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Gas-liquid two-phase flow behavior in terrain-inclined pipelines for gathering transport system of wet natural gas

Yan Yang¹, Jingbo Li¹, Shuli Wang¹, Chuang Wen².*

¹School of Petroleum Engineering, Changzhou University, Zhonglou District,
Changzhou, 213016, China

²Department of Mechanical Engineering, Technical University of Denmark, Nils Koppels Allé, 2800 Kgs. Lyngby, Denmark

*Corresponding author. Tel.: +45 45254168; fax: +45 45251961. E-mail address: cwen@mek.dtu.dk (C. Wen).

Abstract: The Volume of Fluid method and Re-Normalisation Group (RNG) k-ε turbulence model were employed to predict the gas-liquid two-phase flow in a terrain-inclined pipeline with deposited liquids. The simulation was carried out in a 22.5 m terrain-inclined pipeline with a 150 mm internal diameter. The flow parameters were numerically analyzed in detail including the phase distribution in pipes, the velocity and pressure around the elbow, the liquid flow rate and liquid holdup in different cross-section and the volume of liquid outflow. The numerical results presented that a wave crest formed on the liquid level under the suction force which caused by the negative pressure around the elbow, and then it touched to the top of the pipe. When the liquid blocked the pipe, the pressure drop between the upstream and downstream of the elbow increased with the increase of the gas velocity. At larger gas velocity, more liquid was carried out to the pipeline. The liquid periodically flowed and returned along the uphill section when the liquid was no longer flowing out of the pipeline.

Keywords: terrain-inclined pipeline; natural gas; liquid loading; gas-liquid two-phase flow; Volume of Fluid
1. Introduction

During the development of condensate natural gas fields, such as the Puguang Gas Field in China’s Sichuan Basin, there are many uphill and downhill pipes due to the large rolling terrain. The liquid often accumulates in the low section of the pipeline during the wet natural gas transportation. The deposited liquid may block the pipeline under the action of the gaseous phase and cause the sharp fluctuation of pressure [1]. It affects the normal operation of the pipeline, even damages the conveying equipment. In addition, the pipeline is more easily subjected to the corrosion issues [2, 3], because the deposited liquid is difficult to be removed. Therefore, it is significant to study the flow behavior of the deposited liquid in the wet natural gas pipelines.

The gas-liquid two-phase flows in the horizontal and inclined pipes have been studied for decades. Taitel et al. [4] proposed a model to calculate the flow behavior under the transient flow conditions in a hilly-terrain pipeline system. The model was applied to the low liquid and gas flow rates, where the frictional pressure losses could be neglected. Grolman and Fortuin [5] modified the apparent rough surface model to predict the pressure gradient and liquid holdup in slightly inclined pipes. Salhi, et al. [6] improved the one-dimensional two-fluid model to study the stability of stratified gas–liquid two-phase flows in an inclined pipe. Goldstein et al. [7] proposed the exact solutions for the laminar stratified flows in the inclined pipes. The solution could be used for investigating the influence of the pipe inclination and flow geometry on the liquid holdup and pressure gradient. Gawas et al. [8] proposed a new wave celerity correlation for the gas-liquid two-phase stratified flow using the low viscosity fluids and compared with a mechanistic model proposed by Watson. The Poisson probability theory was employed to predict the slug frequency in the gas-liquid horizontal pipelines [9]. The results indicated that the theoretical model resulted in a great improvement...
with the average error of 15% for 1.2 m/s and 0.4 m/s superficial liquid velocities. Ferrari et al. [10] developed a novel two-fluid model to capture the slug flow in pipes using a one-dimensional transient hyperbolic five-equation. This theoretical model successfully captured the slug onset, growing, and development from a stratified flow in horizontal pipes.

Barnea and Shoham [11] carried out an experimental study in an air-water system with 2.55 cm and 1.95 cm pipes and compared with the theoretical prediction of Taitel and Dukler. The results showed that the experiment data agreed with the theoretical prediction. Tzotzi and Andritsos [12] gave a modified form of Andritsos-Hanratty model to study stratified gas–liquid two-phase flow and estimate pressure drop and interfacial friction factor in a horizontal pipe. Jia et al. [13] used two absolute pressure sensors to measure differential pressure and obtained the void fraction from differential pressure model in a horizontal gas-liquid two-phase pipe. Arunkumar et al. [14] used the dielectric sensors to identify the gas-liquid two-phase flow regime. The bubbly, slug and stratified flow regimes could be identified in their experiments. Abdulkadir et al. [15] investigated the unsteady hydrodynamic behaviour of slug flow in a horizontal pipe using the an air–silicone oil mixture as the working fluid. Their results showed that the drift velocity component of the bubbles in horizontal pipes far exceeds the value found for vertical riser pipes. Dinaryanto et al. [16] experimentally studied the initiation and flow development mechanisms of the gas-liquid two-phase slug flow in a horizontal pipe. The initiation frequency of slug flow and the evolution of passing slug frequency along the pipeline were also observed by using two high video cameras in their experiments. Bouyahiaoui et al. [17] employed intermediary of the conductance probe technique to study the slug flow behavior in the vertical downward air-water two-phase
pipelines. The experimental results showed that the mean averaged void fraction increased with the gas superficial velocity.

The numerical simulation approach was employed to predict the flow regime and flow characteristics of gas-liquid two-phase flows. Mouza et al. [18] used computational fluid dynamics (CFD) modeling to simulate the wave stratified two-phase flow in a horizontal pipe. The distribution of the shear stresses and the velocity profiles of both phases were calculated by using CFD modeling. Loilier [19] applied the two-fluid model to simulate the gas-liquid two-phase flow in vertical and terrain-inclined pipelines consisting of the uphill, downhill and horizontal sections. The vertical bubble flow, stratified and terrain-inclined slug flows were simulated and the avoid fraction, slug frequency, slug length and other parameters were obtained in this work. Ekambara et al. [20] performed a numerical study of the bubbly two-phase flow in a horizontal pipeline. The numerical results agreed well with the experimental data. Vallée et al. [21] investigated the stratified and slug flow by using the CFD modeling, particle image velocimetry and high-speed video observation experiments. The main flow characteristics obtained by the CFD modeling were successfully validated against the experiments. Verdin et al. [22] used the CFD modeling for simulating the transport behavior of water droplets in 38 in. diameter pipes and compared the oil flow behavior and droplets with that of water. Santim et al. [23] compared different methodologies in transient isothermal gas-liquid two-phase slug flows in a horizontal pipeline. The comparison studies indicated that the Drift-Flux Model presented the better agreement with the pressure wave velocity by the experiments. Wang et al. [24] developed a fluid-structure interaction dynamic model for the conveying severe slugging flow in a pipeline-riser system. Their results showed that the dynamic response of the riser is
closely related to the characteristics of severe slugging, which can be used to eliminate severe slugging phenomenon to improve the stability of the riser system.

The gas-liquid two-phase flows in horizontal and vertical pipelines attract a great many of experimental and numerical studies, while the flow mechanism of the gathering liquids in the elbow of hilly-terrain pipes is still not understood very well. The purpose of this work is to study the gas-liquid two-phase flow in the terrain-inclined pipelines to obtain local and global flow characteristics using the CFD modeling, which can provide the detailed information for removing the deposited liquid effectively and correspondingly protecting natural gas pipeline and equipment. The Volume of Fluid (VOF) method and Re-Normalisation Group (RNG) $k$-$\varepsilon$ turbulence model are used to predict the deposited liquid motion under the action of the gaseous phase in the terrain-inclined pipeline system including an uphill section and a downhill section. The flow parameters are analyzed in detail, including the velocity and pressure distribution, phase fraction and cross-section liquid holdup.

2. Terrain-inclined pipelines

The terrain-inclined pipeline is employed for our current studies, which contains a downhill section and an uphill part. The inclination angles of downhill and uphill segments are both assigned to be $5^\circ$. It is assumed that a certain amount of liquids is gathering at the bottom of the inclined pipeline system, as shown in Fig. 1. The three-dimensional geometry of terrain-inclined pipeline is established for the computational domain, including the pipe diameter of $D=150$ mm and the length of every section is $75 \, D$, which ensures a fully developed flow for this terrain-inclined pipeline system.
The flow data for the gas-liquid two-phase is presented in Table 1, obtained from Puguang gas field in Sichuan Basin of China. The working fluids are methane and water, respectively. In this case, the solubility of methane in water can be neglected, we, therefore, assume that the water and methane are the immiscible fluids in the numerical studies.

### Table 1 Flow data for gas-liquid two-phase numerical simulation

<table>
<thead>
<tr>
<th>Case</th>
<th>$U_G$ (m s$^{-1}$)</th>
<th>$\rho_G$ (kg m$^{-3}$)</th>
<th>$\mu_G$ (m$^2$ s$^{-1}$)</th>
<th>$h/D$</th>
<th>$V_L$ (m$^3$)</th>
<th>$\rho_L$ (kg m$^{-3}$)</th>
<th>$\mu_L$ (m$^2$ s$^{-1}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>5.5</td>
<td>0.6679</td>
<td>1.087×10$^{-5}$</td>
<td>0.75</td>
<td>0.01637</td>
<td>998.2</td>
<td>0.001003</td>
</tr>
<tr>
<td>2</td>
<td>6.5</td>
<td>0.6679</td>
<td>1.087×10$^{-5}$</td>
<td>0.75</td>
<td>0.01637</td>
<td>998.2</td>
<td>0.001003</td>
</tr>
<tr>
<td>3</td>
<td>7.5</td>
<td>0.6679</td>
<td>1.087×10$^{-5}$</td>
<td>0.75</td>
<td>0.01637</td>
<td>998.2</td>
<td>0.001003</td>
</tr>
</tbody>
</table>

### 3. Computational methods

#### 3.1 Governing equations

The gas-liquid two-phase flow in the terrain-inclined pipeline represents a distinct phase interface. The interface catching is a key issue for our simulation of this kind of flow behavior. The VOF model [25] uses the surface-tracking technology based on the fixed Eulerian mesh, which can be employed to model two or more immiscible fluids. Therefore, we utilize the VOF model here to track the gas-liquid phase interface in the terrain-inclined pipelines.
The continuity equation is as follows [26]:
\[
\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_i}(\rho u_i) = 0
\]  
(1)

The momentum equation [27] is solved throughout the computational domain, and the velocity field is shared among all the phases.
\[
\frac{\partial}{\partial t} (\rho u_i) + \nabla \cdot (\rho u_i u_j) = -\nabla p + \nabla \cdot \left( \mu \left( \nabla u_i + \nabla u_i^T \right) \right) + \rho g + F
\]  
(2)

where \( \rho \) is the density, \( u \) is the velocity, \( p \) is the static pressure, \( \mu \) is the dynamic viscosity, \( \rho g \) is the gravitational body force and \( F \) is external body force.

For the gas-liquid two-phase flow, the \( \rho \) and \( \mu \) in each computational cell are given by the following equations:
\[
\rho = \alpha \rho_2 + (1 - \alpha) \rho_1
\]  
(3)
\[
\mu = \alpha \mu_2 + (1 - \alpha) \mu_1
\]  
(4)

The volume fraction of each phase in each grid cell is calculated throughout the domain. The interface between two phases is tracked by solving the continuity equation for the volume fraction of one (or more) phases. The volume fraction equation is as follows [28]:
\[
\frac{1}{\rho_1} \left[ \frac{\partial}{\partial t} (\rho_1 \alpha) + \nabla \cdot (\rho_1 \alpha u_i) \right] = \sum_{p=1}^{c} (\dot{m}_{pq} - \dot{m}_{qp})
\]  
(5)

where \( \dot{m}_{pq} \) is the mass transfer from phase \( q \) to phase \( p \) and \( \dot{m}_{qp} \) is the mass transfer from phase \( p \) to phase \( q \), \( \alpha_q \) is the volume fraction of phase \( q \).

The primary phase volume fraction is solved based on the following constraint:
\[
\sum_{q=1}^{c} \alpha_q = 1
\]  
(6)

3.2 Continuum surface force model

The effect of surface force along the interface is included in the VOF model. The continuum surface force model proposed by Brackbill et al. [29] is used in this paper.
It is implemented as a source term in the momentum equation. The pressure drop across the surface depends on the surface coefficient $\sigma$ and the surface curvature are measured by two radii in orthogonal directions $R_1, R_2$:

$$p_i - p_s = \sigma \left( \frac{1}{R_1} + \frac{1}{R_2} \right)$$  \hfill (7)

The surface force $F_{vol}$ is expressed as follows:

$$F_{vol} = \sigma \frac{\rho_k \nabla \alpha}{2 (\rho_i + \rho_s)}$$  \hfill (8)

### 3.3 Turbulence model

Among the $k-\varepsilon$, $k-\omega$ and Reynolds stress models [30-32], the RNG $k-\varepsilon$ turbulence model [33] involves an additional term to its dissipation rate equation that can improve the accuracy for rapidly strained flow. The swirling effect on the turbulence is also included in the RNG $k-\varepsilon$ model. Both of these two modifications show that the RNG $k-\varepsilon$ model is more suitable for the prediction of the large curvature and strain rate flow in the terrain-inclined pipelines. Therefore, the RNG $k-\varepsilon$ turbulence model is employed here, because the flow turns at the elbow of the pipe, which connects the uphill section and downhill section in terrain-inclined pipelines. The turbulence kinetic energy, $k$, and its rate of dissipation, $\varepsilon$, are as follows:

$$\frac{\partial (\rho k)}{\partial t} + \frac{\partial (\rho u_i k)}{\partial x_i} = \frac{\partial}{\partial x_i} \left[ \alpha_k \mu_d \frac{\partial k}{\partial x_i} \right] + G_k + G_b - \rho \varepsilon - Y_M$$  \hfill (9)

$$\frac{\partial (\rho \varepsilon)}{\partial t} + \frac{\partial (\rho u_i \varepsilon)}{\partial x_i} = \frac{\partial}{\partial x_i} \left[ \alpha_\varepsilon \mu_d \frac{\partial \varepsilon}{\partial x_i} \right] + \frac{C_{\varepsilon}}{k} (G_k + C_{\varepsilon} G_\varepsilon) - C_{\varepsilon} \rho \frac{\varepsilon^2}{k} - Y_M$$  \hfill (10)

where $G_k$ represents the generation of turbulence kinetic energy due to the mean velocity gradients, $G_b$ is the generation of turbulence kinetic energy due to buoyancy, $Y_M$ represents the contribution of the fluctuating dilatation in compressible turbulence to the overall dissipation rate. $\alpha_k$ and $\alpha_\varepsilon$ are the inverse effective Prandtl numbers for $k$ and $\varepsilon$, respectively.
3.4 Numerical schemes

The pressure-based VOF method and the RNG \( k-\varepsilon \) turbulence model are employed for the computation of the gas-liquid two-phase flow in the terrain-inclined pipeline. The VOF explicit scheme is performed for accurately tracking the surface between the gaseous and liquid phases [34]. In addition, the effect of the body force is considered, including the gravity and the surface force. The pressure-velocity coupling is the PISO algorithm. The continuity equation, momentum equation and turbulence equations are discretized with the QUICK method, while the Geometric Reconstruction scheme (Geo-Reconstruct) is chosen for the volume fraction equation.

The structured mesh is adopted as the strategy of the grid system for the whole computational domain. The velocity-inlet and outflow boundary conditions are set for the inlet and outlet of the calculation model. The residual of the calculation is set to \( 10^{-5} \), and the maximal iterations per time step is 20. The time step size is \( 10^{-4} \) s for proving the calculation convergence in every time step.

4. Results and discussion

4.1 CFD validation

The numerical method was validated with the experiments carried out by Heywood and Richardson [35] before we applied it to the gas-liquid two-phase flows in the terrain-inclined pipeline. The liquid holdup of a slug flow was measured by \( \gamma \)-ray absorption method in a horizontal pipeline with 42 mm inner diameter and 457 m length. The average results of the experimental data were estimated by the probability density functions. The working fluids were air and water, respectively. For the CFD model validation, we utilized six group experimental data in the different gas superficial velocities, while the liquid superficial velocity was fixed at 0.978 m/s. The comparison of the slug flow in a horizontal pipe was described in Fig. 2. The similar flow behavior
of liquid holdup in the slug flow was presented between the simulations and experiments. The maximum relative error was approximate 5.9% occurring at the gas superficial velocities of 4.145 m/s. Therefore, our numerical simulation agreed well with the experimental data.

Fig. 2. Liquid holdup in a slug flow

4.2 Process of liquid touched to the top of pipe

Fig. 3 shows the phase distribution near the elbow of the terrain-inclined pipe at different times with inlet gas velocity of $U_G = 5.5$ m/s. The blue contour presents the gaseous phase, while the red one reveals the liquid phase. The gas-liquid phase interface was disturbed when the gaseous phase flowed through the interface of the liquid phase. A wave crest gradually formed in the gas-liquid interface at $t = 0.005$ s – 0.085 s. Then the wave crest continued to become larger and moved along the gas flow direction at the same time ($t = 0.095$ s – 0.105 s in Fig. 3). At $t = 0.115$ s, the wave crest touched to the top of the pipe and then filled the whole cross-section of the pipe at this point.
Fig. 3. Gas and liquid phase distribution in the terrain-inclined pipes

The distribution of velocity magnitude and static pressure near the elbow are shown in Fig. 4. It represents that the gas velocity obviously increased when the gaseous phase flows through the region above the liquid, and the maximum velocity can reach 36 m/s at $t = 0.005$ s. The gas flow area decreased due to the liquids assembled at the bottom of the pipe, which caused the increase of the gas velocity when the gaseous phase went through the region. At $t = 0.085$ s, the gaseous phase flowed through the top of the wave crest and then it was accelerated. At $t = 0.115$ s, the distribution of velocity in the pipe was different from the first two time steps due to the gaseous phase was separated to two discontinuous regions by the liquid phase.

The pressure difference appeared around the elbow of the terrain-inclined pipe at $t = 0.005$ s due to the increase of the gas velocity. The gas-liquid phase interface was disturbed and then the wave crest formed due to the large suction force overcoming the effect of the gravity. At $t = 0.085$ s, the pressure difference region moved to the
downstream of the elbow and the suction force was directed to this region. The wave crest correspondingly continued to become larger and moved to the downstream of the elbow. At \( t = 0.115 \) s, some amount of the liquids filled the whole cross-section of the pipe and the gaseous phase was divided into two parts. The upstream pressure is higher than the downstream one. The liquids moved by the action of the pressure difference between the upstream and downstream.

![Diagram](image)

(a) Velocity magnitude  (b) Static Pressure

Fig. 4 Velocity and static pressure in the terrain-inclined pipes

4.3 Liquid motion along the uphill section

The gaseous and liquid phase distributions along the uphill section are shown in Fig. 5. The blue and red contours present the gaseous and liquid phase, respectively. We can see that the liquid phase blocked the pipe cross-section like a water plug after the liquid touched to the top of the pipe. It continued to move under the action of the pressure difference between the upstream and downstream, and some liquids moved along the top of the pipe and fell under the action of the gravity (\( t = 0.12 \) s - 0.40 s). At \( t = 0.80 \) s - 5.0 s, the liquid motioned to the pipe outlet and some liquids flowed out the pipe exit under the action of the inertia and the shearing force of the gaseous phase. At \( t = 6.0 \) s - 9.8 s, the remanent liquids refluxed to the elbow and the gravity played a dominant role in this fluid flow process.
Fig. 5 Gas and liquid phase distribution along the uphill pipe

Fig. 6 shows the flow rate of the liquid phase at four cross-sections of the uphill pipe \((x = 1000\, \text{mm}, 3000\, \text{mm}, 6000\, \text{mm}, 7000\, \text{mm})\) under the superficial gas velocity of \(U_G = 5.5\, \text{m/s}\). The flow rate presented the periodic alternation between the positive and negative in Fig. 6 (a) - (d). The positive flow rate means that the liquid motioned along the uphill section, while the negative flow rate represents the liquid refluxed. In Fig. 6 (c), the flow rate is equal to 0 kg/s in this cross-section \((x = 6000\, \text{mm})\) at \(t = 11\, \text{s} - 17\, \text{s}\), and then a small amount of the liquid flowed through this cross-section of the uphill pipe. It illustrates that the liquid discontinuously flowes through this pipe cross-section. We can see from the Fig. 7 (d) that there was almost no liquid going through this pipe cross-section \((x = 7000\, \text{mm})\) after \(t = 10\, \text{s}\). We, therefore, can conclude that the periodic reflux occurs between the elbow and the pipe cross-sections less than \(x = 6000\, \text{mm}\).
Fig. 6 Liquid flow rate in different cross-sections of the uphill pipe

Fig. 7 shows the liquid holdup in different cross-sections of the uphill pipe under the superficial gas velocity of $U_G = 5.5$ m/s. The fluctuation of the liquid holdup is shown in Fig. 7 (a) - (d). The liquid holdup showed the maximum peak of 1.0 in Fig. 7 (a). In other words, the liquids filled the whole cross-section of the uphill pipe at this point. As shown in Fig. 7 (b), the minimum peak of the liquid holdup was around 0.17 and the highest one was close to 0.40. In Fig. 7 (c) - (d), the maximum liquid holdup could reach 0.14 and the others were below 0.05. The average thickness of the liquid layer tended to become thinner from the uphill cross-section of $x = 1000$ mm to $x = 7000$ mm. The shearing force of the gaseous phase on the liquid was gradually weakened, and the liquids no longer flowed out the pipe exit.
The volume of the liquid was monitored in the numerical simulation and the results were presented in Fig. 8. The similar trends were obtained under three different inlet gas velocities. It can be seen that some liquids flowed out the pipe exit and the volume of liquid rapidly decreased in about 2 seconds. Later, the volume of the liquid was no longer changed. The volume percentages of the outflow liquid are approximately 20.16%, 28.4%, 31.58% under the superficial gas velocity of $U_G = 5.5 \text{ m/s}$, 6.5 m/s, 7.5 m/s, respectively.
5. Conclusions

The VOF method and the RNG $k$-$\varepsilon$ turbulence model were used to simulate the gas-liquid two-phase flows within the terrain-inclined pipelines with deposited liquids. The flow area decreased because of the liquids gathered at the bottom of the pipe. It caused the increase of the gas velocity and the decrease of the static pressure, which generated the suction force above the gas-liquid interface. A wave crest formed and touched the top of the pipe under the action of the suction force around the elbow. The liquid blocked the cross-section of the pipe and correspondingly resulted in a large pressure drop. A certain amount of liquids can flow out the pipe exit carried by the gaseous phase. The volume percentages of the outflow liquids were 20.16%, 28.4%, 31.58% under the superficial gas velocity of 5.5 m/s, 6.5 m/s, 7.5 m/s, respectively. The liquid periodically flowed and returned along the uphill section in the pipe when the liquid was no longer flowing out the pipe exit.

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References


