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VARIABLE EXPONENT CALDERÓN'S PROBLEM IN ONE DIMENSION

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Abstract. We consider one-dimensional Calderón's problem for the variable exponent $p(\cdot)$ -Laplace equation and find out that more can be seen than in the constant exponent case. The problem is to recover an unknown weight (conductivity) in the weighted $p(\cdot)$ -Laplace equation from Dirichlet and Neumann data of solutions. We give a constructive and local uniqueness proof for conductivities in L^∞ restricted to the coarsest sigma-algebra that makes the exponent $p(\cdot)$ measurable.

1. Introduction

Calderón's problem [10] asks if an electric conductivity γ in an object Ω can be reconstructed from boundary measurements of current and voltage, which are given by the Dirichlet-to-Neumann map (DN map) $u|_{\partial\Omega} \rightarrow \gamma \nabla u \cdot \nu|_{\partial\Omega}$, where ν is the unit outer normal. In one dimension, the answer to Calderón's problem is negative; only the resistance, that is, the total resistivity $\int_I \gamma^{-1} dx$, can be recovered [17, Section 1.1], where $I \subset \mathbf{R}$ is an open bounded interval. A similar result holds for p -Calderón's problem, where the forward model is given by the weighted p -Laplace equation $-\operatorname{div}(\gamma |u'|^{p-2} u') = 0$: it is only possible to recover the value of the integral

$$\int_I \gamma^{1/(1-p)} dx$$

from the DN map [4, Theorem 2.2]. This is true for any constant $1 < p < \infty$, and is also a special case of corollary 10 in this paper. We describe what can be recovered in the case of a variable exponent $p(\cdot)$. This is the first investigation of an inverse problem related to the variable exponent $p(\cdot)$ -Laplacian. The problem in the constant exponent case was introduced by Salo and Zhong [31] in 2012, after which other theoretical results have been published [3, 5, 7, 8, 18, 24], as well as a numerical study [19]. The works of Sun and others consider the problem of recovering the dependence of A in

$$\operatorname{div}(A(x, u, \nabla u) \nabla u) = 0$$

on all of x, u , and ∇u , but they either do not admit degenerate equations [29, 34] or assume the equivalent of constant p [21].

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One can think of the variable exponent conductivity equation

$$-\operatorname{div} \left(\gamma |\nabla u|^{p(x)-2} \nabla u \right) = 0$$

as arising from a non-linear Ohm's law

$$-j = \gamma |\nabla u|^{p(x)-2} \nabla u,$$

which has a power law dependence between the current j and the gradient of the electric potential u at every point, but where the exponent in the power law varies from point to point. We use this terminology and intuition in the article. An example of a power-law type Ohm's law is certain polycrystalline compounds near the transition to superconductivity [9, 15], where the exponent p is a function of temperature. However, it is not clear how relevant electrical impedance tomography of such materials would be.

1.1. Results. We use two approaches. The first is to consider the limit of the Dirichlet-to-Neumann map (hereafter DN map) as the difference of Dirichlet boundary values grows without bound or approaches zero; this gives information about the value of the conductivity at the maximum or minimum of p , when the extreme value is reached on a set of positive measure.

The second is to consider the DN map as a dual pairing in an L^2 space of essentially the conductivity and another function, and then determine what one can say about the conductivity based on this information. This gives uniqueness for the conductivity in L^∞ restricted to the coarsest sigma-algebra that makes $p(\cdot)$ measurable. We present two proofs. The simpler one requires knowing the full DN map and gives a non-constructive proof. The other proof is constructive and requires knowledge of the DN map only on an open set, but requires working with fairly explicit formulae for arbitrary order derivatives of composite functions and inverse functions. This approach is similar to a classical moment problem [32]; whereas the Hausdorff moment problem asks if there exists a measure μ such that a given sequence m_n satisfies

$$m_n = \int_I x^n d\mu(x),$$

we ask what kinds of functions f satisfy

$$m_n = \int_I (g(x))^n f(x) dx$$

for a specific sequence m_n arising from the forward problem, and for functions $g(x) = 1/(p(x)-1)$ that are derived from the variable exponent p . The results are very similar to what is achieved for one dimensional inverse source problem under attenuation and using multiple frequencies [6]. The numerical methods used there could also be applied to the problem discussed in this paper.

In the case of nonconstant p , even though we only choose our input data from an essentially one-dimensional space (difference of the two Dirichlet data points is a real number), we can recover information on an infinite-dimensional space, if the power p is not piecewise constant. Here the changing non-linear nature of the forward problem makes the inverse problem easier. Nonlinearity of the forward problem has also been used for advantage in the study of non-linear hyperbolic equations [12, 26].

Our main theorems follow. First, let I be equipped with the Lebesgue measure on the Lebesgue sigma-algebra. Define $\sigma(p)$ to be the sigma-algebra on I generated by sets of the form $p^{-1}(A)$ where $A \subseteq \mathbf{R}$ is a Borel set in \mathbf{R} . We consider p to be a function, even though we may write $p \in L^r(I)$. The spaces $L^r(I, \sigma(p))$

are the Lebesgue spaces of almost everywhere equal $\sigma(p)$ -measurable functions of finite L^r -norm with respect to the Lebesgue measure restricted to $\sigma(p)$. It is important to emphasize that $L^r(I, \sigma(p))$ is not a true subspace of $L^r(I)$ because $\sigma(p)$ does not contain all null-sets of the Lebesgue sigma-algebra on I . However, the map $L^r(I, \sigma(p)) \rightarrow L^r(I): [g]_{\sigma(p)} \mapsto [g]$ is a well-defined (possibly not surjective) isometry, so we may regard $L^r(I, \sigma(p))$ as a complete and therefore closed subspace, but emphasize that the equivalence classes are different. This subspace is characterized by the property that every equivalence class has a $\sigma(p)$ -measurable representative. Finally, we define $P: L^2(I) \rightarrow L^2(I, \sigma(p))$ as the orthogonal projection (conditional expectation in probabilistic terms) onto the closed subspace identified with $L^2(I, \sigma(p))$.

For more on sigma-algebras generated by functions or sets we refer to the book of Dellacherie and Meyer [13, definition 5].

Theorem 1. *Consider an open bounded interval I . Suppose the conductivity $\gamma \in L^{\infty}_+(I) \subset L^2(I)$, and suppose there exist constants p^{\pm} such that almost everywhere $1 < p^- < p(x) < p^+ < \infty$. Then the nonlinear projection*

$$\tilde{P}: L^2(I) \cap L^{\infty}_+(I) \rightarrow L^2(I, \sigma(p)) \cap L^{\infty}_+(I, \sigma(p))$$

defined by

$$\tilde{P}(\gamma)(x) = \left(P \left(\gamma^{-1/(p-1)} \right) (x) \right)^{-(p(x)-1)}$$

can be reconstructed from the Dirichlet-to-Neumann map. Furthermore, there exists an arbitrarily small open set of Dirichlet boundary values which are sufficient for reconstructing the projection.

Corollary 2. *Suppose $\gamma_1, \gamma_2 \in L^{\infty}_+(I)$. If $p: I \rightarrow \mathbf{R}$ is a measurable injection and $\tilde{P}(\gamma_1) = \tilde{P}(\gamma_2)$, then $\gamma_1 = \gamma_2$.*

Proof. In this case the projection P is the identity map and the powers in \tilde{P} cancel, so \tilde{P} is also the identity map. □

Remark 3. The following properties hold:

- If γ is $\sigma(p)$ -measurable, then $\tilde{P}(\gamma) = \gamma$.
- We have $\tilde{P} \circ \tilde{P} = \tilde{P}$.
- If $p \equiv 2$ and $|I| = 1$, then \tilde{P} is the harmonic mean.
- If p is constant, then $\tilde{P}(\gamma)$ is a type of average:

$$\tilde{P}(\gamma) = \left(\frac{1}{|I|} \int_I \gamma^{-1/(p-1)} dx \right)^{-(p-1)}.$$

Proof. If γ is $\sigma(p)$ -measurable, then so is $\gamma^{-1/(p-1)}$. Hence, the projection $P \left(\gamma^{-1/(p-1)} \right) = \gamma^{-1/(p-1)}$ and the exponents inside and outside the projection cancel. This proves the first point.

To prove the middle point, one observes that in $\tilde{P} \circ \tilde{P}$ the powers between the two projections P cancel, after which $P \circ P = P$, since P is a projection.

If $p \equiv 2$, then inside the projection P there is γ^{-1} . The sigma-algebra generated by the constant function $p \equiv 2$ is the trivial one; that is, $\sigma(p) = \{\emptyset, I\}$. The only functions measurable with respect to this sigma-algebra are constant functions, whence $P \left(\gamma^{-1} \right)$ must be a constant function; call the constant C . Since integrals

over $\sigma(p)$ -measurable sets are conserved by the projection, we have

$$\int_I C \, dx = \int_I \gamma^{-1} \, dx \iff C = \frac{1}{|I|} \int_I \gamma^{-1} \, dx.$$

Raising this to the power $-(p-1) \equiv -1$ gives the third claim, and the proof of the fourth claim is similar. \square

This theorem is proven in a non-constructive way in Section 3.2 using the multiplicative system theorem (Theorem 17). It turns out that the finite linear combinations of functions $K^{p(x)/(p(x)-1)}$ indexed by the real numbers $K \geq 0$ are dense in the space of functions that are both $\sigma(p)$ -measurable and square integrable. A constructive and local proof is developed in Section 3.3. The constructive proof requires explicit formulae for higher order derivatives of composite functions (Faà di Bruno's formula [22]) and inverse functions [23]. Aside from the difficulty of calculating the derivatives, the proof is quite similar to a related proof of Brander, Ilmavirta and Tyni [6].

The next theorem is proven in Section 3.1. The proof is constructive, easy to implement numerically and reasonably elementary.

Theorem 4. *Consider an open bounded interval I . Suppose the conductivity $\gamma \in L^\infty(I)$ is almost everywhere bounded from below by a positive constant. Suppose further that the exponent $p(\cdot)$ reaches its essential minimum p^- on a set of positive measure Q_+ and essential maximum p^+ on a set of positive measure Q_- . Then the DN map Λ_γ , suitably scaled, gives*

$$\begin{aligned} \lim_{m \rightarrow \infty} \overline{K}_m^{-p^-/(p^- - 1)} \frac{1}{|Q_+|} \Lambda_\gamma(m) &= \left(\int_{Q_+} (\gamma(x))^{-1/(p(x)-1)} \, dx \right)^{-(p^- - 1)} \quad \text{and} \\ \lim_{m \rightarrow 0} \overline{K}_m^{-p^+/(p^+ - 1)} \frac{1}{|Q_-|} \Lambda_\gamma(m) &= \left(\int_{Q_-} (\gamma(x))^{-1/(p(x)-1)} \, dx \right)^{-(p^+ - 1)}, \end{aligned}$$

where m is the difference between the Dirichlet boundary values and \overline{K}_m can be computed from knowledge of m and $p(\cdot)$.

This theorem allows the recovery of the average of the conductivity to a known power over the set where the exponent is largest or smallest. The situation is far more delicate if the exponent does not reach its maximum/minimum or does so in a set of zero measure. The behaviour seems to depend on $|\{x \in I; p(x) = a\}|$ as $a \rightarrow p^\pm$. We do not investigate the matter in more detail in this article.

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2. Forward problem

In this section we discuss the existence and uniqueness for the forward problem in general dimension, and define the voltage-to-current or DN map, also in general dimension. We specialize to the one-dimensional inverse problem in the following sections. Before proceeding, we define the variable exponent Lebesgue space $L^p(\Omega)$,

¹<https://mathoverflow.net/a/292978/>

with $\Omega \subset \mathbf{R}^d$ a bounded open set and $d \geq 1$. The variable exponent Sobolev spaces are defined in terms of it in the usual way. We assume throughout that $p: \Omega \rightarrow [1, \infty]$ is measurable and bounded away from one and infinity. Then, following the book of Diening, Harjulehto, Hästö and Růžička [14, Sections 2 and 3],

$$L^p(\Omega) = \left\{ f: \Omega \rightarrow \mathbf{R} \text{ measurable; } \lim_{\lambda \rightarrow 0} \int_{\Omega} \frac{1}{p(x)} |\lambda f(x)|^{p(x)} \, dx = 0 \right\},$$

where functions which agree almost everywhere are considered identical, and

$$\|f\|_{L^p(\Omega)} = \inf \left\{ \lambda > 0; \int_{\Omega} \frac{1}{p(x)} \left| \frac{f(x)}{\lambda} \right|^{p(x)} \, dx \leq 1 \right\}.$$

Consider a weight function, or conductivity,

$$\gamma \in L^{\infty}_+(\Omega) = \{f \in L^{\infty}(\Omega); \text{ess inf } f > 0\}.$$

The Dirichlet problem for the weighted variable exponent $p(\cdot)$ -Laplacian is

$$\text{div} \left(\gamma |\nabla u|^{p(x)-2} \nabla u \right) = 0 \text{ in } \Omega, \quad u = f \text{ on } \partial\Omega.$$

Uniqueness and existence for the variable exponent function has been investigated in variable exponent Sobolev spaces [14], though one often considers the equation

$$\text{div} \left(p(x) |\nabla u|^{p(x)-2} \nabla u \right) = 0$$

as the basic example [20, Section 2]. That equation arises from minimizing the functional

$$(1) \quad u \mapsto \int_{\Omega} |\nabla u|^{p(x)} \, dx,$$

while we prefer to work with the energy

$$(2) \quad u \mapsto \int_{\Omega} \frac{\gamma}{p(x)} |\nabla u|^{p(x)} \, dx.$$

Since the function $x \mapsto \frac{1}{p} |x|^p$ is the convex conjugate (by the Legendre–Fenchel transform) of $x \mapsto \frac{1}{q} |x|^q$, where q is the Hölder conjugate of p , we see the division by $p(x)$ as natural.

Calderón’s problem asks one to recover the conductivity γ from the DN map Λ_{γ} , which, in the strong form, is given by the formula

$$\Lambda_{\gamma}(f) = \gamma |\nabla u|^{p-2} \nabla u \cdot \nu|_{\partial\Omega},$$

where $u|_{\partial\Omega} = f$, i.e. the input is the potential or Dirichlet boundary value and the output is the current flowing out of the domain. This definition may fail for irregular solutions u or domains Ω . One typically uses the weak DN map instead, which we derive next. We remark that regardless of the energy/equation we use, the energy of the equation and its weak DN map are different. By formally integrating by parts, starting from the strong definition of DN map and multiplying by a test function v

with $v|_{\partial\Omega} = g$, we get

$$\begin{aligned} & \int_{\partial\Omega} \gamma |\nabla u|^{p(x)-2} \nabla u \cdot \nu g \, dS(x) \\ &= \int_{\Omega} \operatorname{div} \left(\gamma |\nabla u|^{p(x)-2} \nabla u \right) v \, dx + \int_{\Omega} \gamma |\nabla u|^{p(x)-2} \nabla u \cdot \nabla v \, dx \\ &= \int_{\Omega} \gamma |\nabla u|^{p(x)-2} \nabla u \cdot \nabla v \, dx. \end{aligned}$$

If we choose $g = f$ and $v = u$, then we have

$$\langle \Lambda_{\gamma}(f), f \rangle = \int_{\Omega} \gamma |\nabla u|^{p(x)} \, dx.$$

We take this “quadratic” form as the definition of the DN map. It can be defined as a functional on $W^{1,p}(\Omega)/W_0^{1,p}(\Omega)$, where $W_0^{1,p}$ is the closure of the space of $W^{1,p}$ functions with compact support [14, Section 8.1].

Lemma 5. *Suppose $1 < p^- \leq p(x) \leq p^+ < \infty$ almost everywhere, and that $\Omega \subset \mathbf{R}^d$, $d \in \mathbf{Z}_+$, is a bounded open set that supports the Poincaré inequality. Consider boundary values $f \in W^{1,p(\cdot)} \cap L^\infty(\Omega)$. Then there exist unique minimizers in $W^{1,p(\cdot)} \cap L^\infty(\Omega) + f$ to the energies (2) and (1).*

Proof. The proof follows the direct method in calculus of variations. The variable exponent Sobolev space is a reflexive Banach space [14, Theorem 8.1.6] and the functional is convex, since $t \mapsto ct^p$ is convex for all $p \geq 1$ and $c \geq 0$. The energies are lower semicontinuous [14, Theorem 3.2.9 and Section 3.2]. Coercivity of the functional requires the Poincaré inequality with $p \equiv 1$ (since we assume bounded boundary values). Therefore the functionals have unique minimizers. \square

This lemma holds for any bounded, open interval in \mathbf{R} . More generally, the 1-Poincaré inequality is satisfied for example in John domains [14, Section 8.2], and in particular in Lipschitz domains.

3. Recovering conductivity

We write $\Omega = I =]a, b[\subset \mathbf{R}$. Assume $1 < p^- \leq p(x) \leq p^+ < \infty$ almost everywhere on I , and that p is a bounded measurable function. In fact, define

$$p^- = \operatorname{ess\,inf}_{x \in I} p(x) \text{ and } p^+ = \operatorname{ess\,sup}_{x \in I} p(x).$$

Later, we will use similar notations for the essential supremum and infimum of another exponent $q(\cdot)$, which is the conjugate Hölder exponent to $p(\cdot)$.

Then, at least formally, the forward problem is

$$\begin{cases} - \left(\gamma |u'|^{p(x)-2} u' \right)' = 0 \text{ in } I, \\ u(a) = A, \\ u(b) = B. \end{cases}$$

We suppose $A \leq B$, whence there should exist a constant $K \geq 0$ such that, for almost every x ,

$$\gamma (u')^{p(x)-1} = K \iff u' = (K/\gamma)^{1/(p(x)-1)}$$

and hence

$$u(x) = A + \int_a^x (K/\gamma(s))^{1/(p(s)-1)} \, ds.$$

Using $u(b) = B$ we get

$$(3) \quad \int_a^b (K/\gamma(s))^{1/(p(s)-1)} \, ds = B - A.$$

Writing $m = B - A$, we have implicitly defined a function $K_m = K : \mathbf{R}_+ \rightarrow \mathbf{R}_+$ by writing as K_m the constant K which satisfies the above equation with $m = B - A$.

The next lemma justifies the previous heuristic discussion, and also implies that the constant K_m is unique, since the minimizer is unique.

Lemma 6. *Suppose that for a given $B - A = m \geq 0$, there exists a function $v \in W^{1,p}(I)$ satisfying the boundary values (A, B) , and a constant K_m for which the following equality is true almost everywhere in x :*

$$\gamma(x) (v'(x))^{p(x)-1} = K_m.$$

Then v is the unique minimizer of energy (2) with boundary values A and B , and thus solves the variable exponent conductivity equation.

We use the same proof as Diening, Harjulehto, Hästö and Růžička [14, Lemma 13.1.4], but they consider a slightly different equation.

Proof. By lemma 5 there exists a unique minimizer. As such, we only need to establish that v minimizes the energy.

Let w be such that $w - v \in W_0^{1,p(\cdot)}(I)$. Since $\gamma(x) (v'(x))^{p(x)-1}$ is constant almost everywhere,

$$\int_a^b \gamma(x) (v'(x))^{p(x)-1} (w' - v') \, dx = 0.$$

We use the inequality

$$\frac{1}{p} |y|^p \geq \frac{1}{p} |z|^p + |z|^{p-1} (y - z),$$

which follows from the convexity of the differentiable function $y \mapsto \frac{1}{p} |y|^p$. The inequality implies

$$\int_a^b \frac{\gamma(x)}{p(x)} |w'|^{p(x)} \, dx \geq \int_a^b \frac{\gamma(x)}{p(x)} |v'|^{p(x)} \, dx + \int_a^b \gamma(x) |v'|^{p(x)-1} (w' - v') \, dx,$$

which implies, since the last integral is zero,

$$\int_a^b \frac{\gamma(x)}{p(x)} |w'|^{p(x)} \, dx \geq \int_a^b \frac{\gamma(x)}{p(x)} |v'|^{p(x)} \, dx. \quad \square$$

Lemma 7. *The map $m \mapsto K_m$ is a strictly increasing, continuous bijection.*

Proof. It suffices to prove that the map

$$K \mapsto \int_a^b (K/\gamma(x))^{1/(p(x)-1)} \, dx$$

is a strictly increasing surjection.

Since $1 < p < \infty$, the map is strictly increasing. We have both $0 \mapsto 0$ and $\int_a^b (K/\gamma(x))^{1/(p(x)-1)} \, dx \rightarrow \infty$ as $K \rightarrow \infty$, since γ and $1/(p-1)$ are positive. The integrand is continuous with respect to K for almost every $x \in I$, which implies

continuity via dominated convergence, given the bounded interval and boundedness of γ and p . Hence, we have surjectivity. \square

The DN map is $\Lambda_\gamma: \mathbf{R}^2 \rightarrow \mathbf{R}$, $\Lambda_\gamma(A, B) = \int_a^b \gamma |u'|^p(x) dx$, as established in Section 2. Assuming $A \leq B$ and inserting u' gives the following lemma:

Lemma 8. *Suppose $A \leq B$. Then*

$$\Lambda_\gamma(A, B) = \Lambda_\gamma(m) = \int_a^b \gamma^{-1/(p(x)-1)} K_m^{p(x)/(p(x)-1)} dx,$$

where K_m is also a function of the conductivity γ .

First we observe that we can recover $\int_a^b \gamma^{-1/(p(x)-1)} dx$ from the Dirichlet to Neumann map as its unique fixed point.

Lemma 9. *Suppose $B > A$.*

- *If $B - A > \int_a^b \gamma^{-1/(p(x)-1)} dx$, then $\Lambda_\gamma(B - A) > B - A$. Also, $K > 1$.*
- *If $B - A = \int_a^b \gamma^{-1/(p(x)-1)} dx$, then $\Lambda_\gamma(B - A) = B - A$. Also, $K = 1$.*
- *If $B - A < \int_a^b \gamma^{-1/(p(x)-1)} dx$, then $\Lambda_\gamma(B - A) < B - A$. Also, $K < 1$.*

Proof. The positive number k is a fixed point of the DN map if and only if

$$k = \Lambda_\gamma(k) = \int_a^b \gamma^{-1/(p(x)-1)} K_k^{p(x)/(p(x)-1)} dx$$

where

$$k = \int_a^b (K_k/\gamma)^{1/(p(x)-1)} dx,$$

which implies

$$\int_a^b \gamma^{-1/(p(x)-1)} K_k^{p(x)/(p(x)-1)} dx = \int_a^b \gamma^{-1/(p(x)-1)} K_k^{1/(p(x)-1)} dx.$$

This is true if and only if $K_k = 1$, since $\gamma > 0$ and $p > 1$.

If $K = 1$, then by equation (3) we have $B - A = \int_I \gamma^{-1/(p(x)-1)} dx$. By bijectivity of $m \mapsto K_m$, there only exists one m with $K_m = 1$. This proves the middle point of the claim.

If $m > k$, then $K_m > 1$, and hence

$$K_m^{p(x)/(p(x)-1)} > K_m^{1/(p(x)-1)},$$

which implies

$$\begin{aligned} \Lambda_\gamma(B - A) &= \int_a^b \gamma^{-1/(p(x)-1)} K_m^{p(x)/(p(x)-1)} dx \\ &> \int_a^b \gamma^{-1/(p(x)-1)} K_m^{1/(p(x)-1)} dx = B - A. \end{aligned}$$

The same argument with reversed inequalities holds when $m < k$. \square

By using, for example, the half-interval search we get:

Corollary 10. *The quantity*

$$\int_I \gamma^{-1/(p(x)-1)} dx$$

can be recovered from the Dirichlet to Neumann map.

The next remark concerns the inverse problem with additional interior data of the type that can, under some idealizations, be recovered using hybrid imaging methods such as ultrasound mediated electrical impedance tomography (UMEIT, also called AET for acousto-electric tomography), conductivity density impedance imaging (CDII) and magnetic resonance electrical impedance tomography (MREIT). [1, 25, 27]

Remark 11. (Interior data and variable exponent) If we have knowledge of interior power data of type $\gamma |u'|^{r(x)}$, where $0 \leq r(x) < \infty$, then the conductivity can be recovered at all points where $p(x) - r(x) \neq 1$. Indeed, a simple calculation gives

$$\gamma |u'|^{r(x)} = \gamma (K/\gamma(x))^{r(x)/(p(x)-1)} = (\gamma(x))^{(p(x)-r(x)-1)/(p(x)-1)} K^{r(x)/(p(x)-1)}.$$

We can choose the Dirichlet data $B - A$ so that $K = 1$ by lemma 7. Hence, γ can be recovered whenever it has a nonzero exponent, or, equivalently, whenever $p(x) - r(x) \neq 1$.

This generalizes a result of Straub [33, chapter 3], which was for $p \equiv 2$.

3.1. Identification at extremes. Next we recover the value of γ at the points where $p(x)$ takes its maximum or minimum value. First write

$$q(x) = \frac{p(x)}{p(x) - 1} \quad \text{and} \quad f(x) = (\gamma(x))^{-1/(p(x)-1)}.$$

These are both injective mappings of $p(x)$ and $\gamma(x)$, respectively, and q is the Hölder dual exponent of p . The maxima of p correspond to the minima of q and vice versa.

Lemma 12. *Suppose the exponent q reaches its essential supremum (respectively infimum) value q^+ (q^-) on a set Q_+ (Q_-) of positive measure. Then*

$$\begin{aligned} \lim_{K \rightarrow \infty} K^{-q^+} \int_a^b f(x) K^{q(x)} dx &= \int_{Q_+} f(x) dx \quad \text{and} \\ \lim_{K \rightarrow 0} K^{-q^-} \int_a^b f(x) K^{q(x)} dx &= \int_{Q_-} f(x) dx. \end{aligned}$$

Proof. For the maximum, by monotone (or dominated) convergence

$$\int_{I \setminus Q_+} K^{q(x)-q^+} f(x) dx \rightarrow 0,$$

since $q(x) - q^+ < 0$ on the set. The integral over Q_+ gives what we claim in the lemma. The argument for the minimum has precisely the same idea. \square

Unfortunately, K_m is not something we can recover from the measurements. We define an auxiliary variable \overline{K}_m , which corresponds to conductivity one and thus is characterized by the equation

$$m = \int_a^b \overline{K}_m^{1/(p(x)-1)} dx$$

and can be calculated without knowing the true conductivity. We write as c_m constants that satisfy the inequality

$$1/c_m \leq K_m/\overline{K}_m \leq c_m.$$

Let K_m^l correspond to the constant weight $\max(1, \text{ess sup}_{x \in I} f(x))$ and K_m^u to $\min(1, \text{ess inf}_{x \in I} f(x))$. That is, we have

$$m = \int_a^b \max \left(1, \text{ess sup}_{x \in I} f(x) \right) (K_m^l)^{1/(p(x)-1)} \, dx \quad \text{and}$$

$$m = \int_a^b \min \left(1, \text{ess inf}_{x \in I} f(x) \right) (K_m^u)^{1/(p(x)-1)} \, dx.$$

This implies $K_m^l \leq K_m \leq K_m^u$ and $K_m^l \leq \overline{K}_m \leq K_m^u$, and thereby K_m/\overline{K}_m and its inverse are bounded by K_m^u/K_m^l . To bound this we use a mean value theorem:

Lemma 13. *Suppose $p \in L^\infty([a, b])$ and suppose $h: A \rightarrow \mathbf{R}$, with the range $[\text{ess inf } p, \text{ess sup } p] \subseteq A$, is a continuous function with*

$$\int_a^b h(p(x)) \, dx = 0.$$

Then there exists a real number $p^ \in [\text{ess inf } p, \text{ess sup } p]$ such that*

$$\int_a^b h(p^*) \, dx = 0.$$

Proof. If no p^* with $h(p^*) = 0$ existed, then by continuity of h , $h(p)$ would be either positive for all $p \in [\text{ess inf } p, \text{ess sup } p]$, or negative for all of them. This contradicts the assumption that the integral of $h(p(x))$ is zero. \square

Lemma 14. *For all $m \in \mathbf{R}_+$ there exists $p^* \in [p^-, p^+]$ such that we can choose*

$$c_m \leq \left(\frac{\max(1, \text{ess sup}_{x \in I} f(x))}{\min(1, \text{ess inf}_{x \in I} f(x))} \right)^{p^*-1}.$$

Proof. We define h by

$$h(p) = (K_m^u)^{1/(p-1)} \min \left(1, \text{ess inf}_{x \in I} f(x) \right) - (K_m^l)^{1/(p-1)} \max \left(1, \text{ess sup}_{x \in I} f(x) \right)$$

and use the mean value lemma (lemma 13). The claim follows from

$$\begin{aligned} & \int_a^b \max \left(1, \text{ess sup}_{x \in I} f(x) \right) (K_m^l)^{1/(p^*-1)} \, dx \\ &= \int_a^b \min \left(1, \text{ess inf}_{x \in I} f(x) \right) (K_m^u)^{1/(p^*-1)} \, dx. \end{aligned} \quad \square$$

We have thus established that the ratio K_m/\overline{K}_m is bounded uniformly in m , since p^* is bounded. We use this information to determine the limit of the ratio as $m^{\pm 1} \rightarrow \infty$.

Lemma 15. *Suppose Q_+ (respectively Q_-) has positive measure. Then*

$$\lim_{m \rightarrow \infty} \left(\frac{\overline{K}_m}{K_m} \right)^{1/(p^- - 1)} = \int_{Q^+} f(x) \, dx,$$

$$\lim_{m \rightarrow 0} \left(\frac{\overline{K}_m}{K_m} \right)^{1/(p^+ - 1)} = \int_{Q^-} f(x) \, dx.$$

Proof. First we prove

$$\int_{Q_+} f(x) dx K_m^{1/(p^- - 1)} = |Q_+| \bar{K}_m^{1/(p^- - 1)} + o_{m \rightarrow \infty} \left(K_m^{1/(p^- - 1)} + \bar{K}_m^{1/(p^- - 1)} \right)$$

and

$$\int_{Q_-} f(x) dx K_m^{1/(p^+ - 1)} = |Q_-| \bar{K}_m^{1/(p^+ - 1)} + o_{m \rightarrow 0} \left(K_m^{1/(p^+ - 1)} + \bar{K}_m^{1/(p^+ - 1)} \right).$$

For $m \rightarrow \infty$ we calculate

$$\begin{aligned} \int_{Q_+} f(x) dx K_m^{1/(p^- - 1)} &= \int_{Q_+} f(x) K_m^{1/(p(x) - 1)} dx \\ &= \int_a^b f(x) K_m^{1/(p(x) - 1)} dx + o_{m \rightarrow \infty} \left(K_m^{1/(p^- - 1)} \right) \\ &= \int_a^b \bar{K}_m^{1/(p(x) - 1)} dx + o_{m \rightarrow \infty} \left(K_m^{1/(p^- - 1)} \right) \\ &= |Q_+| \bar{K}_m^{1/(p^- - 1)} + o_{m \rightarrow \infty} \left(K_m^{1/(p^- - 1)} + \bar{K}_m^{1/(p^- - 1)} \right). \end{aligned}$$

Next we divide by $|Q_+| K_m^{1/(p^- - 1)}$ and use lemma 14 to get the claim. The argument when $m \rightarrow 0$ is essentially the same. \square

Theorem 4 directly follows from the following proposition.

Proposition 16. *Suppose Q_+ or Q_- is of positive measure. Then*

$$\begin{aligned} \lim_{m \rightarrow \infty} \bar{K}_m^{-q^+} \frac{1}{|Q_+|} \Lambda_\gamma(m) &= \left(\int_{Q_+} f(x) dx \right)^{-(p^- - 1)} \quad \text{or} \\ \lim_{m \rightarrow 0} \bar{K}_m^{-q^-} \frac{1}{|Q_-|} \Lambda_\gamma(m) &= \left(\int_{Q_-} f(x) dx \right)^{-(p^+ - 1)}. \end{aligned}$$

Proof.

$$\bar{K}_m^{-p^-/(p^- - 1)} \frac{1}{|Q_+|} \Lambda_\gamma(m) = \left(\frac{\bar{K}_m}{K_m} \right)^{-q^+} K_m^{-q^+} \frac{1}{|Q_+|} \int_a^b f(x) K_m^{q(x)} dx.$$

By lemma 12 we have, as $m \rightarrow \infty$,

$$K_m^{-q^+} \frac{1}{|Q_+|} \int_a^b f(x) K_m^{q(x)} dx \rightarrow \int_{Q_+} f(x) dx.$$

The proof then follows from lemma 15. \square

3.2. Characterization of recognizable functions. In this section we characterize the space of functions with which we pair the unknown function f in the DN map

$$m \mapsto \int_I f(x) K_m^{q(x)} dx.$$

We write $k_m = \log(K_m)$ and consider the set of functions

$$S = \{ \exp(k_m q(x)); m \in \mathbf{R} \}.$$

Since S is closed under pointwise multiplication, any product of linear combinations of elements from S remains a linear combination of elements from S . In particular, the space $\text{span}(S)$ is an algebra.

In the special case that q is continuous and injective, the Stone-Weierstrass theorem implies that $\text{span}(S)$ is dense in $C(I)$ with the usual topology of uniform convergence. Since the continuous functions are dense and continuously embedded in $L^2(I)$ (the interval I is bounded), $\text{span}(S)$ must also be dense in $L^2(I)$ by approximating in L^2 -norm and choosing a diagonal sequence from $\text{span}(S)$. This result we shall now generalize to the setting where q is merely measurable by using the multiplicative system theorem [13, Theorem 21].

Theorem 17. (Multiplicative system theorem) *Suppose H is a vector space of real-valued bounded measurable functions on a measurable space X . Suppose H contains constant functions and is closed under the pointwise convergence of uniformly bounded increasing sequences of functions. Let $M \subseteq H$ be closed under pointwise multiplication, and let \mathcal{G} be the σ -algebra generated by M . Then H contains all bounded \mathcal{G} -measurable functions.*

Suppose $r > 1$ is a positive real number and $s > 1$ is its Hölder conjugate. Note that, when p and q are Hölder conjugates, $\sigma(p) = \sigma(q)$, since the map taking p to q is a homeomorphism from $]1, \infty[$ to itself. Hence p and q generate the same σ -algebra.

Lemma 18.

$$\overline{\text{span}(S)}^{L^r(I)} = L^r(I, \sigma(p)).$$

The proof follows a proof of Eldredge [16] for a similar lemma. We omit the space L^r from the notation of the closure.

Proof. The σ -algebra generated by $\text{span}(S)$ is exactly the sigma-algebra $\sigma(p)$. One consequence is that $\text{span}(S)$ is a subspace of $L^r(I, \sigma(p))$.

First, we show $\overline{\text{span}(S)} \subseteq L^r(I, \sigma(p))$. Pick an element in the L^r -closure. Then there is a Cauchy sequence in $\text{span}(S)$ with this point as its limit. Any Cauchy sequence in $L^r(I)$ has an almost everywhere convergent subsequence [30, Theorems 3.11 and 3.12], and since these functions are measurable with respect to $\sigma(p)$, so is the limit. This gives the first inclusion.

The reverse inclusion will follow from the multiplicative systems theorem. Define $M = \text{span}(S)$ and let H consist of all bounded measurable functions belonging to the equivalence classes of functions inside $\overline{\text{span}(S)} \cap L^\infty(I)$. We note that $M \subset H$ is closed under pointwise multiplication.

- (1) H contains constant functions, since $1 \in S \subset H$ by taking $k_m = 0$.
- (2) H is closed under pointwise convergence of uniformly bounded increasing sequences. To see this, take such a sequence $(f_j)_{j=1}^\infty$ in H converging pointwise to a measurable function f . Then f is bounded and $|f - f_j|^r$ converges pointwise to zero. Uniform boundedness and the dominated convergence theorem together imply that $f_j \rightarrow f$ in the L^r -norm as $j \rightarrow \infty$. Thus the equivalence class of f is in $\overline{\text{span}(S)}$. Hence $f \in H$.

Theorem 17 ensures H contains all bounded $\sigma(p)$ -measurable functions. By this construction, $L^r(I, \sigma(p)) \cap L^\infty(I, \sigma(p)) \subset \overline{\text{span}(S)} \cap L^\infty(I)$ holds. Note that we here use that the sigma-algebra generated by M is actually $\sigma(p)$. Take $h \in L^r(I, \sigma(p))$ and let h also signify a $\sigma(p)$ -measurable representative. Construct a sequence $(h_j)_{j=1}^\infty$ by setting

$$h_j = \max\{-j, \min\{h, j\}\}.$$

These are bounded and $\sigma(p)$ -measurable. By the above, they belong to H . The dominated convergence theorem implies $\overline{h_j \rightarrow h}$ as $j \rightarrow \infty$ in the L^r -norm. Since each h_j is in some equivalence class of $\text{span}(S)$, then $h \in \overline{\text{span}(S)}$ also. This gives the reverse inclusion. \square

Above, we demonstrated that the functions $\exp(k_m q)$, spanning a dense subspace of $L^r(I, \sigma(p))$, suffice to determine $f \in L^s(I, \sigma(p))$ uniquely. Hence f can, in principle, be recovered from measurements of the DN map across all $m \in \mathbf{R}$, provided that it belongs to this space. In general it belongs to the space $L^\infty(I)$, as γ and p are bounded, and thereby to all the L^r spaces.

When q is continuous and increasing, we are unable to recover f in sets where q is constant. This is because $K_m^{q(\cdot)}$ restricted to any such set remains constant upon varying $m \in \mathbf{R}$, so testing against these yields no insight beyond the average of f inside such a set. On the other hand, if q is even on an interval symmetric about the origin, then we can only hope to determine f up to its even part, because all $\exp(k_m q)$ are even in this case.

In abstract language, what we have determined is the projection of $f \in L^\infty(I) \subset L^2(I)$ onto the (complete and therefore closed) subspace $L^2(I, \sigma(p))$. This projection can also be understood in terms of the conditional expectation $\mathbf{E}(f | \sigma(p))$ given $\sigma(p)$ as the algebra of observable events [2, Theorem 3.2.6].

Let $P: L^2(I) \rightarrow L^2(I, \sigma(p))$ be the orthogonal projection onto $L^2(I, \sigma(p))$. We formulate the above as a statement about the reconstructibility of Pf from a collection of measurements of the DN map:

Proposition 19. *There exists an orthonormal sequence $\{s_n\}_{n=1}^\infty \subset \text{span}(S)$ determining the projection $Pf = \sum_{n=1}^\infty \langle f, s_n \rangle s_n$ for any $f \in L^2(I)$.*

Proof. Since $L^2(I, \sigma(p))$ identifies with a subspace of $L^2(I)$, it is also separable. By lemma 18 there is a linearly independent countable dense subset of $\text{span}(S)$. Orthonormalization by the Gram–Schmidt process gives the vectors. \square

Every s_n above is some finite linear combination of the functions $\exp(k_m q)$, and the coefficients are determined from finite combinations of $\langle f, \exp(k_m q) \rangle$, the measurements of the DN map.

By the above, we only need countably many measurements of the DN map, but since the functions s_n depend implicitly on the unknown conductivity γ , we have no way to determine which Dirichlet data m to use beforehand. It is not possible to explicitly reconstruct Pf using proposition 19.

Functions that can not be detected are simply those belonging to $\ker(P)$, and we have the following simple characterization:

Proposition 20. *The kernel $\ker(P)$ consists of all those functions $f \in L^2(I)$ that integrate to zero on every $\sigma(p)$ -set.*

Proof. A function with the stated properties is in $\ker(P)$ by proposition 19, because any s_n can be approximated in $L^2(I, \sigma(p))$ by using simple functions. Conversely, if $Pf = 0$, then the integral of f taken over any $\sigma(p)$ -set is zero, because $\text{span}\{s_n\}_{n=1}^\infty$ can approximate any characteristic of a $\sigma(p)$ -set. \square

Example 21. If I is symmetric around the origin and p is an even function, then all the $\sigma(p)$ -sets are also symmetric about the origin, and by the above, all odd functions on I must be in $\ker(P)$.

Finally, in terms of the conductivity γ , which is the function of interest, it is natural to use P to define a nonlinear mapping

$$\tilde{P}: L^2(I) \cap L^{\infty}_+(I) \rightarrow L^2(I, \sigma(p)) \cap L^{\infty}_+(I, \sigma(p)): \gamma \mapsto \left(P(\gamma^{-\frac{1}{p-1}})\right)^{-(p-1)},$$

To see that this is well-defined, we define, via representatives g , the mapping

$$\Phi: L^2(I) \cap L^{\infty}_+(I) \rightarrow L^2(I) \cap L^{\infty}_+(I): g \mapsto g^{-\frac{1}{p-1}},$$

and put $\tilde{P} = \Phi^{-1}P\Phi$. Since we have the bounds $1 < p^- < p(x) < p^+ < \infty$, then Φ maps $L^{\infty}_+(I)$ to itself, and hence to $L^2(I)$. We note that Φ is invertible. As an aside, we mention that P and \tilde{P} are topologically conjugate [28, Section 4.7] as continuous maps on $L^{\infty}_+ \cap L^2$ in the topology of L^{∞}_+ , since Φ is a homeomorphism from L^{∞}_+ to itself.

Lemma 22. *Provided that g is $\sigma(p)$ -measurable, then $g^{-\frac{1}{p-1}}$ is also $\sigma(p)$ -measurable. The inverse $g^{-(p-1)}$ has the same property.*

Proof. We consider the function

$$x \mapsto (g(x))^{-\frac{1}{p(x)-1}}.$$

There exist uniformly bounded sequences of $\sigma(p)$ -measurable simple functions, denoted (g_j) and (h_j) , converging pointwise to g and the exponent, respectively. Then $(g_j^{h_j})$ converges pointwise to $g^{-\frac{1}{p-1}}$. The proof is similar for $g^{-(p-1)}$. \square

The projection P also maps $L^{\infty}_+(I)$ into $L^{\infty}_+(I, \sigma(p))$. Indeed, observe that $\text{ess sup}(Pg) \leq \text{ess sup}(g)$ for any $g \in L^{\infty}(I)$, since the averages of Pg and g over any $\sigma(p)$ set must be equal, and the set $\{x \in I; Pg(x) > \text{ess sup}(g) + \varepsilon\}$ has measure zero for every $\varepsilon > 0$. Similarly $\text{ess inf}(g) \leq \text{ess inf}(Pg)$.

Therefore $\tilde{P} = \Phi^{-1}P\Phi$ is well-defined with the desired mapping properties. It inherits the projection property $\tilde{P} \circ \tilde{P} = \tilde{P}$ from P . This proves Theorem 1.

3.3. Derivatives of the Dirichlet to Neumann map. In this section our goal is to give a constructive, if ill-posed and inconvenient, alternative to the nonconstructive result in previous Section 3.2. Furthermore, the result here is local in the sense that we only need to know the Dirichlet-to-Neumann map in a neighbourhood of the value m for which $K(m) = 1$, which we will write as k . That is, we have $K(k) = 1$. We also write K_m as $K(m)$ to emphasize the dependence on m .

We will state explicit formulae for high order derivatives of inverse and compound functions. The formulae are used to calculate $\frac{d^n m}{dK^n}|_{m=k}$ in terms of derivatives of the DN map, which then allows calculating $\int_I f(x) \left(\frac{1}{p(x)-1}\right)^n dx$ and proceeding as in the previous section.

Recall that

$$\Lambda_{\gamma}(m) = \int_I f(x) (K(m))^{q(x)} dx,$$

which we now consider as a function of K :

$$\Lambda_{\gamma}(K) = \int_I f(x) K^{q(x)} dx.$$

Also recall

$$m = \int_I f(x) K^{1/(p(x)-1)} dx,$$

whence

$$\frac{d^j m}{dK^j} = \int_I f(x) K^{1/(p(x)-1)-j} \prod_{l=0}^{j-1} (1/(p(x)-1) - l) \, dx.$$

We record the following formula for higher order derivatives of inverse functions [23]. The facts about the indices follow by elementary manipulation.

Lemma 23. *Suppose $n \in \mathbf{Z}_+$, for all j we have $s_j \in \mathbf{N}$, m is a C^n -function of K , and $\frac{dm}{dK} \neq 0$ on an open interval. Then, on that interval,*

$$\frac{d^n K}{dm^n} = \frac{(-1)^{n-1}}{\left(\frac{dm}{dK}\right)^{2n-1}} \sum_{\substack{s_1+s_2+\dots+s_{n-1} \\ 1 \cdot s_1+2 \cdot s_2+\dots+2n-2}} \frac{(-1)^{s_1} (2n - s_1 - 2)! \prod_{j \in \mathbf{Z}_+} \left(\frac{d^j m}{dK^j}\right)^{s_j}}{\prod_{j=2}^\infty (j!)^{s_j} s_j!}.$$

Furthermore, we have:

- $s_j \neq 0$ implies $j \leq n$.
- $s_n \neq 0$ implies the index tuple $(s_1, s_2, \dots, s_{n-1}, s_n) = (n - 2, 0 \dots, 0, 1)$ (and $s_1 = 1$ when $n = 1$).

We note that the derivative $\frac{dm}{dK}$ is indeed nonzero, though higher order derivatives might not be, if p is constant.

Next we state the well known Faà di Bruno’s formula [22].

Lemma 24. (Faà di Bruno’s formula) *For n times continuously differentiable $K \mapsto \Lambda_\gamma$ and $m \mapsto K$*

$$\frac{d^n}{dm^n} \Lambda_\gamma(K(m)) = \sum \frac{n!}{k_1! k_2! \dots k_n!} \cdot \Lambda_\gamma^{(k_1+\dots+k_n)}(K(m)) \cdot \prod_{j=1}^n \left(\frac{K^{(j)}(m)}{j!}\right)^{k_j},$$

where the sum is over n -tuples of numbers k_j with

$$\sum_{j=1}^n j k_j = n.$$

We have

$$\frac{d^j \Lambda_\gamma}{dK^j}(K) = \int_I f(x) K^{q(x)-j} \prod_{l=0}^{j-1} (q(x) - l) \, dx.$$

Hence, the derivative of the DN map with respect to m at $m = k$ is

$$\frac{d\Lambda_\gamma(K(m))}{dm} \Big|_{m=k} = \frac{\int_I f(x) \frac{p(x)}{p(x)-1} \, dx}{\int_I f(x) \frac{1}{p(x)-1} \, dx}.$$

We observe that since $p(x) > p^- > 1$, the derivative is always greater than 1. We already know $\int_I f(x) \, dx$ due to corollary 10. Using $\frac{p}{p-1} = 1 + \frac{1}{p-1}$ we get

$$\int_I \frac{1}{p(x)-1} f(x) \, dx = \int_I f(x) \, dx \left(\frac{d}{dm} \Lambda_\gamma(1) - 1 \right)^{-1},$$

which is well-defined, since the DN map is strictly greater than one. Hence, we now know the dual pairing of f against 1 and $1/(p(x) - 1)$. This proves the following lemma:

Lemma 25.

$$\int_I f(x) \frac{1}{p(x) - 1} dx$$

can be recovered from

$$\Lambda_\gamma(K(k)) \quad \text{and} \quad \left. \frac{d\Lambda_\gamma(K(m))}{dm} \right|_{m=k}.$$

We prove the analogous claim for higher powers of $\frac{1}{p(x)-1}$ by induction.

Proposition 26. *Suppose $n \in \mathbf{N}$. The quantity*

$$\int_I f(x) \left(\frac{1}{p(x) - 1} \right)^n dx$$

can be recovered from

$$\left. \frac{d^l \Lambda_\gamma(K(m))}{dm^l} \right|_{m=k} \quad \text{for } l \in \{0, \dots, n\}.$$

Proof. Corollary 10 establishes the case $n = 0$ and lemma 25 the case $n = 1$. We proceed by strong induction and suppose $n > 1$. That is, we assume that the following integrals are known for all $j < n$:

$$\int_I f(x) \left(\frac{1}{p(x) - 1} \right)^j dx.$$

In Faà di Bruno’s formula for the derivative of order n there is only one n -tuple where $\frac{d^n \Lambda_\gamma}{dK^n}$ appears, namely $(n, 0, \dots, 0)$. Likewise, in the formula for the higher order inverse function, there is only one term with derivative of order n , $(n - 2, 0, \dots, 0, 1)$. This latter term appears in Faà di Bruno’s formula precisely when the n -tuple is $(0, \dots, 1)$. The other summands in the formulae only contain known terms, not

$$\int_I f(x) \left(\frac{1}{p(x) - 1} \right)^n dx.$$

By Faà di Bruno’s formula we write the derivative as

$$\left. \frac{d^n \Lambda_\gamma}{dm^n} \right|_{m=k} = \left. \frac{d^n \Lambda_\gamma}{dK^n} \right|_{K=1} \cdot \left(\left. \frac{dm}{dK} \right|_{K=1} \right)^{-n} + \left. \frac{d\Lambda_\gamma}{dK} \right|_{K=1} \left. \frac{d^n K}{dm^n} \right|_{m=k} + S_1,$$

where S_1 is known by the induction hypothesis. Rewriting, we have

$$\begin{aligned} \left. \frac{d^n \Lambda_\gamma}{dm^n} \right|_{m=k} - S_1 &= \left. \frac{d^n m}{dK^n} \right|_{K=1} \left(\int_I f(x) \frac{1}{p(x) - 1} dx \right)^{-n} \\ &\quad + \int_I f(x) \frac{p(x)}{p(x) - 1} dx \left. \frac{d^n K}{dm^n} \right|_{m=k}, \end{aligned}$$

where, by lemma 23,

$$\left. \frac{d^n K}{dm^n} \right|_{m=k} = - \left(\left. \frac{dm}{dK} \right|_{K=1} \right)^{-(n+1)} \left. \frac{d^n m}{dK^n} \right|_{K=1} + S_2,$$

where S_2 is known by the induction hypothesis. This gives

$$\begin{aligned} & \frac{d^n \Lambda_\gamma}{dm^n} \Big|_{m=k} - S_1 - S_2 \int_I f(x) \frac{p(x)}{p(x) - 1} dx \\ &= \frac{d^n m}{dK^n} \Big|_{K=1} \left(\int_I f(x) \frac{1}{p(x) - 1} dx \right)^{-n} \\ & \quad - \left(\frac{dm}{dK} \Big|_{K=1} \right)^{-(n+1)} \frac{d^n m}{dK^n} \Big|_{K=1} \int_I f(x) \frac{p(x)}{p(x) - 1} dx. \end{aligned}$$

Using $p/(p - 1) = 1 + 1/(p - 1)$ we get

$$\begin{aligned} & \frac{d^n \Lambda_\gamma}{dm^n} \Big|_{m=k} - S_1 - S_2 \int_I f(x) \frac{p(x)}{p(x) - 1} dx \\ &= - \left(\frac{dm}{dK} \Big|_{K=1} \right)^{-(n+1)} \frac{d^n m}{dK^n} \Big|_{K=1} \int_I f(x) dx. \end{aligned}$$

Because everything else on the right hand side is positive, this gives an explicit formula for

$$\frac{d^n m}{dK^n} \Big|_{K=1} = \int_I f(x) \prod_{l=0}^{n-1} (1/(p(x) - 1) - l) dx,$$

where $\prod_{l=0}^{n-1} (1/(p(x) - 1) - l)$ is a polynomial in $1/(p(x) - 1)$ of order n and with leading coefficient 1 for its highest order term. By the induction hypothesis, the dual pairings of all the terms but the highest order one with f are known. This proves the claim. \square

The following lemma is very similar to Lemma 18 and [6]. The proof also follows the proof of Eldredge [16].

Lemma 27. *Let $1 < r < \infty$. Then*

$$\overline{\text{span} \left(\left\{ \left(\frac{1}{p(x) - 1} \right)^n ; n \in \mathbf{N} \right\} \right)^{L^r(I)}} = L^r(I, \sigma(p)).$$

Proof. We first observe that $p \mapsto 1/(p - 1)$ is a homeomorphism from $]1, \infty[$ to $]0, \infty[$. This implies that $\sigma(p) = \sigma(1/(p - 1))$. Since p is bounded away from one and infinity, $1/(p - 1)$ is likewise bounded and hence

$$\text{span} \left(\left\{ \left(\frac{1}{p(x) - 1} \right)^n ; n \in \mathbf{N} \right\} \right) \subseteq L^r(I).$$

The same argument as in the proof of lemma 18 establishes

$$\text{span} \left(\left\{ \left(\frac{1}{p(x) - 1} \right)^n ; n \in \mathbf{N} \right\} \right) \subseteq L^r(I, \sigma(p)).$$

The proof for the inclusion

$$L^r(I, \sigma(p)) \subseteq \text{span} \left(\left\{ \left(\frac{1}{p(x) - 1} \right)^n ; n \in \mathbf{N} \right\} \right)$$

is essentially the same as in lemma 18, though in this case $n = 0$ gives that the function $1 \in H$. In particular, multiplying two linear combinations of polynomials (of $1/(p - 1)$) still gives a polynomial. \square

Proving Theorem 1 proceeds as in the non-constructive case.

References

- [1] BAL, G.: Hybrid inverse problems and internal functionals. - In: Inverse problems and applications: inside out II (edited by G. Uhlmann), Math. Sci. Res. Inst. Publ. 60, Cambridge Univ. Press, 2013, 325–368.
- [2] BOBROWSKI, A.: Functional analysis for probability and stochastic processes: an introduction. - Cambridge Univ. Press, 2005.
- [3] BRANDER, T.: Calderón problem for the p -Laplacian: First order derivative of conductivity on the boundary. - Proc. Amer. Math. Soc. 144, 2016, 177–189.
- [4] BRANDER, T.: Calderón’s problem for p -Laplace type equations. - PhD thesis, Reports of University of Jyväskylä Department of Mathematics and Statistics 155, 2016.
- [5] BRANDER, T., J. ILMAVIRTA, and M. KAR: Superconductive and insulating inclusions for linear and non-linear conductivity equations. - Inverse Probl. Imaging 12:1, 2018, 91–123.
- [6] BRANDER, T., J. ILMAVIRTA, and T. TYNI: Optimal recovery of a radiating source with multiple frequencies along one line. - In preparation.
- [7] BRANDER, T., M. KAR, and M. SALO: Enclosure method for the p -Laplace equation. - Inverse Problems 31:4, 2015, 045001, 16.
- [8] BRANDER, T., B. VON HARRACH, M. KAR, and M. SALO: Monotonicity and enclosure methods for the p -Laplace equation. - SIAM J. Appl. Math. 78:2, 2018, 742–758.
- [9] BUENO, P. R., J. A. VARELA, and E. LONGO: SnO₂, ZnO and related polycrystalline compound semiconductors: An overview and review on the voltage-dependent resistance (non-ohmic) feature. - J. European Ceramic Society 28:3, 2008, 505–529.
- [10] CALDERÓN, A. P.: On an inverse boundary value problem. - In: Seminar on numerical analysis and its applications to continuum physics (edited by W. H. Meyer and M. A. Raupp), Sociedade Brasileira de Matematica, 1980, 65–73. Reprinted as [11].
- [11] CALDERÓN, A. P.: On an inverse boundary problem. - Comput. Appl. Math. 25:2-3, 2006, 133–138. Reprint of [10].
- [12] DE HOOP, M., G. UHLMANN, and Y. WANG: Nonlinear interaction of waves in elastodynamics and an inverse problem. - Preprint, arXiv:1805.03811, 2018.
- [13] DELLACHERIE, C., and P.-A. MEYER: Probabilities and potential. - North-Holland Math. Stud. 29, North-Holland Publishing Co., Amsterdam – New York, 1978.
- [14] DIENING, L., P. HARJULEHTO, P. HÄSTÖ, and M. RŮŽIČKA: Lebesgue and Sobolev spaces with variable exponents. - Lecture Notes in Math. 2017, Springer, 2011.
- [15] DUBSON, M. A., S. T. HERBERT, J. J. CALABRESE, D. C. HARRIS, B. R. PATTON, and J. C. GARLAND: Non-Ohmic dissipative regime in the superconducting transition of polycrystalline Y₁Ba₂Cu₃O_x. - Phys. Rev. Lett. 60:11, 1988, 1061–1064.
- [16] ELDREDGE, N.: Closure of polynomials of a function in L^2 . - MathOverflow, <https://mathoverflow.net/a/292978/1445>, 2018.
- [17] FELDMAN, J., M. SALO, and G. UHLMANN: The Calderón problem — an introduction to inverse problems. - Draft, <https://ims.nus.edu.sg/events/2018/theo/files/tutn3.pdf>.
- [18] GUO, C.-Y., M. KAR, and M. SALO: Inverse problems for p -Laplace type equations under monotonicity assumptions. - Rend. Istit. Mat. Univ. Trieste 48, 2016, 79–99.
- [19] HANNUKAINEN, A., N. HYVÖNEN, and L. MUSTONEN: An inverse boundary value problem for the p -Laplacian: a linearization approach. Inverse Problems 35:3, 2019, 034001.
- [20] HARJULEHTO, P., P. HÄSTÖ, Ú. V. LÊ, and M. NUORTIO: Overview of differential equations with non-standard growth. - Nonlinear Anal. 72:12, 2010, 4551–4574.
- [21] HERVAS, D., and Z. SUN: An inverse boundary value problem for quasilinear elliptic equations. - Comm. Partial Differential Equations 27:11-12, 2002, 2449–2490.

- [22] JOHNSON, W. P.: The curious history of Faà di Bruno's formula. - Amer. Math. Monthly 109:3, 2002, 217–234.
- [23] KANEIWA, R.: The formula for higher order derivatives of inverse functions. - Otaru University of Commerce Humanities Research 131, 2016, 1–3.
- [24] KAR, M., and J.-N. WANG: Size estimates for the weighted p -Laplace equation with one measurement. - Preprint, <http://www.math.ntu.edu.tw/~jnwang/pub/resources/papers/size0614.pdf>, 2018.
- [25] KUCHMENT, P., and D. STEINHAEUER: Stabilizing inverse problems by internal data. - Inverse Problems 28:8, 2012, 084007, 20.
- [26] KURYLEV, Y., M. LASSAS, L. OKSANEN, and G. UHLMANN: Inverse problem for Einstein-scalar field equations. - Preprint, arXiv:1406.4776, 2018.
- [27] KWON, O., E. J. WOO, J.-R. YOON, and J. K. SEO: Magnetic resonance electrical impedance tomography (MREIT): simulation study of J -substitution algorithm. - IEEE Trans. Biomedical Engineering 49:2, 2002, 160–167.
- [28] MEISS, J. D.: Differential dynamical systems. - Mathematical Modeling and Computation 22, Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, revised edition, 2017.
- [29] MUÑOZ, C., and G. UHLMANN: The Calderón problem for quasilinear elliptic equations. - Preprint, arXiv:1806.09586, 2018.
- [30] RUDIN, W.: Real and complex analysis. - Tata McGraw-Hill Education, 2006.
- [31] SALO, M., and X. ZHONG: An inverse problem for the p -Laplacian: Boundary determination. - SIAM J. Math. Anal. 44:4, 2012, 2474–2495.
- [32] SCHMÜDGEN, K.: The moment problem. - Grad. Texts in Math. 277, Springer, 2017.
- [33] STRAUB, D. I.: Numerical and microlocal analysis of inverse problems with internal data. - PhD thesis, University of Rochester, 2016.
- [34] SUN, Z.: On a quasilinear inverse boundary value problem. - Math. Z. 221:2, 1996, 293–305.