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Exciton dephasing and biexciton binding in CdSe/ZnSe islands

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The dephasing of excitons and the formation of biexcitons in self-organized CdSe/ZnSe islands grown by molecular-beam epitaxy is investigated using spectrally resolved four-wave mixing. A distribution of exciton-exciton scattering efficiencies and dephasing times in the range of 0.5–10 ps are observed. This indicates the presence of differently localized exciton states at comparable transition energies. Polarization-dependent measurements identify the formation of biexcitons with a biexciton binding energy of more than four times the bulk value. With decreasing exciton energy, the binding energy slightly increases from 21.5 to 23 meV, while its broadening decreases from 5.5 to 3 meV. This is attributed to a strong three-dimensional confinement with improving shape uniformity for decreasing exciton energy. [S0163-1829(99)04739-6]

Self-organized semiconductor islands providing three-dimensional quantum confinement presently attract attention because of their expected superiority for applications in opto-electronic and all-optical devices, where improvements such as a reduction of the laser threshold and higher optical nonlinearities due to the modified density of states are predicted and partly already achieved.^{1–3} The optical properties of such islands have been mainly studied by incoherent methods.^{4–7} Driven by the search for blue-green-emitting semiconductor materials, the fabrication^{8–13} and optical characterization^{14–18} of self-assembled II-VI islands have recently been the subject of intensive research. Only little is known about the coherent properties of these structures, giving information about exciton scattering and biexciton formation. These effects are particularly important in II-VI materials because of their large biexciton binding energy, which strongly influences their stimulated emission. Four-wave mixing (FWM) provides a powerful tool to study the fundamental coherent processes in low-dimensional II-VI semiconductor systems, as already shown in quasi-two-dimensional^{19,20} and quasi-one-dimensional structures.²¹

In this paper we study CdSe/ZnSe island structures by femtosecond transient FWM. The excitation-density-dependent dephasing of island excitons and the biexciton binding energy are investigated. The studied CdSe/ZnSe island structures were grown by molecular-beam epitaxy (MBE) on a GaAs(001) substrate. Details about the growth and the characterization are given in Ref. 10. The capped ZnSe island structures were characterized by plan-view transmission electron microscopy (TEM) and high resolution TEM, showing an island density of $5 \times 10^{10} \text{ cm}^{-2}$ and a lateral island size of about $30 \times 15 \text{ nm}^2$ elongated along $[1\bar{1}0]$. Quantitative analysis of the Cd distribution by a digital

analysis of lattice images²² shows $\text{Zn}_{1-x}\text{Cd}_x\text{Se}$ islands of about 9 ML height with an inhomogeneous Cd content close to $x=1$ in the island center.

In the FWM experiments, the sample was excited by the signal wave of an optical parametric amplifier pumped by a regenerative Ti-sapphire amplifier, providing 90-fs pulses of 50 nJ pulse energy at a repetition rate of 300 kHz. The pulses are chirp compensated and spectrally shaped to a duration between 120 fs and 3 ps in a pulse shaper. The FWM experiments were performed in reflection geometry with two incident pulses of the directions $\mathbf{k}_{1,2}$, which were colinear ($\uparrow\uparrow$) or cross-linear ($\uparrow\rightarrow$) polarized. The focus diameter of the pulses at the $1/e^2$ intensity on the sample was 200 μm . The FWM signal in the reflected $2\mathbf{k}_2 - \mathbf{k}_1$ direction was detected time-integrated and spectrally resolved by a combination of a spectrometer and an optical multichannel analyzer as a function of the time delay τ between \mathbf{k}_1 and \mathbf{k}_2 . Typical FWM efficiencies were 10^{-5} . The sample was kept in a helium bath cryostat at a temperature of 5 K.

Microphotoluminescence (PL) measurements (Fig. 1) on the investigated sample reveal a structured band, indicating the presence of individual transitions from localized states with energies distributed around 2.3 eV. The strong blue-shift compared to the band gap of 1.752 eV in bulk zincblende CdSe (Ref. 23) is due to interdiffusion, strain, and confinement effects as discussed in Refs. 14–18. The FWM experiments are performed using different center wavelengths λ of the exciting pulses, indicated by arrows in Fig. 1.

The FWM intensity for ($\uparrow\uparrow$) configuration (Fig. 2) shows a maximum at $\tau > 0$ and a fast quenching for $\tau < 0$, typical for the photon-echo response of an inhomogeneously broadened system. For $\tau > 0$ a nonexponential decay, slowing

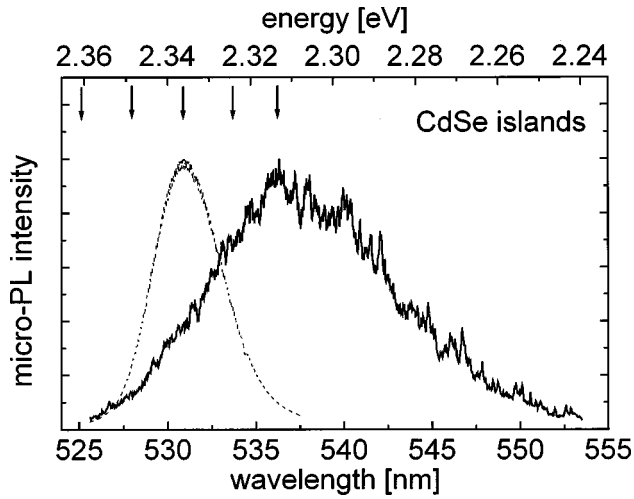


FIG. 1. Micro-PL spectra with $0.5\text{-}\mu\text{m}$ resolution of the CdSe quantum island sample at 10 K excited with $\lambda=514\text{ nm}$. The dashed line is one spectrum of the FWM excitation pulses. The different center positions are indicated by arrows.

down with increasing decay time, is observed. This reveals a distribution of dephasing times within the excited exciton ensemble, which we attribute to different exciton localization. The FWM signal intensity decreases strongly with decreasing excitation photon energy, a consequence of the decreasing density of states. The smallest excitation photon energy feasible in the FWM experiment corresponds to the maximum of the PL (see Fig. 1). Using a double-exponential fit to the data (thick full line in Fig. 2) we extract two components of different dephasing times T_2 . The homogeneous linewidth $2\hbar/T_2$ of the fast component versus excitation intensity is given in the inset of Fig. 2 for different transition energies. The intensity dependence as well as the zero-density extrapolated value decreases with decreasing transition energy. Additionally, the weight of the fast component relative to the slow one is decreasing. From these observa-

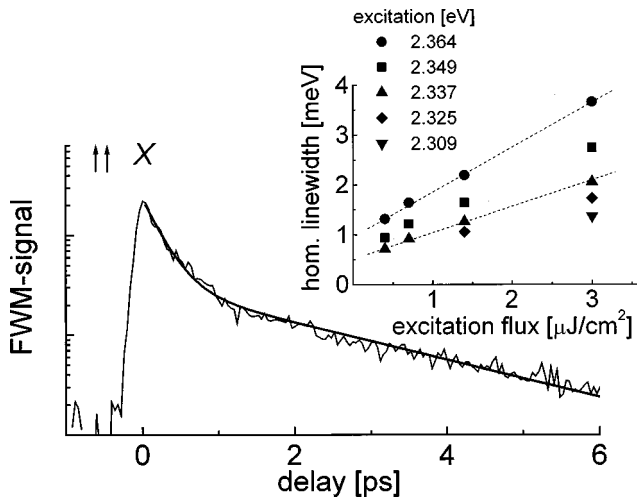


FIG. 2. FWM intensity as a function of the delay time τ in the $(\uparrow\uparrow)$ configuration, for an excitation energy of 2.325 eV and a total \mathbf{k}_1 and \mathbf{k}_2 pulse intensity of $1.4\text{ }\mu\text{J}/\text{cm}^2$. The inset shows the homogeneous linewidth (full width at half maximum) of the fast component as a function of the pulse intensity for different excitation energies, as labeled.

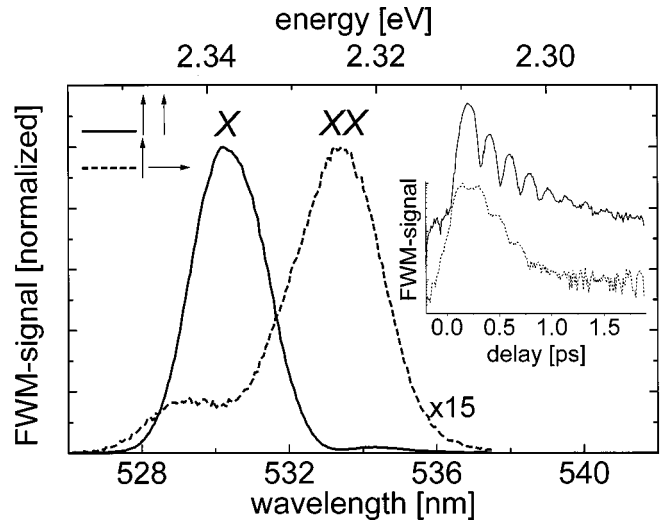


FIG. 3. FWM spectra in the $(\uparrow\uparrow)$ and $(\uparrow\rightarrow)$ configurations at a delay time of 0.4 ps , an excitation energy of 2.337 eV , and a total pulse intensity of $3\text{ }\mu\text{J}/\text{cm}^2$. The $(\uparrow\rightarrow)$ data are magnified by a factor of 15. The inset shows the FWM spectral intensity versus delay time at the position of the biexciton-induced transitions XX for the $(\uparrow\uparrow)$ and $(\uparrow\rightarrow)$ configurations.

tions, we attribute the fast component to more delocalized exciton states, which show more efficient exciton-exciton scattering and in addition phonon-assisted relaxation into lower-lying states. With decreasing transition energy, the fraction of more delocalized exciton states decreases. This is in agreement with the position of the PL maximum, that indicates the transition energy below which radiative recombination is faster than phonon-assisted relaxation. The slow decaying component is accordingly attributed to these more strongly localized excitons. Their homogeneous linewidth decreases slowly with decreasing transition energy ranging from 0.22 to 0.13 meV and shows only a weak density dependence.

The formation of biexcitons is identified by polarization-dependent FWM measurements, for which the selection rules can be derived from four-level²⁴ or five-level²⁵ Bloch equations. Figure 3 shows FWM spectra for $(\uparrow\uparrow)$ and $(\uparrow\rightarrow)$ polarization configurations. In the $(\uparrow\uparrow)$ configuration the FWM signal is dominated by the ground-state to exciton ($0\text{-}X$) transition (labeled X). For the $(\uparrow\rightarrow)$ configuration the signal mainly originates from the exciton to biexciton ($X\text{-}XX$) transition (labeled XX), which is spectrally separated from the ($0\text{-}X$) transition by the biexciton binding energy, and is collinear polarized to \mathbf{k}_1 . The strongly suppressed signal at the ($0\text{-}X$) transition shows that exciton to unbound biexciton transitions are weak in these island structures, which is related to the quantization of the biexciton continuum.²⁶ The FWM traces at the spectral position of the XX transitions for $(\uparrow\uparrow)$ and $(\uparrow\rightarrow)$ configuration are given in the inset of Fig. 3. Pronounced oscillations with a period of $T=185\text{ fs}$ arising from the interference between exciton- and biexciton-induced polarizations are observed, showing the expected phase shift of π between $(\uparrow\uparrow)$ and $(\uparrow\rightarrow)$ configurations.²⁷ The beats are rather weak in the $(\uparrow\rightarrow)$ configuration due to the small contribution of the unbound biexcitons. The beating period corresponds to a biexciton binding energy of $E_{XX}=22.4\pm 0.3\text{ meV}$, which is by far larger than the bulk

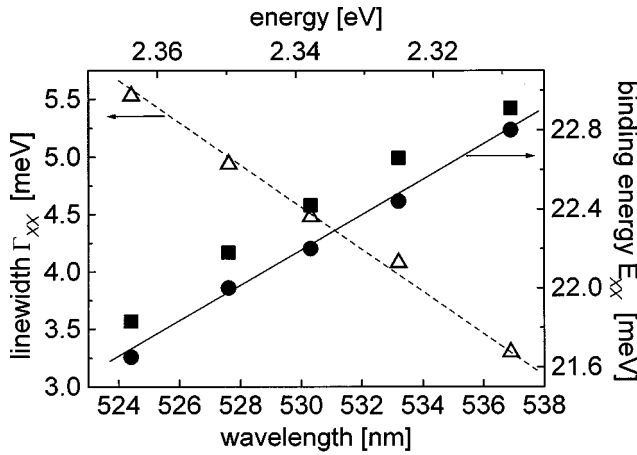


FIG. 4. Biexciton binding energy E_{XX} (full symbols) and inhomogeneous broadening Γ_{XX} (open triangles) of the biexciton-induced XX signal as a function of the exciton transition energy. The circles give the biexciton binding energy E_{XX} determined from the FWM spectra performed with spectrally narrow \mathbf{k}_1 pulses; the squares show the values obtained from the FWM beat period.

CdSe biexciton binding energy of about 4 meV,¹⁸ indicating a strong three-dimensional quantum confinement in the islands. This is in agreement with experiments on similar structures showing biexciton binding energies of about 20 meV (Refs. 17 and 18) and also to theoretical and experimental results for spherical CdSe quantum dots in glasses.^{28,29} We observe a decrease of the oscillation period and thus an increase of the biexciton binding energy of about 5% with decreasing exciton transition energy, plotted as filled squares in Fig. 4. The observed increase of the biexciton binding energy is attributed to the stronger spatial localization in the energetically deeper localized states, similar to what has been found in other localized systems.³⁰

A direct determination of the biexciton binding-energy distribution from the FWM spectra can be achieved by selecting a narrow distribution of $(0-X)$ transitions with a spectrally narrow \mathbf{k}_1 pulse.³¹ FWM spectra for the $(\uparrow\rightarrow)$ configuration and a \mathbf{k}_1 spectral width of 1 meV are given in Fig. 5, where the scattered light intensity from the \mathbf{k}_1 pulse indicates the energy of the selected $(0-X)$ transition. The FWM spectra were normalized by the spectral intensity of the spectrally broad \mathbf{k}_2 pulse in order to isolate the spectral shape of the biexcitonic FWM response. The extracted biexciton binding energies (filled circles in Fig. 4) show a slight increase with decreasing exciton transition energy, similar to the data extracted from the beat period. Furthermore, the $(X-XX)$ transition exhibits an additional broadening, which is decreasing from 5.5 to 3.3 meV with decreasing $(0-X)$ transition energy (Fig. 5). Such an additional broadening of the biexciton binding energy has been observed in localized quasi-two-dimensional systems^{26,31} and is explained as an inhomogeneous broadening due to a nonuniform shape of the localization potentials for excitons even at comparable transition energies. The observed decrease of the broadening (Fig. 4) is accordingly interpreted as improved shape uniformity of the localized states, which are less influenced by the surrounding potential landscape due to their smaller extension. This is consistent with the observed increase of the biexciton binding energy, also indicating stronger localization.

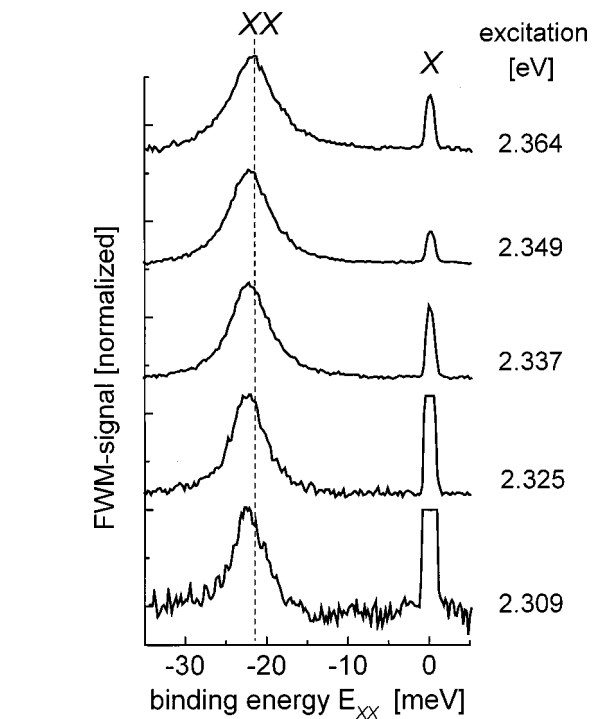


FIG. 5. FWM spectra at $\tau=0.4$ ps using a spectrally narrow \mathbf{k}_1 pulse ($\Delta E=1$ meV) obtained at various excitation wavelengths for the $(\uparrow\rightarrow)$ configuration. The energetic distance between the \mathbf{k}_1 pulse and the biexciton-induced XX -signal directly gives the biexciton binding energy.

ity of strongly localized states, which are less influenced by the surrounding potential landscape due to their smaller extension. This is consistent with the observed increase of the biexciton binding energy, also indicating stronger localization.

In summary, the low-temperature dephasing of excitons in the investigated self-organized ZnSe capped CdSe island structures shows a distribution of dephasing times due to differently localized excitons. The faster decaying components, which show a large exciton-exciton scattering efficiency, are attributed to weakly localized excitons that can relax by phonon emission into more strongly localized states. The slower components, on the other hand, are attributed to these more strongly localized states, and show a weaker exciton-exciton scattering efficiency. The localized biexcitons in the structures have binding energies of 22 meV, more than four times the bulk value, indicating a strong three-dimensional confinement. Since biexcitons with such high binding energies are stable at room temperature they are important for applications in all-optical or opto-electronic devices. With decreasing transition energy, the biexciton binding energy is slightly increasing and its inhomogeneous broadening is decreasing, which is attributed to shrinking size and increasing shape uniformity of the localized states.

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