As shown in Fig. 3, the deviation between the calculations and the experimental results is within a reasonable range. The SrCoO$_3$-$\delta$ perovskite phase is stable only at high temperature. In a temperature range of 1250−1423 K, SrCoO$_3$-$\delta$ is stable at $P_\text{O}_2=10^{-0.5}−10^{-4.8}$ bar, and its stability decreases with decreasing temperature. The type of diagram shown in Fig. 3 can be used to explore the stability range (temperature and oxygen partial pressure) for required phases under a given composition.
Fig. 3. Calculated stability diagram of Sr-Co-O based on the parameters obtained in this research for a composition of \(x\text{Co}/(x\text{Co} + x\text{Sr})=0.5\) with experimental data from Cherepanov et al. [65], Vashook et al. [16] and Takeda et al. [13] included. The symbols represent single-phase or two-phase region or phase boundary between two neighboring regions, e.g. \(\text{Sr}_2\text{Co}_2\text{O}_5/\text{Sr}_6\text{Co}_5\text{O}_{15}\).

Fig. 4 plots site fractions in SrCoO\(_{3-\delta}\) in air and at \(P\text{O}_2 = 1\) Pa calculated in this research. At low temperatures where SrCoO\(_{3-\delta}\) is metastable, Co exists mainly as Co\(^{4+}\) in SrCoO\(_{3-\delta}\). At high temperatures, SrCoO\(_{3-\delta}\) has a large oxygen vacancy at the O site where Co\(^{3+}\) and Co\(^{2+}\) form. The results presented in Figure 4 should be treated with caution, as no validation is conducted due to lack of experimental data on ion distribution and charge disproportion in SrCoO\(_{3-\delta}\) at different temperatures and \(P\text{O}_2\). Fig. 5 shows the oxygen nonstoichiometry in SrCoO\(_x\) calculated in this research, which agrees very well with the experimental results.
Fig. 4. Site fractions in SrCoO$_{3-\delta}$. (a) in air (b) at $P_{O_2} = 1$ Pa.

Fig. 5. Calculated oxygen nonstoichiometry ($x$ in SrCo$_x$O$_3$) as a function of log$P_{O_2}$ in comparison with experimental data from Vashook et al. [17].

4.2. The Sr-Co-Fe-O system

Fig. 6 presents the calculated SrO-Fe$_2$O$_3$-CoO$_x$ isothermal sections at 1373 K and 1173 K in air. The calculated isothermal section at 1373 K agree quite well with most of the experiment data from Aksenova et al. [34], except for the Co solubility in Sr$_3$Fe$_{2-y}$Co$_y$O$_7$. As presented in the isothermal section at 1173 K in Fig. 6b, the phase equilibria at the SrO-CoO side is quite different from 1373 K (Fig. 6a), as disappearing of Sr$_2$Co$_2$O$_5$, Sr$_3$Co$_2$O$_7$, and forming Sr$_6$Co$_5$O$_{15}$. At both temperatures, the SrCo$_{1-x}$Fe$_x$O$_{3-\delta}$ perovskite phase is more stable on
the Fe-rich side (i.e. close to SrFeO$_{3-\delta}$) and forms secondary phases on the Co-rich side. Experimental studies on phase equilibria in SrO-Fe$_2$O$_3$-CoO$_x$ at other temperatures under varying oxygen partial pressures would thus be valuable to validate the calculations and the database of Sr-Co-Fe-O. Fig. 7 plots the phase fraction for the composition SrFe$_{1-x}$Co$_x$O$_3$ at 1273 K and 973 K in air. According to the calculation, the SrFe$_{1-x}$Co$_x$O$_3$ perovskite is stable when $x < 0.64$ at 1273 K, and when $x < 0.35$ at 973 K. Table 4 lists the calculated Co solubility in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13}$ at different temperatures in air, as compared with experimental results. It clearly shows that the Co solubility in Sr$_4$Fe$_{6-x}$Co$_x$O$_{13}$ decreases with decreasing temperature.

![Diagram](image1)

(a)

![Diagram](image2)

(b)
Fig. 6. Calculated isothermal sections of SrO-Fe₂O₃-CoOₓ in air compared with the experimental data from Aksenova et al. [34]. a) at 1373 K, b) at 1173 K.

Fig. 7. Calculated phase fraction for the composition SrFe₁₋ₓCoₓO₃₋δ in air, (a) at 1273 K and (b) 973 K.

Table 4 Co solubility (x in SrₓFe₁₋ₓCoₓO₃₋δ) at different temperatures in air

<table>
<thead>
<tr>
<th>T</th>
<th>Calculated solubility from this research (x)</th>
<th>Experimentally determined solubility from the literature</th>
</tr>
</thead>
<tbody>
<tr>
<td>1123</td>
<td>1.1</td>
<td>1</td>
</tr>
<tr>
<td>1173</td>
<td>1.2</td>
<td>1.5</td>
</tr>
<tr>
<td>1223</td>
<td>1.3</td>
<td>2</td>
</tr>
<tr>
<td>1273</td>
<td>1.4</td>
<td>1.5</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.8</td>
</tr>
<tr>
<td>1373</td>
<td>1.6</td>
<td>1.6</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1</td>
</tr>
<tr>
<td>1423</td>
<td>1.6</td>
<td>1.4</td>
</tr>
<tr>
<td>1473</td>
<td>1.7</td>
<td>1.5</td>
</tr>
</tbody>
</table>

As for the SrFe₁₋ₓCoₓO₃ perovskite, the cation valences (especially the B-site cations) have a significant influence on the magnetic property, electronic conductivity and thermal-conductivity of the perovskite. Fig. 8 presents the site fractions in SrCo₀.₂Fe₀.₈O₃₋δ in air as a function of temperature calculated in this research. Co and Fe exist mainly as cations with a valency of 4+ or 3+ in SrCo₀.₂Fe₀.₈O₃₋δ in air, which is in agreement with
experimental results [18, 20]. With increasing temperature, the concentration of Co$^{3+}$ and of Fe$^{3+}$ increases at the expense of a decrease in the concentration of Co$^{4+}$ and Fe$^{4+}$.

Fig. 8. Calculated site fractions in SrCo$_{0.2}$Fe$_{0.8}$O$_{3-\delta}$ in air as a function of temperature.

5. Conclusions

As a part of a research project aimed at developing a thermodynamic database of the La-Sr-Fe-Co-O system for applications in SOFCs and gas permeation membranes, thermodynamic databases for the Sr-Co-O and Sr-Co-Fe-O systems were developed by applying the CALPHAD method. The resulting database can be used for calculating phase equilibria and thermodynamic properties at temperatures of 298 – 3000 K and oxygen partial pressure of $10^{-15}$ – 1 bar using Gibbs energy minimization software. Despite scarce experimental data for some of the occurring phases, the assessed thermodynamics of the perovskite phase allows for some important predictions: In Sr-Co-O, the perovskite phase is stable only at high temperatures; In Sr-Co-Fe-O, at $T \leq 1373$ K, the perovskite phase is stable with high Fe content and its stability decreases with decreasing temperature.

Acknowledgment

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References