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Synthetic image generator for defocusing and astigmatic PIV/PTV

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This technical note introduces a synthetic image generator (SIG), referred to as MicroSIG, to be used for particle image/tracking velocimetry (PIV/PTV) analysis involving defocusing or astigmatic particle images. Target applications are experimental setups with volume illumination and small depth of field, in which the defocusing of particles plays a major role. This includes for instance μ PIV experiments or 3D PTV methods using defocusing or astigmatism to determine the out-of-plane particle position. The software uses an approximated model of spherical lens and conventional ray tracing to create realistic defocused/astigmatic particle images within a reasonable computational time. MicroSIG can be used to compare or optimize different PIV/PTV algorithms using arbitrary flows and different conditions including, non-monodispersed particles, non-uniform illumination, astigmatic optics, non-spherical particles (spheroids). A version of the software, available in Matlab and Python, can be found in the supplementary material and at this link: <https://gitlab.com/defocutracking>.

I. INTRODUCTION

Synthetic image generators (SIG) are commonly used in particle image/tracking velocimetry (PIV/PTV) research studies to create artificial sets of images that simulate experimental PIV/PTV recordings [1]. The obvious advantage of synthetic images is that the positions and displacements of particles are known *a priori* and different experimental conditions (such as seeding density, illumination, particle sizes) can be easily realized by changing few parameters in the code. Synthetic images have been extensively used in PIV/PTV research: from classical studies to determine the precision of the method and the optimal design rules [2–5], to more recent works assessing the resolution and accuracy of Tomographic-PIV [6] or single-pixel correlation methods [7], just to give few examples. An attempt to provide a standardized tool for generating synthetic images was the EUROPIV SIG [8], developed within the EUROPIV 2 project. Standardized synthetic images were used in the various editions of the PIV Challenge, a world-wide periodical initiative to assess the current state of the art of PIV techniques [9].

Most SIGs use 2D Gaussian functions to model particle images, which is the classical approach in PIV analysis [10]. This approach provides realistic results in case of in-focus particles with diameter of few pixels, which represents a large part of conventional and multi-camera PIV/PTV experiments. On the other hand, for methods using large magnification optics and volume illumination, such as μ PIV [11] or single-camera 3D methods [12, 13], it is necessary to model also particles which are out of focus. Several models can be used for this purpose, for instance using again the 2D Gaussian function approximation [14, 15], or the Huygens-Fresnel diffraction integral [16]. All these approaches, however, neglect the spherical aberration that has a significant effect on

these setups. Spherical aberration is a consequence of the spherical shape of real lenses which are in general not able to focus rays on a single spot (see Fig. 1).

This technical note introduces a SIG specifically designed to provide realistic images of in- and out-of-focus particles taken with large magnification optics. The SIG consists of a ready-to-use software, named MicroSIG, that is available in Matlab and Python. MicroSIG relies on the following assumptions:

- Particles are sphere or spheroids emitting light uniformly from their surface (like fluorescent particles).
- The objective lens is a spherical lens infinitely thin.
- Only geometric optics is used (diffraction is neglected).

MicroSIG can be used as a standardized tool to test and develop PIV and PTV methods dealing with defocused particle images. Additionally MicroSIG allows

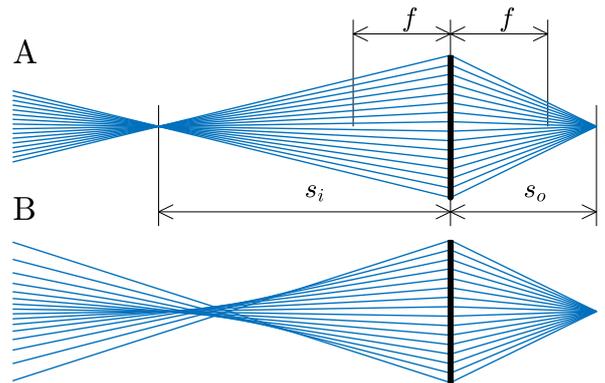


FIG. 1. (A) Ray tracing using a perfect thin lens: Rays originated from a point source at distance s_o from the lens converge to a single point at distance s_i , following the formula $1/s_i + 1/s_o = 1/f$, being f the focal lens of the lens. (B) Ray tracing using a spherical lens: Rays originated from a point source do not converge to a single point.

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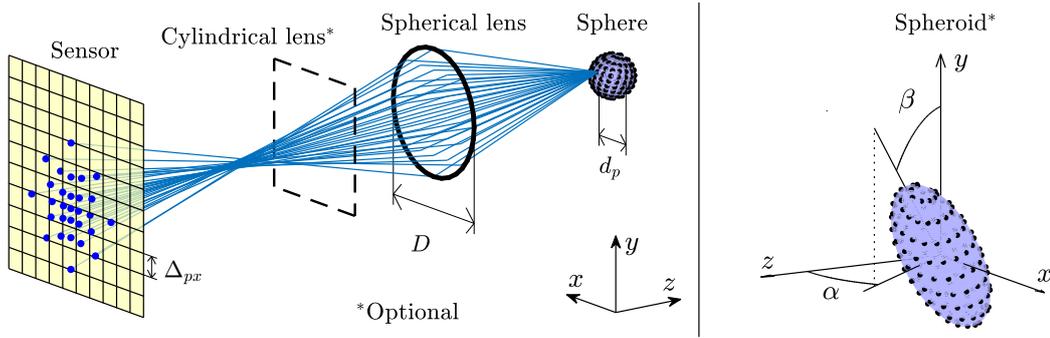


FIG. 2. General concept of MicroSIG. Tracer particles are spheres or spheroids, with light point-sources uniformly distributed on the surface. The path of each ray going through a spherical lens (and optionally also a cylindrical lens) is calculated. The intensity of one pixel is calculated as the count of rays hitting its area. A sketch of a spheroid particle with $e_p = 3$ and its respective orientation angles α and β is presented on the right panel.

to simulate spheroidal particles and astigmatic optics. Astigmatic imaging emerged as one robust method for 3D PTV [15, 17] and it has been successfully applied in many applications [18–21], however up to now a standardized SIG suitable for astigmatic particle images was not available yet.

II. IMPLEMENTATION

MicroSIG uses a three-dimensional ray-tracing approach, needed to deal with non-spherical particles and astigmatic lenses (systems that are not axisymmetric). The ray-tracing algorithm uses a matrix implementation, in which each ray crossing a plane perpendicular to the optical axis is described by a four-dimensional vector [22]:

$$\mathbf{r} = \begin{bmatrix} q_x \\ q_y \\ p_x \\ p_y \end{bmatrix} \quad (1)$$

where the vector $\mathbf{q} = (q_x, q_y)$ represents the position of the intersection point of the ray with the plane, and the vector $\mathbf{p} = (p_x, p_y)$ its orientation.

If the ray travels undisturbed, the ray vector \mathbf{r}' corresponding to a different plane can be calculated as:

$$\mathbf{r}' = \begin{bmatrix} q_x + tp_x \\ q_y + tp_y \\ p_x \\ p_y \end{bmatrix} = \begin{bmatrix} 1 & 0 & t & 0 \\ 0 & 1 & 0 & t \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} q_x \\ q_y \\ p_x \\ p_y \end{bmatrix} = \mathbf{T}\mathbf{r} \quad (2)$$

where t is the distance between the two planes. A lens can be considered as a function λ , depending on the focal length f , that modifies the orientation vector:

$$\mathbf{r}' = \begin{bmatrix} q_x \\ q_y \\ p_x' \\ p_y' \end{bmatrix} = \lambda(\mathbf{r}, f) \quad (3)$$

In MicroSIG, the function λ is constructed to simulate a infinite-thin, biconvex spherical lens. Rays are deflected using the Snell's law, first with respect to the air-glass interface, and then to the glass-air interface. The radius of curvature R of the interface is calculated from the lensmaker's equation for thin lens:

$$R = 2(n_l/n_m - 1)f \quad (4)$$

where n_l and n_m are the refractive index of the lens and of the lens immersion medium, respectively.

The first step for creating a particle image is to generate a discrete light field (i.e. a discrete ensemble of ray vectors) originated by a sphere (or spheroid) lying at a certain distance z from the lens. Practically, a number N_s of equally-spaced points are located on the particle surface, and for each point a number N_r of ray vectors with random orientation is created. In this way, the discrete light field \mathbf{r}_i is created, with $i = 1, 2, \dots, N$ and $N = N_s \cdot N_r$.

The light field is projected into the sensor plane with the following transformations:

$$\mathbf{r}_i' = \mathbf{T}_1 \mathbf{r}_i, \quad i = 1, 2, \dots, N \quad (5)$$

$$\mathbf{r}_j'' = \lambda(\mathbf{r}_j', f), \quad j \subseteq i \quad (6)$$

$$\mathbf{r}_k^{\text{ccd}} = \mathbf{T}_2 \mathbf{r}_k'', \quad k \subseteq j \quad (7)$$

where \mathbf{T}_1 is the translation matrix from the particle to the lens, and \mathbf{T}_2 from the lens to the sensor. The index j indicates the rays that hit the lens and are a subset of i , k indicates the rays that hit the sensor and are a subset of j . The sensor is composed by squared pixels with fill factor equal to 1. Finally, the intensity (counts) measured by a single pixels is calculated by the number of rays hitting its area.

If the astigmatic optics is activated, Eq. (7) is substituted by:

$$\mathbf{r}_k^{\text{ccd}} = \mathbf{T}_2 \mathbf{L}_{\text{cyl}} \mathbf{T}_{2\text{cyl}} \mathbf{r}_k'', \quad k \subseteq j \quad (8)$$

where \mathbf{L}_{cyl} is a linear transformation corresponding to a perfect cylindrical lens with focal length f_{cyl} :

$$\mathbf{L}_{cyl} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ -1/f_{cyl} & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad (9)$$

and \mathbf{T}_{2cyl} and \mathbf{T}_2 are translation matrices from the lens to the cylindrical lens, and from the cylindrical lens to the sensor. The astigmatic aberration created by the cylindrical lens can be adjusted either by modifying its position or its focal length [15]. To reduce complexity, it was decided to arbitrarily place the cylindrical lens at 2/3 the distance between the main lens and the sensor, and leave as free parameter only the focal length f_{cyl} .

III. THE MICROSIG SOFTWARE

The MicroSIG software takes as first input a *setting file*, which contains a list of 14 parameters defining the optics and the camera settings:

- *magnification*: Magnification of the optics.
- *numerical_aperture*: Numerical aperture of the optics.
- *focal_length*: Focal length of the virtual spherical lens.
- *ri_medium*: Refractive index of the immersion medium of the lens, typically air.
- *ri_lens*: Refractive index of the of the lens.
- *pixel_size*: Size of the pixels in μm .
- *pixel_dim_x*: Sensor width in pixels.
- *pixel_dim_y*: Sensor height in pixels.
- *background_mean*: Constant image background value.
- *background_noise*: Standard deviation of the Gaussian noise added to the images.
- *points_per_pixel*: Number of point sources per unit area, calculated as N_s divided by the theoretical area of an in-focus particle image (optimal values: 10–20).
- *n_rays*: Number of rays on each point source, corresponding to N_r (optimal values: 100–500).
- *gain*: Multiplicative factor to adjust the final image intensity.
- *cyl_focal_length*: Focal length of the virtual cylindrical lens. If set to 0, no astigmatism is present.

focal_length and *cyl_focal_length* are virtual parameters of the optical model and must be determined empirically. Recommended values, obtained from comparison with experimental images in the literature [13, 23] are 350 μm and 4000 μm , respectively. Examples are provided in the next section.

Subsequent inputs are the *data files*, one for each image to generate. The *data files* must be provided in ASCII format and are row arranged (one particle, one line).

MicroSIG allows four different input formats:

- x, y, z, d_p
- x, y, z, d_p, c_{int}
- $x, y, z, d_p, e_p, \alpha, \beta$
- $x, y, z, d_p, e_p, \alpha, \beta, c_{int}$

For each particle there are four mandatory variables: x, y, z and d_p . x and y are the particle's center coordinates in the image and are given in pixels as units. z is the out-of-plane component and is given in μm . $z = 0$ corresponds to the theoretical coordinate of the focal plane, which however may not correspond to the position where the particle looks more in focus due to the spherical aberration. d_p is the particle size (in μm).

To simulate spheroidal particles (i.e. ellipsoids of revolution), three additional parameters are needed: e_p , a factor which indicates the elongation along one axis of the particle (oblates for $e_p < 1$, prolates $e_p > 1$), α and β , which are the orientation angles as defined in Fig. 2. Finally, c_{int} is an optional multiplication factor to simulate light intensity inhomogeneities.

IV. RESULTS

Two examples of comparison between experimental and simulated particle image scans along the z -direction are reported in Fig. 3. In the first example in Fig. 3A, the experimental data are taken from [13] and correspond to a 2- μm -diameter particle observed with a microscope objective with magnification $M = 10\times$ and numerical aperture $\text{NA} = 0.3$. The scan is taken across a total height of 120 μm , with steps of 2 μm for a total of 61 images. The corresponding synthetic images were obtained using the setting parameters listed in Table I (Setting 1). The *data files* use the format $[x, y, z, d_p]$, with $d_p = 2$ μm and z ranging from -80 to 40 μm with steps of 2 μm . It should be noted that MicroSIG considers a particle at $z = 0$ always at the theoretical focal plane, and that this does not correspond exactly to the point of maximum intensity as a consequence of the spherical aberration (as shown in Fig. 3C). Also the choice of *focal_length* affects the rate of change of defocusing across the z -direction.

The second example is shown in Fig. 3B and corresponds to a 1- μm -diameter particle observed with a mi-

Parameter	Unit	Setting		
		1	1a	2
magnification	-	10	10	63
numerical_aperture	-	0.3	0.3	0.75
focal_length	μm	350	350	350
ri_medium	-	1	1	1
ri_lens	-	1.5	1.5	1.5
pixel_size	μm	6.45	6.45	6.45
pixel_dim_x	pixels	752	752	141
pixel_dim_y	pixels	366	366	141
background_mean	counts	350	350	52
background_noise	counts	10	10	5
points_per_pixel	1/pixels	18	18	12
n_rays	-	500	500	250
gain	-	3.2	3.2	1.1
cyl_focal_length	μm	0	4000	0

TABLE I. Values for the *setting files* used in the simulations in Section IV

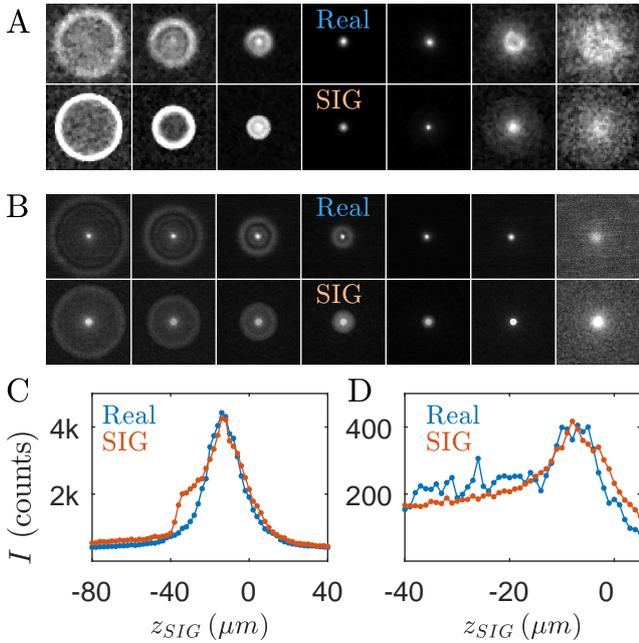


FIG. 3. Comparison between experimental and simulated particle image scans. (A) Real and synthetic images of 2- μm spheres taken with a 10 \times /0.3 objective lens. (B) Real and synthetic images of 1- μm spheres taken with a 63 \times /0.75 objective lens. (C) and (D) Corresponding maximum particle image intensity across the z direction.

lescope objective lens with $M = 63\times$ and $\text{NA} = 0.75$. The scan is taken across a total height of 45 μm , with steps of 1 μm for a total of 46 images. The experimental data are taken from [23]. The simulated images are taken using the setting parameters listed in Table I (Setting 2) and z ranging from -40 to 5 μm with steps of 1 μm . For both examples, a good qualitative and quantitative agreement between experimental and simulated images can be observed. Some finer structures, originated from the diffraction patterns, cannot clearly be recreated by MicroSIG, however the overall structure of the defocused images is well reproduced.

Further examples are shown in Fig. 4. Fig. 4A shows experimental and simulated images of 2- μm particles randomly distributed in a 100- μm -height microchannel. Fig. 4B shows the same case but this time using astigmatic optics. In both cases the experimental images are taken from [13] and more detail about the experimental setup can be found there. The synthetic images were obtained using the settings 1 and 1a in Table I. For the astigmatic case, it should be noted that the value $\text{cyl_focal_length} = 4000$ μm is a virtual parameter that cannot directly be related to the focal length of the cylindrical lens employed in the experiment (300 mm). Finally, Fig. 4C shows simulations of spheroidal particles with different positions and orientations. This feature allows to develop and test new algorithms able to recover the position and orientation of non-spherical objects, which are

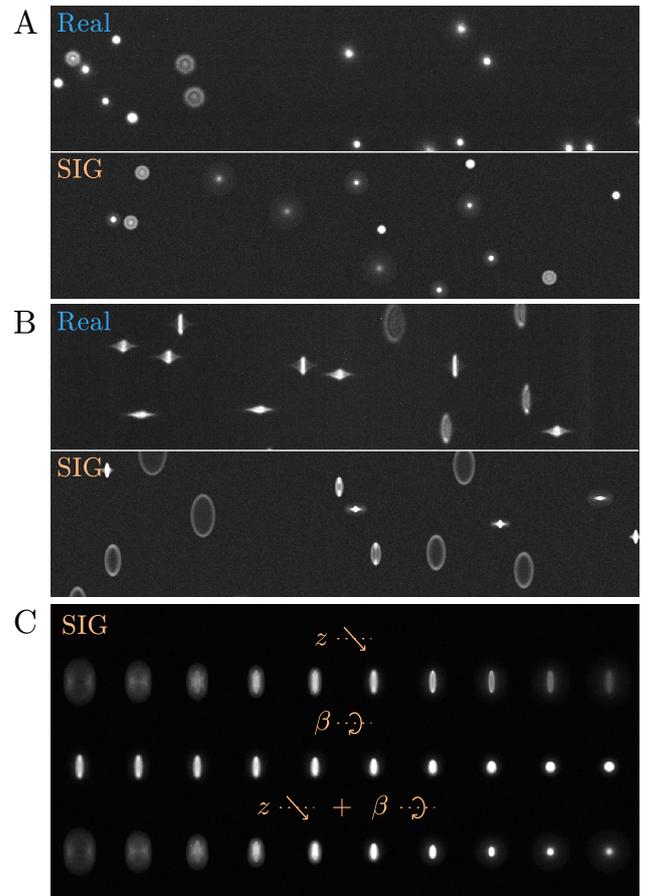


FIG. 4. (A) Real and synthetic images of 2- μm particles randomly distributed in a 100- μm -height microchannel observed with a 10 \times /0.3 objective lens, and (B) same setup with in addition an astigmatic optics. (C) Simulations of spheroidal particles with different positions and orientations.

relevant in many fields from cell measurements [24, 25] to anisotropic colloidal particles [26, 27].

V. CONCLUSIONS

This note introduces the main features of MicroSIG, a synthetic image generator (SIG) that provides realistic defocused and astigmatic particle images. MicroSIG can be used for the assessment of the uncertainty, under different conditions, of several PIV/PTV methods in which defocusing (or astigmatism) plays a role, including μPIV or 3D-PTV single-camera methods. Furthermore, it can be used to develop methods to measure position and orientation of non-spherical objects, which is currently an open field of research in colloids or biological flows. The MicroSIG software, available in Matlab and Python, is included in the supplementary material and can be downloaded at this link: <https://gitlab.com/defocustracking>.

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- [1] K. Okamoto, S. Nishio, T. Saga, and T. Kobayashi, *Measurement Science and Technology* **11**, 685 (2000).
 - [2] R. D. Keane and R. J. Adrian, *Measurement science and technology* **1**, 1202 (1990).
 - [3] R. D. Keane and R. J. Adrian, *Applied scientific research* **49**, 191 (1992).
 - [4] J. Westerweel, *Experiments in Fluids* **29**, S003 (2000).
 - [5] F. Scarano, *Measurement science and technology* **13**, R1 (2001).
 - [6] K. Lynch and F. Scarano, *Experiments in Fluids* **56**, 66 (2015).
 - [7] S. Scharnowski, R. Hain, and C. J. Kähler, *Experiments in fluids* **52**, 985 (2012).
 - [8] B. Lecordier and J. Westerweel, in *Particle image velocimetry: recent improvements* (Springer, 2004) pp. 145–161.
 - [9] C. J. Kähler, T. Astarita, P. P. Vlachos, J. Sakakibara, R. Hain, S. Discetti, R. La Foy, and C. Cierpka, *Experiments in Fluids* **57**, 97 (2016).
 - [10] R. J. Adrian and C.-S. Yao, *Applied optics* **24**, 44 (1985).
 - [11] J. G. Santiago, S. T. Wereley, C. D. Meinhart, D. Beebe, and R. J. Adrian, *Experiments in fluids* **25**, 316 (1998).
 - [12] C. Cierpka and C. Kähler, *Journal of visualization* **15**, 1 (2012).
 - [13] R. Barnkob, C. J. Kähler, and M. Rossi, *Lab on a Chip* **15**, 3556 (2015).
 - [14] M. Olsen and R. Adrian, *Experiments in fluids* **29**, S166 (2000).
 - [15] M. Rossi and C. J. Kähler, *Experiments in fluids* **55**, 1809 (2014).
 - [16] T. Fuchs, R. Hain, and C. J. Kähler, *Optics letters* **39**, 1298 (2014).
 - [17] C. Cierpka, R. Segura, R. Hain, and C. J. Kähler, *Measurement Science and Technology* **21**, 045401 (2010).
 - [18] W. Qiu, J. T. Karlsen, H. Bruus, and P. Augustsson, *Physical Review Applied* **11**, 024018 (2019).
 - [19] M. Rossi, A. Marin, N. Cevheri, C. J. Kähler, and M. Yoda, *Microfluidics and nanofluidics* **23**, 67 (2019).
 - [20] F. Kiebert, S. Wege, J. Massing, J. König, C. Cierpka, R. Weser, and H. Schmidt, *Lab on a Chip* **17**, 2104 (2017).
 - [21] R. Barnkob, N. Nama, L. Ren, T. J. Huang, F. Costanzo, and C. J. Kähler, *Physical Review Applied* **9**, 014027 (2018).
 - [22] A. Gerrard and J. M. Burch, *Introduction to matrix methods in optics* (Courier Corporation, 1994).
 - [23] M. Rossi, R. Segura, C. Cierpka, and C. J. Kähler, *Experiments in fluids* **52**, 1063 (2012).
 - [24] K. Taute, S. Gude, S. Tans, and T. Shimizu, *Nature communications* **6**, 8776 (2015).
 - [25] J. S. Guasto, R. Rusconi, and R. Stocker, *Annual Review of Fluid Mechanics* **44**, 373 (2012).
 - [26] J.-C. Loudet, A. M. Alsayed, J. Zhang, and A. G. Yodh, *Physical review letters* **94**, 018301 (2005).
 - [27] M. Cavallaro, L. Botto, E. P. Lewandowski, M. Wang, and K. J. Stebe, *Proceedings of the National Academy of Sciences* **108**, 20923 (2011).