



NQR Investigation of Phase Transitions in Cesium Plumbochloride

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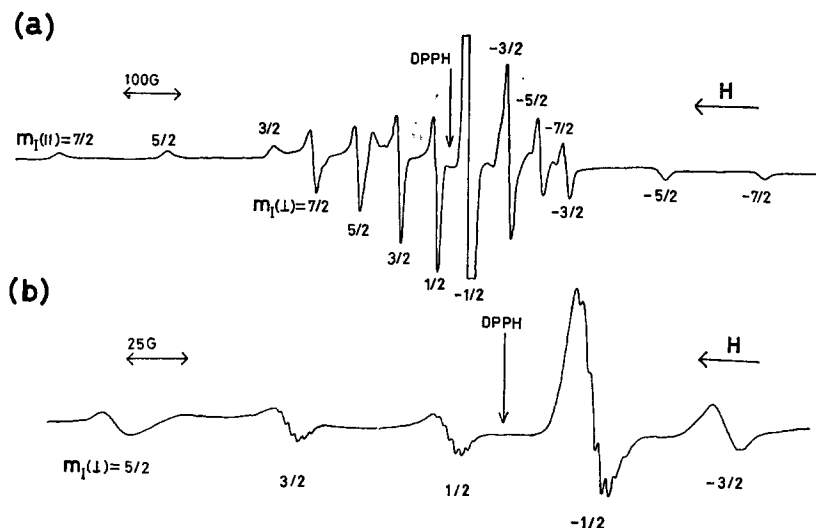


FIG. 1. Electron spin resonance spectra of VOPC in conc H_2SO_4 matrix (1.75×10^{-3} mol/liter) at 77°K measured by JEOL p-10 type spectrometer (100 kc/sec modulation) at X band. Field modulation is 3.0 G wide for (a) and 0.7 G for (b).

Ingram and Bennett¹ on polycrystalline specimens, and more recently by Assour *et al.*² on magnetically diluted VOPC samples, respectively. The latter authors, showing vanadium nuclear hyperfine (hf) structure but with the absence of the nitrogen superhyperfine (shf) lines, suggested that the unpaired electron in the vanadium $3d_{xy}$ orbital is localized on the metal core, and that there is no *in-plane* π bonding between the metal and the ligand-nitrogen atoms.

We have studied the ESR spectra of VOPC in dilute glassy matrix of conc sulfuric acid, α -Cl-naphthalene, or of quinoline at 77°K and found contrary to the above authors' results that the spectra give rise to nitrogen shf structure in all solvents investigated. Thus, the unpaired electron is believed to interact with the ligand-nitrogen nuclei as in the case of vanadyl porphyrin.³

The ESR spectrum obtained with large width for the field modulation is shown in Fig. 1(a). Such a spectrum may be explained on the basis of the following spin Hamiltonian with axial symmetry:

$$\mathcal{H} = \beta [g_{\parallel} H_x S_x + g_{\perp} (H_x S_x + H_y S_y)] + A_{\parallel} I_x S_x + A_{\perp} (I_x S_x + I_y S_y) \quad (S = \frac{1}{2}, I = \frac{7}{2}), \quad (1)$$

where β is the Bohr magneton and A is the vanadium nuclear hf interaction constant. After Sands⁴ and Bleaney,⁵ the spin-Hamiltonian parameters were calculated as

$$\begin{aligned} g_{\parallel} &= 1.968 \pm 0.002, & g_{\perp} &= 1.987 \pm 0.001, \\ A_{\parallel} &= -(161.2 \pm 2.0) \times 10^{-4} \text{ cm}^{-1}, \\ A_{\perp} &= -(57.3 \pm 1.0) \times 10^{-4} \text{ cm}^{-1}. \end{aligned}$$

When a small width for the field modulation such as 0.7 G was applied, some of the perpendicular hf lines ($m_I = \frac{3}{2}, \frac{1}{2}, -\frac{1}{2}$) were found to split into at least seven lines spaced with separations of 2.5–3.0 G [Fig. 1(b)]. Such splittings were not observed, however, on the other perpendicular lines ($m_I = \frac{7}{2}, \frac{5}{2}, -\frac{3}{2}, -\frac{5}{2}, -\frac{7}{2}$) nor on any parallel lines. Splitting of this sort may be

interpreted qualitatively if the unpaired electron of the vanadium interacts with the ligand-nitrogen nuclei.

The nitrogen shf structure can be explained by adding the following term to Eq. (1)

$$\sum [A_{\parallel}^N S_x I_z^N + A_{\perp}^N (S_x I_x^N + S_y I_y^N)],$$

where I^N refers to nitrogen nuclear spin ($I=1$) and A^N to nitrogen shf constant. Accordingly, if the four nitrogen ligands are magnetically equivalent, nine shf lines should be observed with relative intensities of

$$1:4:10:16:19:16:10:4:1.$$

The observed shf structure is consistent with this expectation, if two weakest lines are assumed to be smeared out. On the other hand, it is not obvious why not all the vanadium hf lines give rise to nitrogen shf structure.

The spin density around the nitrogen ligands of VOPC is expected to arise from the delocalization of in-plane π bonding or from the configuration interaction of such orbital as in-plane σ bonding.³ To discuss the bonding scheme along the line accurate data for A_{\parallel}^N are at least required. Unfortunately we were unable to detect A_{\parallel}^N .

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NQR Investigation of Phase Transitions in Cesium Plumbobichloride

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As has been demonstrated by Møller,¹⁻² cesium plumbobichloride (CsPbCl_3) has the cubic perovskite

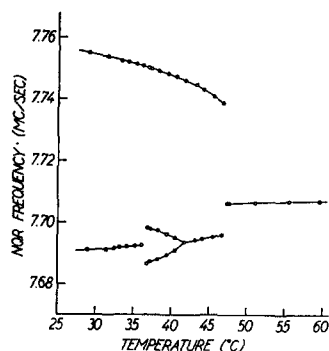


FIG. 1. The ^{35}Cl NQR frequency in CsPbCl_3 as a function of temperature. The linewidth is about 1.5 kc/sec.

structure above 47°C and a slightly distorted tetragonal structure below this temperature. Furthermore it is suggested¹ that the cesium atoms and the chlorine atoms are displaced from the ideal positions expected for the perovskite structure. A Cl^{35} pure quadrupole resonance experiment has been carried out on CsPbCl_3 powder to get further information on this structure, the results being shown in Fig. 1.

The NQR experiment indicates clearly that the phase transition at 47°C is a first-order transition. At 42°C the low-frequency line, the intensity of which is about twice as high as that of the other line, splits into two lines; this indicates a phase transition of the second order which has not been observed previously. At 37°C the splitting of the low-frequency line disappears through a first-order transition. No hysteresis has been detected in any of the transitions.

Above 47°C all chlorine atoms are placed at the same field gradient corresponding to crystallographically identical positions. Between 47° and 42°C the crystal is tetragonal¹ with two crystallographically nonidentical chlorine positions; the low-intensity line originates in the chlorine atoms placed at the fourfold axes and the high-intensity one in the chlorine atoms placed at the twofold axes. As the main contribution to the field gradient originates in the Pb ions, the relative positions of the two lines in the diagram are the opposite of what might be expected for the simplest possible transition from the cubic phase to the tetragonal one (i.e., an elongation of the c axis and a simultaneous contraction of the two axes perpendicular to the c axis). This indicates that some rearrangement of the chlorine atoms must take place during the phase transition at 47°C .¹ Between 42° and 37°C , three lines of approximately the same intensity are observed corresponding to three different chlorine positions. It should be noted that the two transitions at 42° and 37°C take place in such a way that the chlorine atoms placed at the fourfold axes obviously are unaffected.

Taking into account the two nearest Pb^{2+} ions, the four nearest Cs^+ ions, and the eight nearest Cl^- ions, the experimentally determined Sternheimer antishielding factor is -23 . The experiment thus indicates that the quadrupole coupling is entirely ionic.

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Raman Spectra from Dilute Solutions of HDO in D_2O

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Raman spectra from a $6.2M$ (11 mole%) solution of D_2O in H_2O were reported recently.¹ The OD stretching contour from the $6.2M$ solution was observed to be markedly asymmetric at 25°C . A high-frequency shoulder was evident near $2660 \pm 20 \text{ cm}^{-1}$ which was separated by an inflection from the contour maximum at $2515 \pm 5 \text{ cm}^{-1}$. Inflections and shoulders have also been observed in the OD stretching contour from a $3.1M$ (6 mole%) solution of D_2O in H_2O and in the OH stretching contours from $6.5M$ (12 mole%) and $2.8M$ (5 mole%) solutions of H_2O in D_2O .² Some evidence for weak components arising from intermolecular coupling between HDO vibrations was obtained, however, and thus it seemed desirable to investigate solutions more dilute than 5 mole%. Accordingly, Raman spectra from a $0.56M$ (1 mole%) solution of H_2O in D_2O were obtained in this work.³

Photoelectric Raman spectra from the $0.56M$ or 1 mole% solution were obtained in the OH stretching region with a Spectrophysics 140 argon-ion-laser having maximum intensity approaching 1 W at 4880 \AA . A Cary Model 81 Raman spectrophotometer modified to utilize an ITT FW 130 (S-20) phototube was used with the laser, and a 25- μl iter capillary Raman cell was employed with the Cary image slicer. A typical spectrum is shown in Fig. 1.

Examination of Fig. 1 (upper section) indicates pronounced asymmetry in the OH stretching contour from the dilute HDO solution. A shoulder is evident near $3630 \pm 20 \text{ cm}^{-1}$, and it is separated from the contour maximum by an inflection near $3600 \pm 20 \text{ cm}^{-1}$.

Detailed analysis of the OH contour was accomplished by determining the base line from extrapolation of curvatures above 3700 cm^{-1} and below 3100 cm^{-1} (see the dashed line in Fig. 1). (The slope and curvature resulted from proximity to the extremely intense OD stretching contour, and from other sources, e.g., from the capillary cell, and from the required high amplification.) The spectral contour was then transferred to a horizontal base line (lower section of Fig. 1), where a pronounced high-frequency shoulder is evident. The quantity $|dI/d\nu|$ from the transferred curve, the absolute value of the derivative of contour intensity with respect to frequency shift in cm^{-1} , was found to be minimal near $3600 \pm 20 \text{ cm}^{-1}$, and zero near $3435 \pm 10 \text{ cm}^{-1}$. Analog decomposition of the transferred curve by means of the duPont 310 curve Resolver also revealed two major Gaussian components near 3435 and 3628 cm^{-1} , indicated in Fig. 1 by peak height and peak width.

The shape of the OH stretching contour indicated