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Vankó, György ; Bordage, Amélie ; Glatzel, Pieter ; Gallo, Erik ; Rovezzi, Mauro ; Gawelda, Wojciech ; Galler, Andreas ; Bressler, Christian ; Doumy, Gilles ; March, Anne Marie

Total number of authors:
22

Published in:
Journal of Electron Spectroscopy and Related Phenomena

Link to article, DOI:
[10.1016/j.elspec.2012.09.012](https://doi.org/10.1016/j.elspec.2012.09.012)

Publication date:
2013

Document Version
Early version, also known as pre-print

[Link back to DTU Orbit](#)

Citation (APA):

Vankó, G., Bordage, A., Glatzel, P., Gallo, E., Rovezzi, M., Gawelda, W., Galler, A., Bressler, C., Doumy, G., March, A. M., Kanter, E. P., Young, L., Southworth, S. H., Canton, S. E., Uhlig, J., Smolentsev, G., Sundström, V., Haldrup, K., Brandt van Driel, T., ... Lemke, H. T. (2013). Spin-state studies with XES and RIXS: From static to ultrafast. *Journal of Electron Spectroscopy and Related Phenomena*, 188, 166-171.
<https://doi.org/10.1016/j.elspec.2012.09.012>

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Spin-state studies with XES and RIXS: From static to ultrafast

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Abstract

We report on extending hard X-ray emission spectroscopy (XES) together with resonant inelastic X-ray scattering (RIXS) to study ultrafast phenomena in a pump-probe scheme at MHz repetition rates. The investigated system consists of low-spin (LS) Fe^{II} complex compounds, where optical pulses induce a spin-state transition to their \approx nanosecond-lived high-spin (HS) state. Time-resolved XES clearly reflects the spin-state variations with very high signal-to-noise ratio, in agreement with HS–LS difference spectra measured via thermally induced spin crossover, and reference HS–LS systems in static experiments. 1s2p RIXS, measured at the Fe 1s pre-edge region shows variations after laser excitation, which are consistent with the formation of the HS state. Our results demonstrate that now X-ray spectroscopy experiments with overall rather weak signals, such as RIXS, can be reliably exploited to study chemical and physical transformations on ultrafast time scales.

Keywords: molecular switching, time-resolved spectroscopy, X-ray spectroscopy, spin transition, photoinduced transition, ultrafast phenomena, pump-probe experiments

1. Introduction

The spin state of transition metal ions affects the properties of their compounds both on macroscopic

and microscopic levels, as it has a large influence on ionic and thus molecular size, lattice unit cell volume, density, transport, magnetic and vibrational properties. Particularly interesting are cases when the spin state can be changed by external stimuli, with implications in atomic, molecular and solid state physics, geophysics, chemistry, and biochemistry. Molecular systems that can be switched between their different spin states are especially intriguing for their potential applications in high-density molecular storage devices, and the photochemical properties may prove to become useful in novel light energy converters based on Fe-containing molecules. At low temperatures many low-spin (LS) Fe^{II} complex compounds can be switched by light to a high-spin (HS) state [1, 2]; however, the detailed mechanism of the LS to HS switching process is still under debate [3]. Tools sensitive to the electronic structure and the spin state, which can also be utilized on the femtosecond (fs) to nanosecond (ns) time scales, can shed light on the yet unknown details of the relevant intermediate steps involved in the spin-state switching process. Previously, fs laser spectroscopy has been used for this purpose, which have delivered a rich picture of the ultrafast events taking place on the first few hundred fs [4, 5, 6, 7, 8]. These tools are, however, rather insensitive to the changes in the spin state next to the concomitant geometric structure changes.

Hard X-ray spectroscopies can provide unique insights into the electronic and geometric structure, combined with element-, orbital-, and orientation sensitivity [9, 10, 11]. These techniques probe the bulk properties of the sample, and are even compatible with extreme conditions due to the large penetration depths [12]. X-ray Absorption Spectroscopy (XAS) with its variants X-ray Absorption Near Edge Structure (XANES) and Extended X-ray Absorption Fine Structure (EXAFS) have proven their utility in determining the unoccupied electron states and the local structure, respectively [13, 14, 15, 16]. X-ray Emission Spectroscopy (XES) and Resonant Inelastic X-ray Scattering (RIXS) are very powerful in elucidating the electronic structure, providing information on the electron energies, local geometry, spin and valence state, etc. [9, 17, 18, 19, 20, 21, 22, 23]. These tools can thus deliver complementary information to

the usually employed ultrafast laser spectroscopies, when implemented in a pump-probe experiment with a laser pump and a time-delayed X-ray probe pulse.

The wide-spread extension of hard X-ray techniques to the ultrafast time domain has been initiated at synchrotron radiation sources about a decade ago, and has remained limited in its time resolution to ca. 50-100 ps, due to the intrinsic X-ray pulse width in this range. In nearly all these cases, such studies on molecular systems have employed X-ray absorption spectroscopy [24, 25, 26] or X-ray scattering (XRS) techniques [25, 27, 28, 29, 30, 31, 32] only and have been, with few exceptions, limited to low repetition rates (ca. 1 kHz) by using Ti:sapphire-based amplified laser systems, which deliver a sufficient number of optical photons to match the number of ground state molecules in the laser beam [24]. In this way, the photoexcitation of a sufficient fraction of the sample (5-50 %) was ensured for each laser pulse. This was necessary, since the number of probing X-ray photons/pulse are orders of magnitude lower (and their MHz rate cannot be exploited), which so far hampered a wide-spread use of such structural X-ray techniques. More flux demanding experiments, such as X-ray emission spectroscopy (XES) or even the very low cross section inelastic X-ray scattering techniques (IXS), thus become nearly impossible to envision. However, a recent pioneering experiment successfully recorded transient XES spectra on photoexcited aqueous $[\text{Fe}(\text{bipy})_3]^{2+}$ at the Fe $K\alpha$ emission line, which reflected a line shape variation characteristic for the LS to HS transition [33].

The above works also showed that the integrated number of X-ray photons is the bottleneck for these time-resolved X-ray techniques, and in consequence we have developed a setup employing a high-power MHz laser, which allows for a substantial increase in detection efficiency due to a closer match between the repetition rates of both the synchrotron and the laser source. This strategy now permits ca. 10^{12} X-ray photons per data point to be collected in a few seconds, which is at least 3 (4?) orders of magnitude larger than with conventional kHz laser systems after extensive data acquisition times of several days. With this approach one can now envision to employ multiple X-ray structural tools simultaneously on the

sample during the course of a light-triggered chemical reaction or physical transition.

2. Experimental methods

We have extended the capabilities of such laser-X-ray studies towards MHz pump-probe repetition rates, both at ID26 of the European synchrotron radiation facility (ESRF) and at sector 7-ID of the Advanced Photon source (APS). These high-energy synchrotron radiation sources offer the advantage to tailor the laser pump rate to that of all electron bunches within the storage ring, so that one can perform such studies with quasi-static S/N.

Both beamlines use undulator radiation, which is then monochromatized using a double crystal monochromator equipped with either diamond (APS) or silicon (ESRF) crystals. At 7-ID beamline, the X-rays are then focussed down to $5\ \mu\text{m}$ FWHM spot size at the sample by Kirkpatrick-Baez mirrors [34], whereas at ID26 the beam size is tailored using a toroidal (?) mirror and further downstream slit system to a spot of $100 \times 500\ \mu\text{m}^2$. Time-resolved experiments are most conveniently performed at both synchrotrons in special equidistant bunch filling modes, namely in the 24-bunch (APS) and 16-bunch modes (ESRF). In the 24-bunch mode at a ring current of 102 mA and for the X-ray energies between 7.5 keV and 8.2 keV, the X-ray incident flux on 7-ID is ca. 2×10^5 photons that are delivered to the sample in the ~ 80 ps FWHM X-ray pulse at a repetition rate of 6.52 MHz. In 16-bunch mode (90 mA, 5.68 MHz), $\approx 2 \times 10^6$ photons are delivered to the sample in ~ 100 ps FWHM for each X-ray pulse on ID26.

Aqueous solution of $[\text{Fe}(\text{bipy})_3]^{2+}$ (with concentrations between 12 mM and 25 mM) in form of a $100\ \mu\text{m}$ thick sheet from a liquid jet was used as sample in both experiments. From the APS experiment, we also report here time-resolved 1s2p RIXS taken on a 20 mM solution of $[\text{Fe}(\text{terpy})_2]^{2+}$ (terpy:2,2':6',2''-terpyridine). The laser installed in the APS 7-ID-D hutch [35] was operated at a 3.26 MHz repetition rate, emitting 532 nm radiation with 5 W power corresponding to pulse energies of $2\ \mu\text{J}$; the pulse width was 10 ps, and was focussed to a $25\ \mu\text{m}$ FWHM spot size on the sample. Spatial and temporal stability

were $5\ \mu\text{m}$ and 10 ps over several hours [35]. At ID26, we used a mobile fs laser system described by the following parameters: wavelength 515 nm, repetition rate 1.42 MHz, pulse length 280 fs. Cylindrical lenses were used to achieve an elliptical beam profile of ca. $125 \times 600\ \mu\text{m}^2$, which provided a good overlap with the X-ray spot size on the sample.

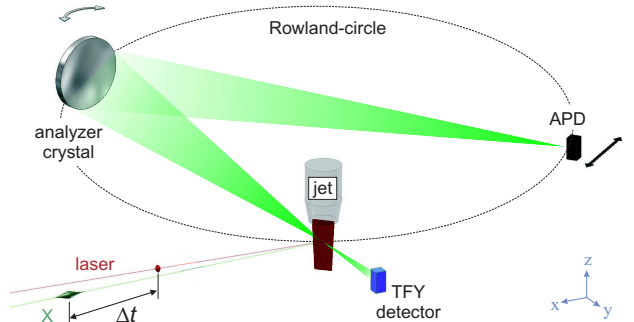


Figure 1: Schematic setup at APS 7-ID; see the text for details.

Figure 1 shows the schematic experimental setup. The laser and X-ray beams impinge almost collinearly (5 degrees angle of separation) on a fast-flowing $100\ \mu\text{m}$ thick liquid jet mounted at 45° to the beam propagation direction. A wavelength-dispersive crystal spectrometer was constructed at 7-ID (in the horizontal plane) to measure the XES signal using a spherically bent analyzer crystal of 10 cm diameter positioned at a 90° scattering angle, and an avalanche photodiode (APD) detector at its focus, such that the sample-analyzer-detector setup constitutes a 1m Rowland circle. At ID26, the permanent five-crystal spectrometer of the beamline, built in the vertical plane, was used. A thin plastic bag filled with He was installed inside both spectrometers to minimize absorption loss of the XES intensity. The spectra of the emitted X-rays were recorded by scanning the Ge(440) spherically bent analyzer crystal through the appropriate Bragg angle range to collect the Fe $K\alpha$ spectra, and the APD was moved accordingly. Total energy resolution measured through the spectrometer on 7-ID (resp. ID26) was 1.0 eV FWHM at 6404 eV (resp. ??? at ??? eV), a resolution composed of 0.4 eV (resp. ??? eV) and 0.9 eV (resp.

???? eV) contributions from the monochromator and the spectrometer, respectively.

The relative timing between the laser and X-ray pulses is controlled with an electronic phase shifter and it allows choosing a particular time delay Δt , at which we recorded XES spectra or RIXS planes. We used a data acquisition scheme where two spectra are collected alternately, following the laser ON/OFF pattern (at APS every second X-ray pulse was pumped with the laser, whereas at ESRF every 4th X-ray pulse), to acquire data on the ground and laser-excited states. This permitted us to optimally exploit the incident X-ray flux by matching the repetition rates of both pump and probe sources, similarly to our previously employed strategy with kHz systems [7, 36].

3. Results and discussion

3.1. The high-spin state of $[\text{Fe}(\text{bipy})_3]^{2+}$ as reflected by XES

In an X-ray absorption study, Gawelda *et al.* determined the structure of the light-induced ≈ 0.6 ns-lived excited state of $[\text{Fe}(\text{bipy})_3]^{2+}$ and showed that the structural variations correspond to a 0.2 \AA elongation of the Fe–N bond length [8], typical for a LS to HS transition in Fe^{2+} systems [37]. In our recent extension towards ps time-resolved XES (using a low-repetition pump laser at 1 kHz), we characterized the spin momentum change of the Fe^{II} , which was in a good agreement with the expected $\Delta S = 2$ [33]. In the experiments reported here, the MHz pump-probe rates have allowed us to collect very high quality data at acquisition times similar to conventional static experiments. Figure 2 displays the transient $\text{K}\alpha$ signals obtained on the aqueous solution of the $[\text{Fe}(\text{bipy})_3]^{2+}$ molecules, with additional HS–LS reference spectra. The latter includes the theoretical difference spectrum using the charge transfer multiplet approach (Fig. 2c), as well as difference in static experiments: the difference of the room temperature $\text{K}\alpha$ spectra of HS $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ and LS $[\text{Fe}(\text{bipy})_3]^{2+}$ (Fig. 2d), and the difference measured at thermal spin crossover in $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$, well above and below the transition temperature (Fig. 2e).

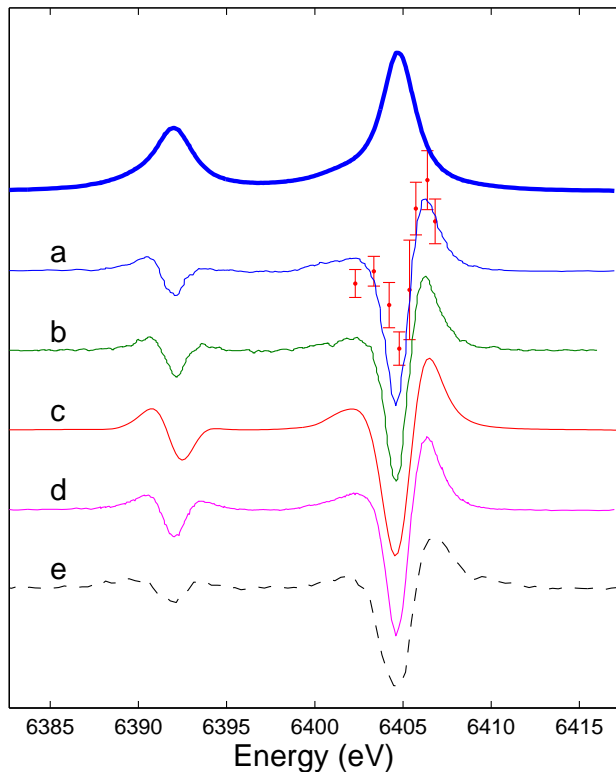


Figure 2: $\text{K}\alpha$ spectra of the aqueous solution of (LS) $[\text{Fe}(\text{bipy})_3]^{2+}$ (top thick blue curve, scaled down) and several HS–LS difference spectra. **a**: $[\text{Fe}(\text{bipy})_3]^{2+}$ laser on – laser off transient from APS 7-ID (taken at a time delay of $\Delta t = 80$ ps); the circles with error bars show our previous kHz experiment at the SLS[33] for comparison. **b**: $[\text{Fe}(\text{bipy})_3]^{2+}$ transient spectrum from ESRF ID26 ($\Delta t = 77$ ps). **c**: Difference curve of the $\text{K}\alpha$ spectra calculated with the charge transfer multiplet approach for HS and LS Fe^{2+} in O_h symmetry, from Ref. [33]. **d**: difference of the room temperature $\text{K}\alpha$ spectra of the HS $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ and the LS $[\text{Fe}(\text{bipy})_3]^{2+}$. **e**: static experiment at the thermal spin crossover of $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$, the difference of the $\text{K}\alpha$ spectra taken at the room temperature (HS form) and at 80 K (LS form) is shown.

The excellent agreement between the time-resolved and the static experiments, as well as with the theory, demonstrates that XES can reliably determine the 3d spin momentum (or, more precisely, the number of unpaired electrons), and can be used as a powerful tool for ultrafast studies.

3.2. Time-resolved XES: Validation of the methodology for line shape analysis

In conventional XES experiments aiming at the determination of the spin state, energy shifts and even line shape distortions can hinder the quantitative analysis. This is particularly true for a horizontal spectrometer arrangement, when avoiding these effects would require keeping the sample in the same position along the X-ray beam very precisely. In the same way, in cryostats, pressure cells, etc., it is often difficult to ensure that the sample remains in the same position when the temperature or the pressure is changed. Therefore, in many previous experiments the spectra were aligned on the main peak [38, 39, 40]. In order to circumvent this problem, in Ref. [18] we suggested aligning the spectra on their center of mass for cases when energy shifts can take place. A further typical assumption at the line shape analysis is that the fluorescence yield is not affected by the transition, so the XES intensity at a fluorescence feature is identical for both states; however, it is known that the $K\beta : K\alpha$ intensity ratio can show some chemical dependence [41]. Since the IAD approach is used more and more widely [18, 42, 43, 44, 45, 46], it is important to test the validity of these assumptions.

The alternating acquisition scheme of the present pump-probe experiment provides us with a unique opportunity to do so, since the two spectra are taken quasi-simultaneously. Therefore the sample position, the geometrical arrangement and all other parameters are identical for the two spectra in the laser ON/OFF (or pumped/unpumped) channels, that probe the LS ground state and the (partly) HS excited state, respectively. For $[\text{Fe}(\text{bipy})_3]^{2+}$, the $K\alpha$ spectra were of excellent statistical quality: 20 pairs of laser ON/OFF scans, with a total of 39×10^6 counts were taken, in two sets. We compare the difference of their center of mass ($\Delta\text{COM}=\text{COM}(\text{HS})-\text{COM}(\text{LS})$) values, as well as the possible intensity variation for the laser ON/OFF channels, with the ON values scaled to 100 % conversion.

For the two sets, taken at average step size of 125 meV and 280 meV, ΔCOM values of $-5 (\pm 22)$ meV and $4 (\pm 18)$ meV were determined. However, the ΔCOM values of the individual scans scattered between ± 1 step size. Moreover, 1-2 orders of magni-

tude smaller errors are expected from counting statistics with error propagation, suggesting that the spectra are affected by some small instabilities during the pump-probe experiment. Nevertheless, the COM seems unperturbed and thus can be used for aligning the spectra at line shape analysis. In fact, the $K\alpha$ difference spectrum of $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ is only similar to the others because the HS and LS spectra (taken from Ref. [18]) were now shifted to the same COM before making the difference. The such-treated spectrum displays a good similarity with the others, only it is broader due to the large beam size used in that previous experiment.

The differences of integrated intensity in the laser ON —normalized to the X-ray beam intensity monitor, and scaled to 100% HS conversion— and laser OFF channels is $\Delta I = -0.4 \pm 0.3\%$, when compared to the laser OFF channel, for the first set of scans, and it is $-1.6 \% \pm 0.9 \%$ for the other. In certain scans evident crosstalk between the two channels were observed affecting a not more than one data point; however, a smaller electronics problem of this kind might have gone unnoticed. The scans on the $[\text{Fe}(\text{terpy})_2]^{2+}$ result in $\Delta I = -0.9 \pm 0.3\%$, all these indicate a small $K\alpha$ yield loss of around $-0.7 \% \pm 0.2 \%$ at 100 % HS conversion. In the same experiment, we also measured the $K\beta$ spectra of the laser excited $[\text{Fe}(\text{bipy})_3]^{2+}$, and an effect of $\Delta I = +0.5 \pm 0.4\%$ was observed (scaled to 100 % spin conversion). We do not wish to conclude here on the possible effect on the electronic structure; however, we can claim that these small possible intensity variations will not undermine the accuracy of the spin-state determination.

3.3. Time-resolved RIXS at the spin-state transition

Resonant inelastic X-ray scattering (RIXS) can also give valuable insights into the electronic structure [10, 9, 11]. For spin-state transitions the variations of the 3d state energies and populations are substantial. Consequently, the RIXS planes taken at the 1s pre-edge resonances, which (at least partly) stem from $1s \rightarrow 3d$ transitions, are expected to show large variations at spin state changes. This has been demonstrated for the thermal SCO in $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ in a static experiment, again well below and above the transition temperature [18]. We

replot these figures in Fig. 3, revealing that while for the LS case there is a single resonance, there are three of them for the HS state (ignoring splittings from the weak S-O interaction for the 3d electrons). These resonances, while they are less accessible in XANES due to the lifetime broadening and the overlap of the peaks, clearly show up in the RIXS planes. Their difference is also shown; however, this old set of spectra are again of limited resolution, and a 2D splining was necessary to produce the difference. Nevertheless, it should be sufficient to lay down the topological expectations on what we should expect for a time-resolved difference spectrum for a Fe^{II} spin-state transition. The lower row of Fig. 3 presents the 1s2p RIXS data for the light-induced spin-state transition in [Fe(terpy)₂]²⁺ taken at 7-ID. The changes of the affected 3d states are quite clear from this figure: the intensity of the single peak of the ground state drops, and two additional peaks appear, indicating a more complex electron configuration following laser excitation (HS population was determined to 40% here). The difference spectra and the HS state spectrum reveal a structure with 3 resonances. This finding agrees very well with what is expected to be observed for the quadrupolar 1s3d resonances at LS→HS transition in an octahedral Fe(II) complex, and is in good agreement with what was observed for the pseudooctahedral [Fe(phen)₂(NCS)₂] in static measurements at a thermal spin-state transition [18, 47]. One should keep in mind, however, that the [Fe(terpy)₂]²⁺ has a pronounced distortion from the *O_h* symmetry, and lacks an inversion center; therefore, mixing of 3d and 4p states is expected, and a relevant dipole contribution to the pre-edge might appear. Therefore, the complete understanding of the spectra requires an elaborate theoretical analysis, which is beyond the scope of the present paper.

4. Conclusions

We presented here a successful transition from static to dynamic XES and RIXS to investigate low-spin (LS) Fe^{II} complex compounds and their photoinduced spin-transition. The excellent agreement between the results of the different experiments enabled to validate the assumptions previously made for line

shape analysis during XES static experiments. In addition, the results demonstrated the power and reliability of using an X-ray spectrometer in combination with ultrafast lasers. With the applied conditions, ps-resolved resonant hard X-ray spectroscopy has practically reached its full potential at synchrotron radiation sources. These successful pump-probe experiments open up novel exciting opportunities to study many physical, chemical, or biological processes on the ultrafast time scale.

5. Acknowledgement

This project was supported by the European Research Council via contract ERC-StG-259709, by the German research community (DFG) via contract SFB925, the European CRISP project, by the European XFEL, and by the Hungarian Scientific Research Fund (OTKA) under contract No. K 72597. A.M.M., G.D., S.H.S., E.P.K., and L.Y. acknowledge support from the U.S. Department of Energy (DOE) Office of Science, Division of Chemical, Geological and Biological Sciences under Contract No. DE-AC02-06CH11357. Use of the Advanced Photon Source, an Office of Science User Facility operated for DOE Office of Science by Argonne National Laboratory, was supported by the U.S. DOE under Contract No. DE-AC02-06CH11357. We acknowledge the assistance of Pablo Fajardo, Christian Hervé, Christophe Lapras, and the staff of ID6 and ID9 from the ESRF as well as the staff of 7-ID from the APS.

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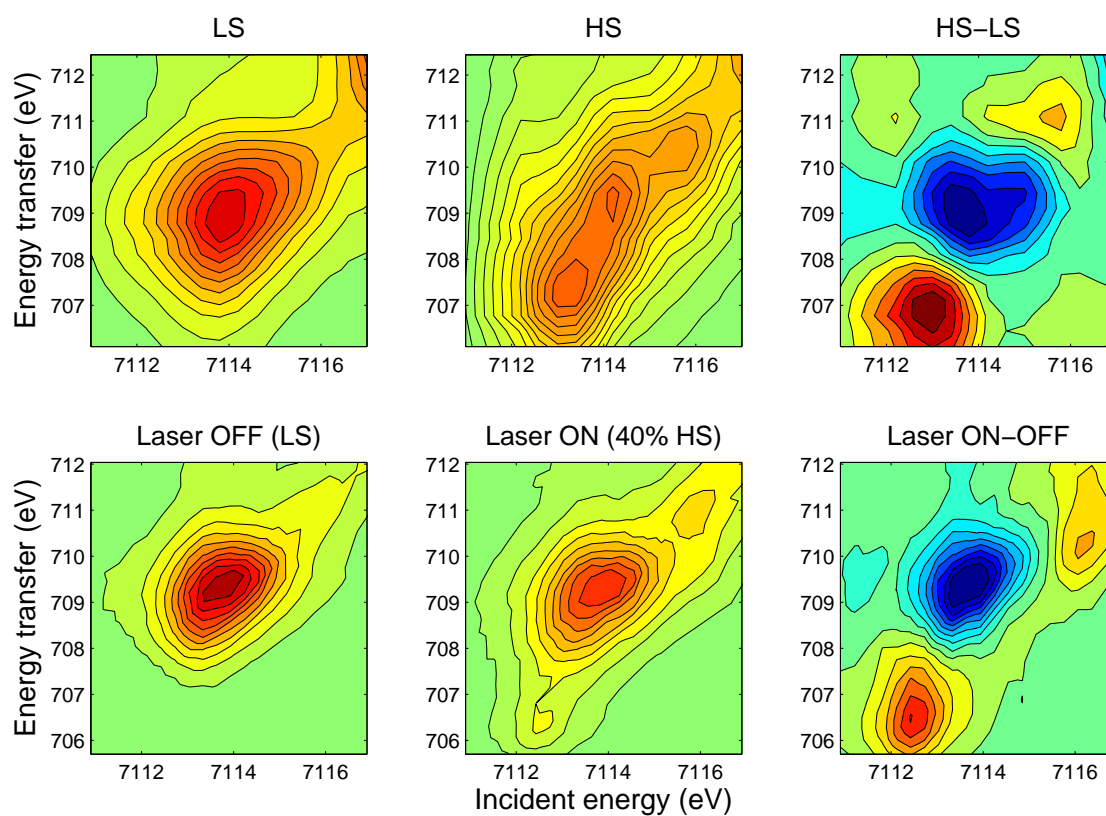


Figure 3: $1s2p_{3/2}$ RIXS of $[\text{Fe}(\text{phen})_2(\text{NCS})_2]$ at thermal spin crossover (upper row), and $[\text{Fe}(\text{terpy})_2]^{2+}$ in aqueous solution at laser excitation (lower row).