



## **A particle simulation code for analysis of nonlinear electron oscillations in a magnetized plasma waveguide**

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<p>Title and author(s)</p> <p>A Particle Simulation Code for Analysis of Nonlinear Electron Oscillations in a Plasma Waveguide</p> <p>by</p> <p>V.A. Turikov</p>	<p>Date August 1978</p> <hr/> <p>Department or group</p> <p>Physics</p> <hr/> <p>Group's own registration number(s)</p>
<p>21 pages + tables + 5 illustrations</p>	
<p>Abstract</p> <p>A description is given of a computer code for simulation of electron oscillations in a magnetized plasma in a cylindrical waveguide. The one-dimensional particle-in-cell method with the reverse linear interpolation of charge density is used. The program has options for treating nonlinear processes in a plasma with periodical and reflecting boundary conditions. For periodical conditions, Poisson's equation is solved by means of the Fourier method. For reflecting conditions, the double recursive procedure is used. The values of the potential derivatives at the space grid points are calculated by means of the parabolic interpolation. The main purpose of the program is to investigate nonlinear phenomena in a plasma column after applying a short localized impulse of an external electric field.</p>	<p>Copies to</p>
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## 1. Introduction

Strongly nonlinear processes in a collisionless plasma are being intensively investigated at the present time both experimentally and theoretically. The Q-machine is a very suitable device for treating such general plasma properties. In recent experiments<sup>1),2)</sup> in a single-ended Q-machine with a magnetized plasma in a cylindrical waveguide, formation of various kinds of stable pulses due to localized external excitations was observed. These pulses have been identified as solitons, rarefactive waves and phase-space holes similar to the BGR-modes<sup>3)</sup>. It is very difficult to investigate such kinds of problem theoretically because one has to solve the nonlinear system of Vlasov and Poisson equations with a strong external perturbation.

The simulation code SIMULA has developed to investigate nonlinear electron oscillations in a magnetized cylindrical plasma column to analyze the experimental results from the Q-machine plasma. Computer calculations performed with this code<sup>2),4)</sup> have confirmed the fact that stable positive pulses of the potential propagating in a plasma after an external impulse in axial direction appear as vortexes or "holes" in phase space, like the vortexes forming in the two-stream instability<sup>5),6)</sup>. Also some important properties of solitons and rarefactive pulses have been studied by means of these computer simulations.

The particle-in-cell simulation scheme<sup>7),8)</sup> with the reverse linear interpolation method<sup>5)</sup> of charge density calculation is used in the simulation code. The program is written in FORTRAN and runs have been made on a Burroughs 6700 at the Risø Computer Installation. The purpose of this paper is to describe this simulation code.

## 2. Physical Assumptions and Equations

The computer program described here simulates the properties of a magnetized plasma in a cylindrical waveguide. We assume infinite axial magnetic field, therefore electrons can move only in the axial direction (x-direction), and we can consider the problems to be one-dimensional. The ions are considered as a stationary positive background.

Space-charge waves (azimuthally symmetric modes) in a cylindrical waveguide have the maximum phase velocity for small wavenumbers<sup>9)</sup>

$$v_{ph} = \frac{\omega_{pe}}{k_1}, \quad (1)$$

where  $k_1 = 2.404/r_0$ ,  $r_0$  is the radius of the plasma column, and  $\omega_{pe}$  is the electron plasma frequency.

For the azimuthally symmetric modes, the potential can be calculated from Poisson's equation

$$\frac{d^2\phi}{dx^2} - k_1^2\phi = -4\pi\rho. \quad (2)$$

We assume boundary conditions with a zero electric field at the ends of the plasma column

$$\left. \frac{d\phi}{dx} \right|_{x=0} = \left. \frac{d\phi}{dx} \right|_{x=L} = 0, \quad (3)$$

where  $L$  is the column length.

The electric field energy can be calculated using Poisson's equation (2)

$$W_E = \frac{\pi r_0^2}{8\pi} \int (E_x^2 + k_1^2\phi^2) dx, \quad (4)$$

where  $E_x$  is the electric field in the plasma.

We assume that the localized impulse of the external potential can be presented in the form:

$$e\phi_{ext}(x,t) = W_{ph} \cdot \eta(x)\sigma(t), \quad (5)$$

where the functions  $\eta(x)$ ,  $\sigma(t)$  describe the space and time variations of the external potential, and  $W_{ph}$  is the kinetic energy of the electron with the phase velocity (1). The concrete profiles of  $\eta(x)$  and  $\sigma(t)$  used in the simulation are presented in section 5.9.

### 3. General Algorithm of the Simulation Scheme

The particle-in-cell method (PIC method) is employed in the simulation code. In this method it is convenient to use dimensionless variables defined by the following relations<sup>8)</sup>

$$x = \frac{x}{\Delta}, \quad v = v \frac{2\Delta t}{\Delta},$$

$$E = -E_x \frac{4(\Delta t)^2}{\Delta} \frac{q}{m_0}, \quad F = \phi \frac{2(\Delta t)^2}{\Delta^2} \frac{q}{m_0}, \quad (6)$$

$$Q = -G \frac{\rho}{en_0}, \quad G = 2u_{pe}^2 (\Delta t)^2,$$

where  $\Delta$  is the space step,  $\Delta t$  is the time step, and  $n_0$  is the initial plasma density.

The leap-frog scheme<sup>8)</sup> is used for moving all simulation particles the time step  $2\Delta t$ :

$$v_m^{k+1} = v_m^{k-1} + E^k(x_m^k) + E_{\text{ext}}^k(x_m^k), \quad (7)$$

$$x_m^{k+2} = x_m^k + v_m^{k+1}, \quad m = 1, 2, \dots, N,$$

where  $k$  is the time step number,  $N$  is the total number of simulation particles, and  $E_{\text{ext}}(x)$  is the external electric field.

The first time step is realized by the Euler method<sup>8)</sup>

$$x_m^1 = x_m^0 + v_m^0/2. \quad (8)$$

After the calculation of the coordinates of all particles at moment  $t^k = k\Delta t$  we can obtain the values of the charge density  $Q_i^k$  in the nodes of the space grid. The charge density is calculated by distributing the charge of each particle between the two nearest grid points according to a reverse linear interpolation<sup>5)</sup>

$$Q_i^k = G (N_i(t^k)/N_c - 1), \quad (9)$$

$$N_i(t^k) = \sum_{m=1}^N \delta(\text{Int}|x_m^k - i|) \cdot (1 - |x_m^k - i|),$$

where  $N_i(t^k)$  is the number of charges in  $i$  units in grid point  $i$  at time moment  $t^k$ ,  $N_c$  is the initial number of particles in each cell,  $\text{Int}(z)$  is a truncation from real value to integer and

$$\delta(m) = \begin{cases} 1, & m = 0 \\ 0, & m \neq 0. \end{cases}$$

Having the charge density distribution  $Q_i^k$ , we can calculate the new values of the self-consistent electric field in eqs. (7) used to move all particles the next time step. For the electric field calculation, Poisson's equation is solved with the boundary conditions (3) (or with periodical conditions) and charge density values  $Q_i^k$  using the three-point difference scheme

$$F_{i+1}^k - (2 - k_{\perp}^2 \Delta^2) F_i^k + F_{i-1}^k = Q_i^k. \quad (10)$$

Equations (7)-(10) define the general time loop of the PIC method used in the simulation code.

If a simulation particle is outside the interval  $0 < X < J$  ( $J = L/\Delta$ ) in some time step and reflecting boundary conditions (3) are used, the coordinate and the velocity of this particle are altered according to:

$$x_{m, \text{refl}}^k = \begin{cases} -x_m^k, & x_m^k < 0, \\ 2J - x_m^k, & x_m^k > J, \end{cases}$$

$$v_{m, \text{refl}}^k = -v_m^k.$$

When the periodical conditions are used, the coordinates and velocities of such particles are changed in the following way:

$$x_{m, \text{per}}^k = \begin{cases} J + x_m^k, & x_m^k < 0, \\ x_m^k - J, & x_m^k > J, \end{cases}$$

$$v_{m, \text{per}}^k = v_m^k.$$

We assume the initial distribution of the simulation particles is homogeneous in x-space and random in v-space with Gaussian statistics (see sec. 5.7)

The electric field energy (4) and the total plasma energy calculated per one particle are computed during the simulated process

$$W_E = \frac{1}{2GJ} \int_0^J [(E + E_{\text{ext}})^2 + 4k_{\perp}^2 \Delta^2 F^2] dx, \quad (11)$$

$$W_{\text{total}} = W_E + \frac{1}{N} \sum_{m=1}^N v_m^2.$$



We calculate the electric field in an arbitrary point  $X$  by means of the expression:

$$E(X) = 2 \frac{dF}{dX} = 2(2A_j X + B_j),$$

where  $A_j$  and  $B_j$  are coefficients of the quadratic interpolation of the potential between the three nearest space grid points

$$F(X) = A_j X^2 + B_j X + C_j, \quad j = \text{Int}(X) + 1.$$

For the calculation of the integral  $\int_0^J F^2 dX$  in formula (11), we use parabolic interpolation for the values  $F_i^2$ .

It appeared during the simulations that the total energy conserves after the action of the external field with an error of about 1-2%, which means that the model can be considered quite accurate.

The general flow of the program is presented in Figs. 1a, 1b, and concrete features of the main program SIMULA are discussed in sec. 5.

#### 4. Algorithm for Solving Poisson's Equation

There are two options for solving Poisson's equation in the simulation code - for periodical and nonperiodical conditions.

General nonperiodical boundary conditions can be written in the form:

$$(A_1 \frac{dF}{dX} + B_1 F) |_{X=0} = C_1, \quad (12)$$

$$(A_r \frac{dF}{dX} + B_r F) |_{X=J} = C_r, \quad (13)$$

where  $A_1$ ,  $B_1$ ,  $C_1$ ,  $A_r$ ,  $B_r$ , and  $C_r$  are prescribed constants.

For solving Poisson's equation with boundary conditions (12) and (13), the double recursive procedure is employed<sup>8)</sup> with the parabolic interpolation of the potential for calculating  $\frac{dF}{dX}$  at the ends of the plasma column:

$$F_j = c_j F_{j-1} + s_j, \quad j = 1, 2, \dots, J,$$

where the coefficients  $c_j$  and  $s_j$  are also defined by the recurrence relations:

$$c_{j-1} = \frac{1}{c_{j-k_1^2 \Delta^2 - 2}},$$

$$s_{j-1} = \frac{Q_{j-1}^{-s_j}}{c_{j-k_1^2 \Delta^2 - 2}}, \quad j = J, J-1, \dots, 2.$$

The value  $F_0$  can be found using the boundary condition (12) for the left boundary of the interval

$$F_0 = \frac{C_1 + A_1 (c_2 s_1^{-s_2 - 4s_1})/2}{A_1 (2c_1 - c_1 c_2 / 2 - 3/2) + B_1}.$$

For the periodical boundary conditions, we use the Fourier method for solving the Poisson equation (8), (10):

$$F_j = \sum_{i=1}^{J_2} (\hat{P}_i^c \cos \frac{2\pi i j}{J} + \hat{P}_i^s \sin \frac{2\pi i j}{J}), \quad (14)$$

where  $J_2 = (J-1)/2$  for odd  $J$  and  $J_2 = J/2$  for even  $J$ , and  $\hat{P}_i^c, \hat{P}_i^s$  are the cosine and sine Fourier coefficients. These coefficients can be obtained by means of a Fourier analysis of the Poisson equation (10):

$$\hat{P}_i^c = \frac{\hat{Q}_i^c}{\lambda_i}, \quad \hat{P}_i^s = \frac{\hat{Q}_i^s}{\lambda_i}, \quad (15)$$

$$\lambda_i = 2 \left( \cos \frac{2\pi i}{J} - 1 - \frac{k_1^2 \Delta^2}{2} \right), \quad i = 1, 2, \dots, J_2,$$

where  $\hat{Q}_i^c, \hat{Q}_i^s$  are the cosine and sine Fourier coefficients of the charge density

$$\hat{Q}_i^c = \frac{2}{J} \sum_{m=1}^J Q_m \cos \frac{2\pi m i}{J},$$

$$\hat{Q}_i^s = \frac{2}{J} \sum_{m=1}^J Q_m \sin \frac{2\pi m i}{J}, \quad i = 1, 2, \dots, J_2 \text{ (odd } J)$$

For even  $J$ , the coefficient  $\hat{Q}_{J/2}^c$  can be written in the form:

$$\hat{Q}_{J/2}^c = \frac{1}{J} \sum_{m=1}^J Q_m \cos \pi m = \frac{1}{J} \sum_{m=1}^J Q_m (-1)^m$$

The coefficient  $\hat{Q}_{J/2}^s$  is equal to zero as  $\sin \pi m = 0$ .

Hence the Fourier procedure for solving Poisson's equation consists of two steps. At first, we perform the Fourier analysis



- Card 6 (E10.3) :  $DL = \frac{\Delta}{\lambda_D}$ , where  $\Delta$  is the space step and  $\lambda_D$  is the Debye length.
- Card 7 (E10.3) :  $DT = \Delta t \cdot \omega_{pe}$ , where  $\Delta t$  is the time step and  $\omega_{pe}$  is the electron plasma frequency.
- Card 8 (E10.3) : DEN = initial plasma density in  $\text{cm}^{-3}$ .
- Card 9 (E10.3) : E0 = initial plasma temperature in eV.
- Card 10 (E10.3) : R0 = radius of the plasma column in cm.
- Card 11 (E10.3) : UD =  $v_d/v_{th}$ , drift velocity of the initial  $v$ -distribution in thermal velocity units (for the case of two-stream instability).
- Card 12 (5E10.3): XMIN ) = boundary values of X for the phase  
XMAX ) space plot.  
HX = value of the step in X units between  
two lines of the phase space plot.  
UMIN ) = boundary values of  $v/v_{th}$  for the  
UMAX ) phase space plot and for the  
velocity distribution plot.
- Card 13 (6E10.3): AL ) boundary conditions parameters  $A_L$ ,  
BL )  $B_L, C_L, A_R, B_R, C_R$  (see eqs. (12),  
CL ) = (13) for the Poisson equation. If  
AR ) all these parameters are equal to  
BR ) zero the Poisson equation will be  
CR ) solved with periodical boundary  
conditions.
- Card 14 (6E10.3): PE(1) )  
- ) amplitudes and space parameters  
- ) = of the external potential (see  
- ) sec. 5.9).  
PE(6) )
- Card 15 (5E10.3): PE(7) )  
- ) time parameters of the external  
- ) = potential (see sec. 5.9).  
- )  
PE(11) )

If the parameter  $NV > 1$ , pairs of cards similar to cards 14 and 15 have to be added with parameters for other cases of external potential. The total number of such pairs is equal to  $NV$ .

### 5.3 Program Output

After every  $KD$  time step, the total energy and the electrostatic energy (in dimensionless units, see sec. 3) are printed. If  $XMAX > XMIN$  the phase-space plot and the velocity distribution function are drawn by the line-printer. When  $XMAX \leq XMIN$  these output data will be omitted.

After the final time step,  $KF$ , for every case of the external potential, a plot of the temporary development of the potential space variation is drawn by the plotter. The potential is plotted in  $A_p W_{ph}/e$  units, where  $A_p = |\eta(x)|_{\max}$  (see eq. (5)). The test run output for nonperiodical boundary conditions is shown in Figs. 2a, 2b.

### 5.4 Work Constants

- $N$  = total number of simulation particles.
- $G$  =  $2\omega_{pe}^2 (\Delta t)^2 = 2DT^2$  = normalization factor for the charge density calculation (see eqs. (6)).
- $LD$  = the Debye length  $\lambda_D$  in cm.
- $CF\emptyset$  =  $W_{ph} \frac{2(\Delta t)^2}{m_e \Delta^2} = 4 \cdot \left(\frac{DT}{DL} \frac{R\emptyset}{LD}\right)^2 \frac{1}{23.116}$  = normalization factor for the external potential calculation (see eqs. (5), (6)).
- $VN$  =  $V \cdot v_{th}/v = 2\sqrt{2} \frac{DT}{DL}$  = velocity normalization factor for the initial distribution generation, where  $v_{th}$  is the electron thermal velocity.
- $UTN$  =  $2GJ$  ) factor for the electric field energy
- $DKE$  =  $4k_{\perp}^2 \Delta^2$  ) (8) calculation.
- $SF$  =  $\max(|AP1|, |AP2|) \cdot CF\emptyset$  = scale factor for plotting the potential space variations, where  $AP1, AP2$  are the external potential amplitudes (see sec. 5.9).
- $KBOUND$  = a sign for the type of boundary condition.  $KBOUND = 1$  for periodical boundary conditions and  $KBOUND = 0$  for reflecting boundary conditions.
- $JEV$  = 1 for even  $J$  and  $\emptyset$  for odd  $J$ .

$J_2$  =  $J_2$  = maximum index for the potential values Fourier synthesis (14).

$C(1000)$  =  $\cos\frac{2\pi j}{J}$  constants for the fast Fourier transform

$S(1000)$  =  $\sin\frac{2\pi j}{J}$  ) = in a case of periodical conditions.

### 5.5 Principle Variables

$X(50000)$  ) arrays of the normalized coordinates and velocities of all particles in each time step.

$Q(1000)$  ) = arrays of the normalized charge density and potential in the space grid points  $X = 1, 2, \dots, J$ .

$F0$  = value of the potential in the point  $X = 0$ .

$A(500)$  ) = arrays of the cosine and sine Fourier coefficients (15) for the Fourier method of solving the Poisson equation.

$K$  = number of the time step in the leap-frog scheme (7).

$NEF$  = number of cases with the concrete external potential parameters.

$KTP$  = output interval counter (the program output is realized when  $KTP = KD$ ).

$T$  =  $DT \cdot (2K-1)/2\pi$  = time moment in plasma periods, which corresponds to the time step  $K$ .

$UT$  = normalized electric field energy per one particle (7).

$ET$  = normalized thermal plasma energy per one particle.

### 5.6 Subroutine PNSOLV

**Purpose:** This subroutine solves the finite difference Poisson equation (10) for the dimensionless charge density values  $Q(X)$  in the points  $X = 1, 2, \dots, J$ . The solution is presented by the values  $F(X)$  in the same points and the value  $F0$  in the point  $X = 0$ . If the parameter  $KBOUND = 1$  (all parameters in card 13 are equal to zero), the Poisson equation is solved for periodical boundary conditions, and when  $KBOUND = 0$ ,

it is solved for the general nonperiodical conditions (12), (13) (see sec. 4). The subroutine obtains all parameters by means of the COMMON-block |PSN|.

Calling Sequence: CALL PNSOLV.

COMMON-block |PSN|: KBOUND, JEV (see sec. 5.4), AL, BL, CL, AR, BR, CR, J (see eqs. (12), (13) and sec. 5.2), J2, DKT, F0, F(1000), Q(1000), C(1000), S(1000), A(500), B(500) (see secs. 5.4, 5.5).

### 5.7 Subroutine INGEN

**Purpose:** This subroutine generates the initial values of the dimensionless coordinates X and velocities V for all simulation particles. The initial distribution is homogeneous in X-space and random in V-space, according to the Maxwellian distribution. For generation of the V-distribution, the Risc Computer Library subroutine NORMAR is used. When the parameter UD > 0, the subroutine creates a two-Maxwellian distribution in V-space with relative drift velocity  $UD \cdot v_{th}$ .

Calling Sequence: CALL INGEN (X, V, N, J, VN, UD).

**Parameters:** X(N) ) = arrays of the dimensionless coordinates  
V(N) ) and velocity values generated by the  
subroutine.

N = number of points in phase space generated by the subroutine.

J = length of X-space interval.

VN = velocity normalization factor (see sec. 5.4).

UD = relative drift velocity in  $v_{th}$  units of a two-Maxwellian distribution for the two-stream instability case.

### 5.8 Subroutine PCTR XV

**Purpose:** By means of the line-printer the subroutine plots a phase space distribution of the simulation particles and also a velocity distribution function by a simple histogram summing.

Calling Sequence: CALL PCTRXV (X, V, XMIN, XMAX, HX, UMIN, UMAX, VN).

Parameters: X(N) ) = coordinates of the simulation particles  
V(N) ) in phase space.  
N = total number of particles.  
XMIN )  
XMAX ) parameters of plots printed  
HX ) = by the subroutine (see the  
UMIN ) description of card 13 in sec. 5.2).  
UMAX )  
VN = velocity normalization factor (see sec. 5.4).

### 5.9 Function EXTPOT

Purpose: This function calculates the value of the external potential in a point X at a time moment T for studying the nonlinear plasma response to a short localized excitation. The external potential is normalized by the kinetic energy of an electron with velocity  $v_{ph}$ :

$$EXTPOT (X,T) = e\phi_{ext}/W_{ph} = \phi_{ext} \cdot \frac{2e}{m_e} \left(\frac{k_{\perp}}{\omega_{pe}}\right)^2 .$$

The function EXTPOT is constructed as a product of two functions:

$$EXTPOT (X,T) = \eta(X) \cdot \sigma(T) ,$$

where the time-dependent function  $\sigma(T)$  is assumed to satisfy the requirements:

$$\sigma(T) \geq 0 , \quad \max \sigma(T) = 1 .$$

According to the conditions of the experiments with the Q-machine plasma<sup>1), 2)</sup>, the following approximation for the impulse profile is used (see Fig. 3):



$$\eta(X) = \begin{cases} -AP1, & 0 \leq X \leq DX1, \\ 0, & 0 \leq X - DX1 - DP1 \leq DX2, \\ -AP2, & X - DX1 - DP1 \geq DX2 + DP2, \end{cases} \quad (16)$$

$$\sigma(T) = \begin{cases} 0, & T \leq DT1, \\ 1, & 0 \leq T - DT1 - DT2 \leq DT3, \\ STF, & T - DT1 - DT2 \geq DT3 + DT4, \end{cases}$$

where AP1, AP2 are the amplitudes of the external potential jumps in  $W_{ph}/e$  units (see the expression (5)).

For the approximation of the space and time shapes of the external potential in the intervals between the regions mentioned in the definitions of  $\eta(X)$  and  $\sigma(T)$ , half a cosine period is used with corresponding parameters to satisfy the relations (16).

External potential parameters are taken from the COMMON-block |EXTP|.

Calling Sequence: EXTPOT (X,T).

Arguments: X = space point coordinate in  $\Delta$  units,  
T = time moment in plasma periods.

COMMON-block |EXTP|: PE(1) = AP1 )  
PE(2) = DX1/J ) parameters of the  
PE(3) = DP1· $\Delta$ /R $\emptyset$  ) = space profile func-  
PE(4) = AP2 ) tion  $\eta(X)$  (see Fig.3).  
PE(5) = DX2/J )  
PE(6) = DP2· $\Delta$ /R $\emptyset$  )  
  
PE(7) = DT1 )  
PE(8) = DT2 ) parameters of the  
PE(9) = DT3 ) = time profile func-  
PE(10) = DT4 ) tion  $\sigma(T)$  (see Fig.3).  
PE(11) = STF )

## 6. Acknowledgements

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8. Figures

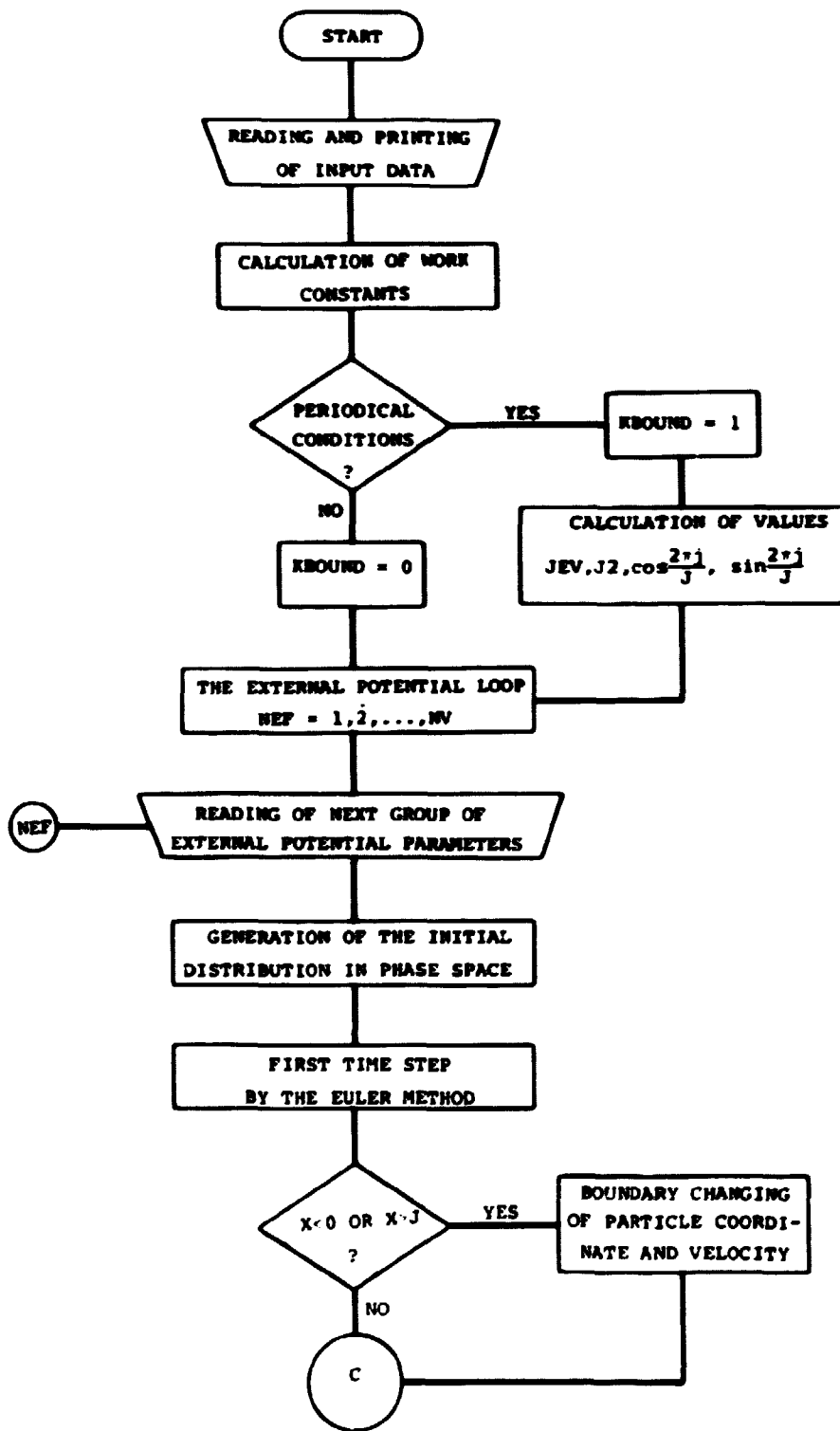


Fig. 1a. General Block Diagram of the Code.

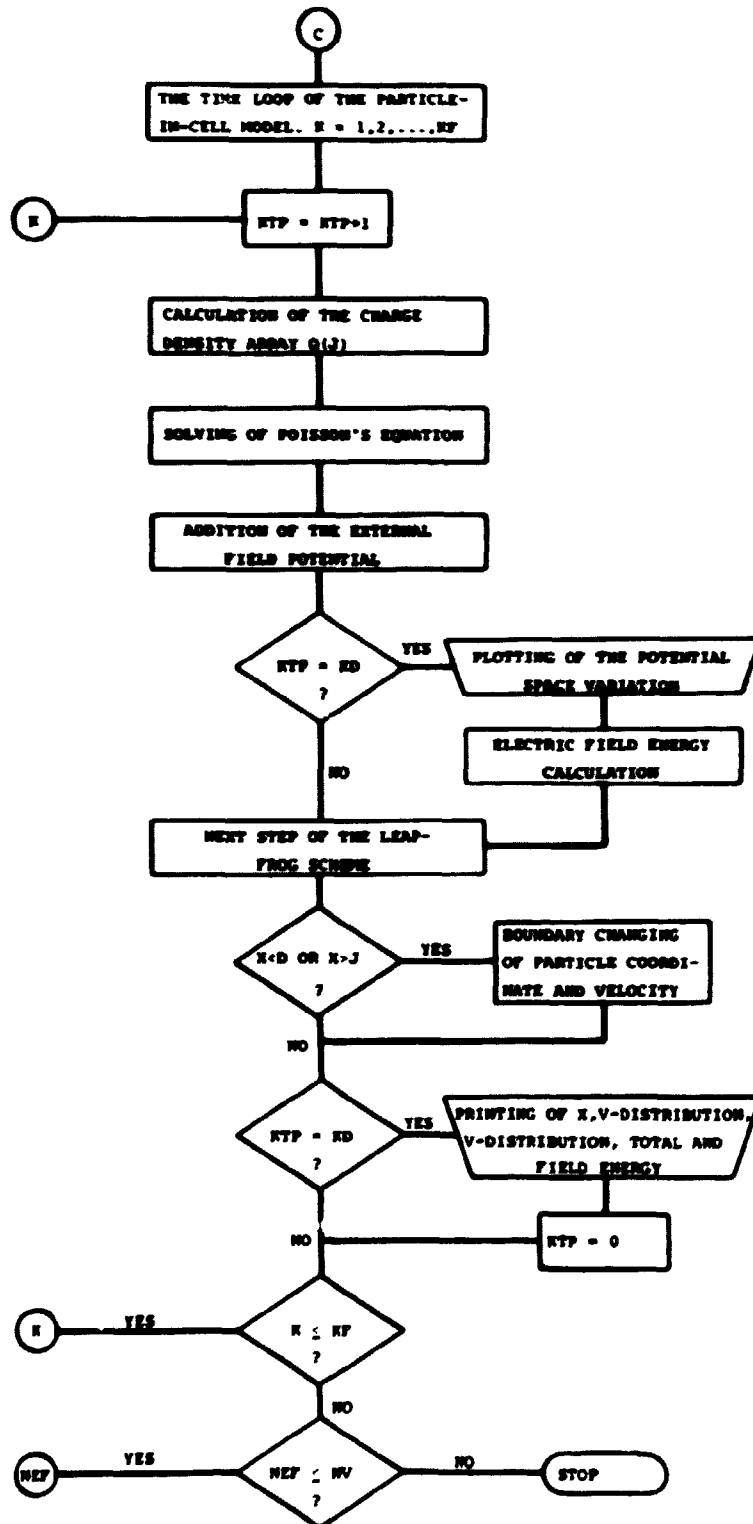


Fig. 1b. General Block Diagram of the Code (cont.)

```
SIMULATION PARAMETERS
NO. OF CELLS = 100  NO. OF PARTICLES = 1000  TIME = 1.00E+01  DT = .0001E+01
CELL SIZE = .001E+01  IONIZATION = .001E+01  ALPHA = .001E+01

PRIMARY SPECTRA PARAMETERS
PHOTON ENERGY = 1.00E+01  ELECTRON ENERGY = 1.00E+01  IONIZATION ENERGY = 1.00E+01

SECONDARY SPECTRA PARAMETERS
PHOTON ENERGY = 1.00E+01  ELECTRON ENERGY = 1.00E+01  IONIZATION ENERGY = 1.00E+01

CASE 1 - IONIZATION PARAMETERS
PHOTON ENERGY = 1.00E+01  ELECTRON ENERGY = 1.00E+01  IONIZATION ENERGY = 1.00E+01
DT = .0001E+01  IONIZATION = .001E+01  ALPHA = .001E+01  IONIZATION ENERGY = 1.00E+01
```

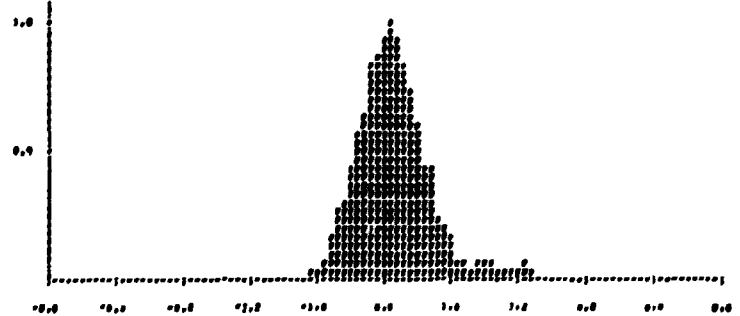
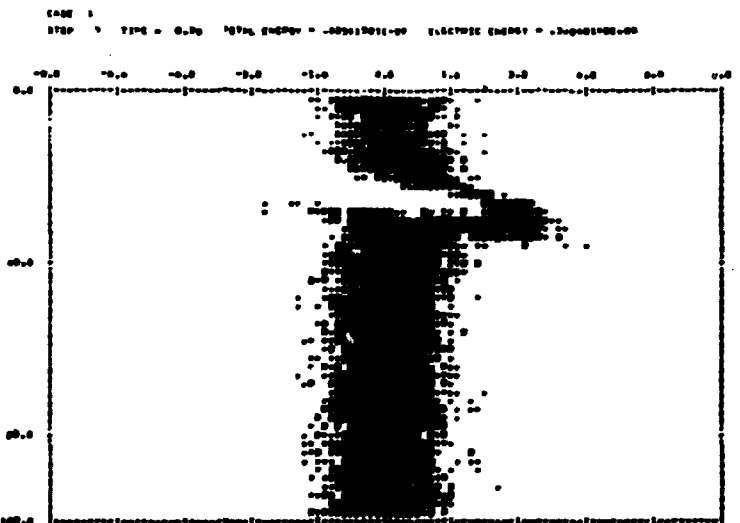


Fig. 2a. Test Run Output by the Line-printer.

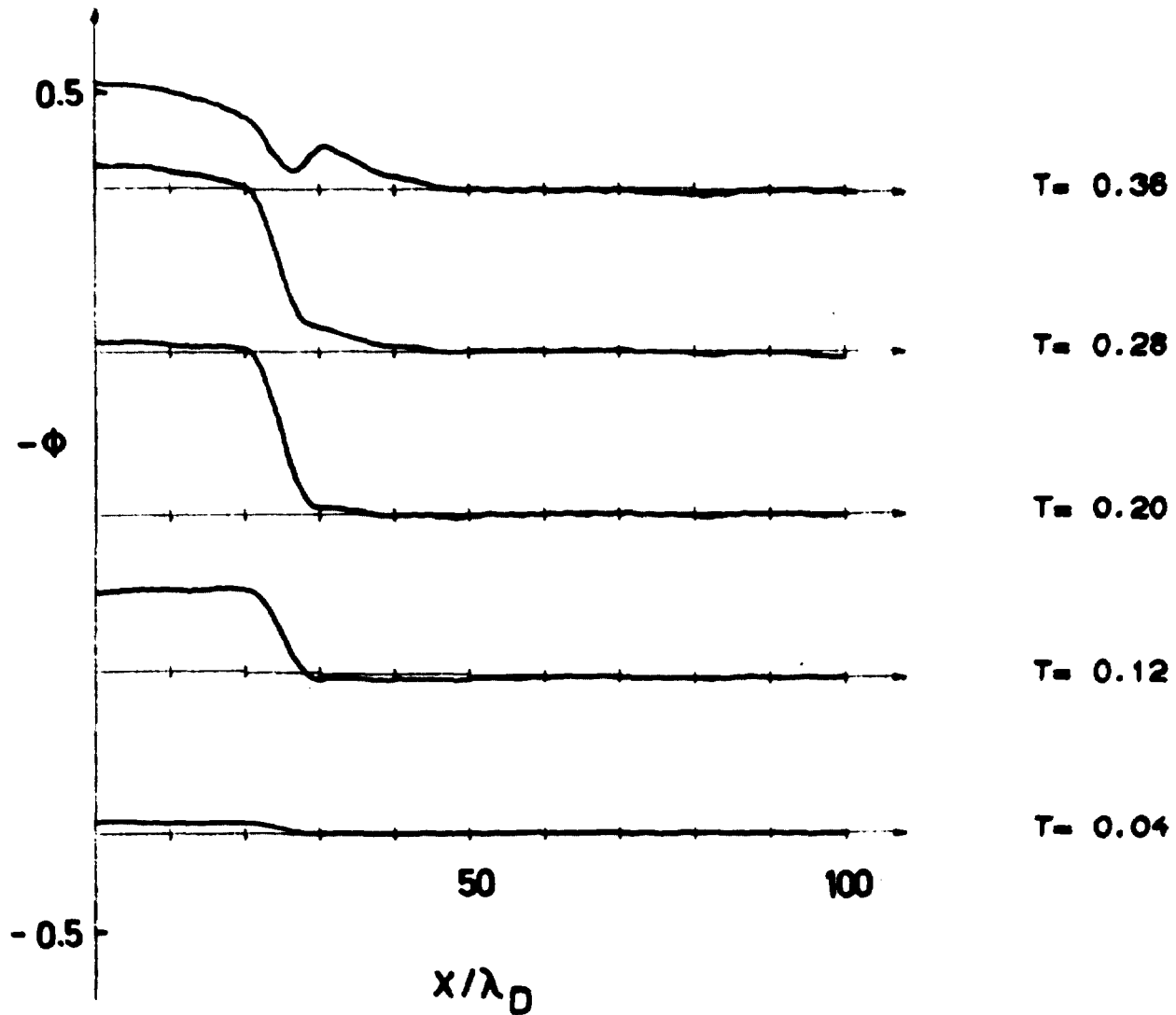


Fig. 2b. Test Run Output by the Plotter. Evolution of the Potential Space Variation.

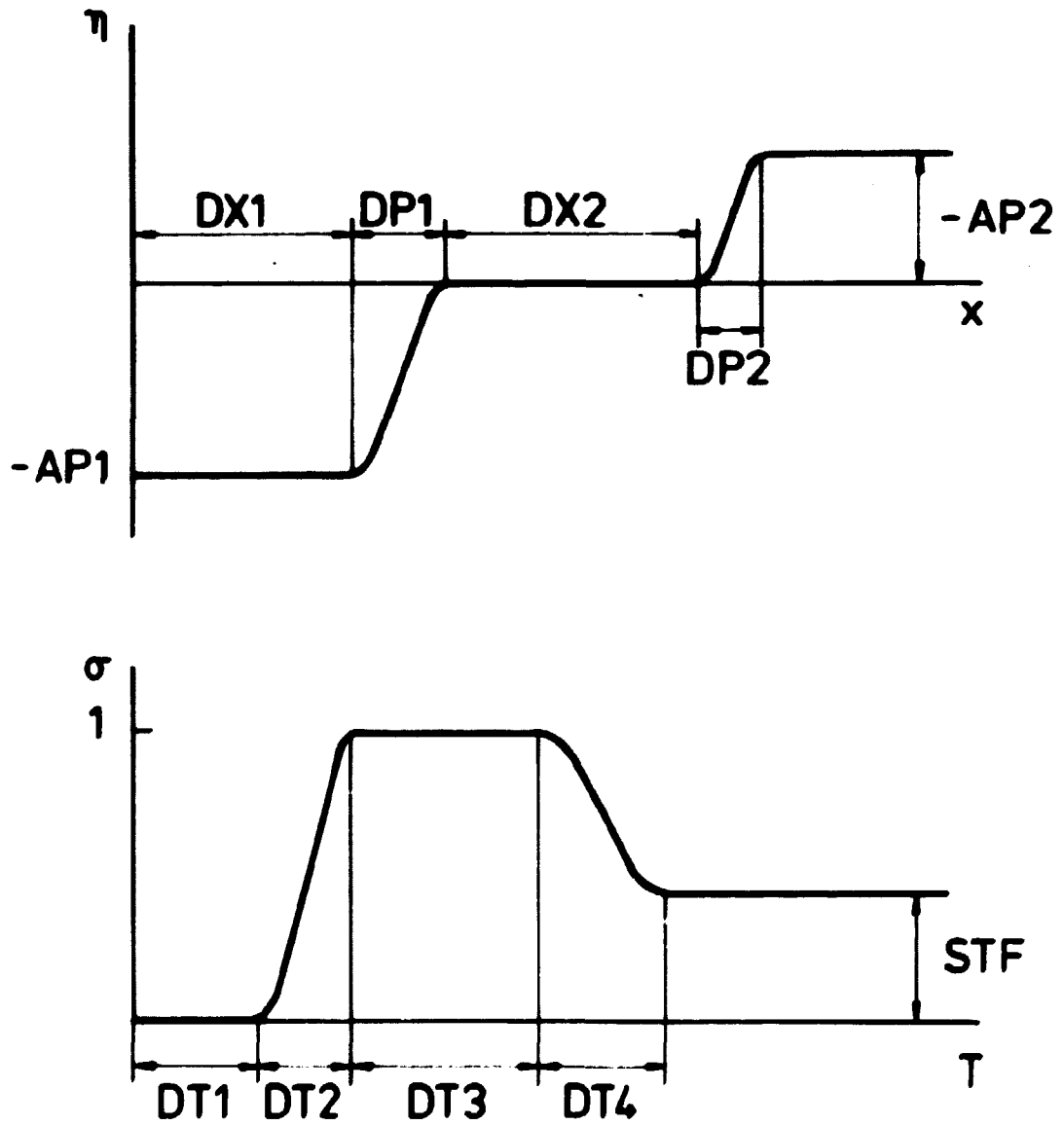


Fig. 3. Space and Time Profiles of the External Potential Function  $EXTPOT(X,T) = \eta(X) \cdot \sigma(T)$ .

9. Program Listings

1. MAIN PROGRAM SIMULA

C MAIN PROGRAM SIMULA

```

REAL PLTEXT(1),LC
DATA PLTEXT/'T'/
DIMENSION X(50000),V(50000)
COMMON /EXIT/PE(11)/PSA/XM(LNU),EV,AL,PL,CL,AR,BN,CH,J2,CN1,
      *FC,F(1000),C(1000),C(1000),S(1000),A(500),R(500)

```

C HEADINGS AND PRINTING OF THE INPUT DATA

```

C
      READ(5,1) LC,J,KF,KD,NV
1  FORMAT(15)
      READ(5,2) DL,DT,DEH,EO,RO,LU,XPIN,XPAX,HP,LPIN,LPAX,
      *AL,BL,CL,AR,BN,CH
2  FORMAT(6(E10,3)/),5E10,3/6E10,3)
      WRITE(6,3) LC,J,KF,KD,NV,DL,DT,DEH,EO,RO,LU
3  FORMAT(///10X,'SIMULATION PARAMETERS: '//
      *'LC=',10,4X,'J=',10,4X,'KF=',10,4X,'KD=',10,4X,'NV=',10,4X,
      *'DL=',E9,3,4X,'DT=',E9,3,4X,'DEH=',E9,3,4X,'EO=',E9,3,4X,
      *'RO=',E9,3,4X,'LU=',E9,3,4X)
      WRITE(6,4) XPIN,XPAX,HP,LPIN,LPAX
4  FORMAT(///10X,'PHASE SPACE PLOT PARAMETERS: '//
      *'1X,'XMIN=',E9,3,4X,'XPAX=',E9,3,4X,'PX=',E9,3,4X,
      *'UPIN=',E9,3,4X,'LPAX=',E9,3,4X)
      WRITE(6,7) AL,PL,CL,AR,BN,CH
7  FORMAT(///10X,'BOUNDARY CONDITIONS PARAMETERS: '//
      *'1X,'AL=',E9,3,4X,'BL=',E9,3,4X,'CL=',E9,3,4X,
      *'AR=',E9,3,4X,'BN=',E9,3,4X,'CH=',E9,3,4X)

```

C CALCULATION OF THE WORK CONSTANTS

```

C
      N=NC*J
      G=2.*DT*DT
      LD=740.*SQRT(EO/DEH)
      CFO=4.0*(DT*RO/DL/LD)**2/23.110
      VM=2.*SQRT(2.)*DT/DL
      UTM=J*4.*DT*DT
      DRT=1.+(2.400/RO*LD*CL)**2/2.
      DRE=0.*(DRT-1.)
      PI=3.141592654
      PI2=6.283185307
      CALL PSCALE(1200,320)

```

C BOUNDARY CONDITIONS ANALYSIS

```

C
      IF(AL.EQ.0.AND.BL.EQ.0.AND.AR.EQ.0.AND.BN.EQ.0)KBCOND=1
      IF(KBCOND.EQ.0) GO TO 4

```

C PREPARATION OF THE WORK CONSTANTS AND ARRAYS FOR PERIODICAL BOUNDARY CONDITIONS

```

C
      C(1)=J/2.
      J2=J/2
      IF(C(1)=J2-0.1) 51,51,52
51  JEV=1
      GO TO 53
52  JEV=0
      J2=(J-1)/2

53  C(1)=PI2/J
      S(1)=SIN(C(1))
      C(1)=COS(C(1))
      DO 5 K=2,J
      C(K)=C(K-1)*C(1)+S(K-1)*S(1)
      S(K)=S(K-1)*C(1)+C(K-1)*S(1)

```





```

C
C SERIAL ELECTRIC FIELD ENERGY CALCULATION
C
47 LT=0.
   GN 20 I=1.
   YI=1
   IF(I.EQ.J) IJ=J-1
   IF(I1=1) A6=05.00
45 F1=U0*G
   GC TU 47
46 F1=F(I1-1)*F(I1-1)
47 F2=F(I1)*F(I1)
   F3=F(I1+1)*F(I1+1)
   AF=0.5*(F1-2.*F2+F3)
   BF=F3-F2-AF*(2.*I1+1)
   CF=F2-AF*I1-AF*I1+1
20 L1=1+AF/3.*(1+1)-(I1-1)*3)*BF/2.*(1+2-(I1-1)*2)*CF
   L1=CKEOUT
C
C CALCULATION OF THE COEFFICIENTS FOR THE ELECTRIC FIELD INTERPOLATION
C
49 EL=2.*(FU-2.*F(1)+F(2))
   Q(1)=2.*(F(2)-F(1))-EL*0.5
   GC 20 I=2.
   EN=2.*(F(I-1)-2.*F(I)+F(I+1))
   F(I-1)=EL
20 Q(I)=2.*(F(I+1)-F(I))-EN*(2.*I+1)/2.
   EL=EN
   F(J-1)=EN
   F(J)=F(J-1)
   Q(J)=Q(J-1)
   IF(NTP=NS) 22.21.22
C
C TOTAL ELECTRIC FIELD ENERGY CALCULATION
C
21 GL 10 I=1.
   GN=1
   EL=1-1
   UT=UT+(I)*Q(I)*GN-EL)
   EN=ER*GN
   EL=EL*EL
   LT=LT+F(I)*Q(I)*(GN-EL)
   ER=EN*I
   EL=EL*(I-1)
10 LT=UT+F(I)*Q(I)/3.*(FR-EL)
   UT=LT/UTN
   ET=0.
C
C NEXT STEP OF THE LEAP-FRAG SCHEME
C
22 GU 20 N=1.
   IX(N)=1
   IF(NTP.EQ.NC) VNC=V(N)
   V(N)=V(N)+F(I)*X(N)+Q(I)
   IF(NTP.EQ.NC) ET=ET+0.25*(VX2+V(N))*0.2
   X(N)=X(N)+V(N)
   IF(X(N).GT.G.O.AND.G.X(N).LT.0) GU TC 20
   X(N)=X(N)-SIGN(J,X(N))
   IF(NBC=NC) A2.44.42
40 X(N)=X(N)
   V(N)=V(N)
42 IF(X(N).EQ.G.O.CF.X(N).EQ.0) X(N)=ABS(X(N)-1.E-10)
20 CONTINUE
C
C OUTPUT OF INFORMATION ABOUT PLASMA STATE
C
   IF(NTP.LT.NC) GC TC 30
   WRITE(A,LP) NTP,ET/UT,LT
17 FORMAT(//////I6.0, 'CASE ',I2//I9.0, 'STEP ',I3.3X, 'TIME ',F6.2.3X,
   ' TOTAL ENERGY ',E14.8.3X, 'ELECTRIC ENERGY ',E14.8//)
   IF(XMIN.LT.XMAX)
   CALL PCTRXV(X,V,IX,ITN,XPAR,PR,UPIN,UP,X,VA)
   NTP=0
30 CONTINUE
   CALL PTEMP
   STOP
   END

```

## 2. SUBROUTINE PNSOLV

```

SUBROUTINE PNSOLV
  COMMON /PSN/ KBOUND,JEV,AL,BL,CL,A1,HN,CM,J,J2,DKT,FG,F(1000),
  *G(1000),C(1000),S(1000),A(500),B(500)
  IF(KBOUND.GT.0) GO TO 3
  CCCC
  NONPERIODICAL BOUNDARY CONDITIONS (SCANNING METHOD).
  CCCC
  DOWNWARDS SCANNING
  CM=A1+HN
  C(J)=A1*(2.-DKT)/CM
  S(J)=(CR-0.5*AR+G(J-1))/CM
  DO 1 I=1,J-1
  N=J-I+1
  CM=C(N)*2.-DKT
  C(N-1)=-1./CM
  1 S(N-1)=(G(N-1)-S(N))/CM
  CCCC
  UPWARDS SCANNING
  FU=0.
  IF(DKT.GT.1.) FU=(CL+AL*(0.5*(C(2)*S(1)+S(2))-2.*S(1)))/
  *(AL*(2.*C(1)+0.5*C(1)*C(2))-1.5*AL)
  F(1)=C(1)+FU+S(1)
  DO 2 I=1,J-1
  2 F(I+1)=F(I)+C(I+1)+S(I+1)
  GO TO 8
  CCCC
  PERIODICAL BOUNDARY CONDITIONS (FOURIER METHOD).
  CCCC
  FOURIER ANALYSIS
  3 DO 5 I=1,J2
  A(I)=0.
  H(I)=0.
  C1=1.
  S1=0.
  DO 4 M=1,J
  CM=C1+C(I)-S1+S(I)
  S1=S1+C(I)+C1*S(I)
  C1=CM
  A(I)=A(I)+G(M)*C1
  4 B(I)=B(I)+G(M)*S1
  CM=J*(C(I)-DKT)
  A(I)=A(I)/CM
  5 B(I)=B(I)/CM
  IF(JEV.GT.0) A(J2)=A(J2)/2.
  CCCC
  FOURIER SYNTHESIS
  DO 6 I=1,J
  F(I)=0.
  C1=1.
  S1=0.
  DO 6 M=1,J2
  CM=C1+C(I)-S1+S(I)
  S1=S1+C(I)+C1*S(I)
  C1=CM
  6 F(I)=F(I)+A(M)*C1+B(M)*S1

  DO 7 I=1,J
  7 F(I)=F(I)-F(J)
  FU=0.
  8 RETURN
  END

```

### 3. SUBROUTINE INGEN

```

SUBROUTINE INGEN(X,V,N,J,VN,UC)
DIMENSION X(N),V(N)
DX=(1.+J)/N
X(1)=DX/2
DO 1 I=2,N
1 X(I)=X(I-1)+DX
Y=174
IV=8479
DO 2 I=1,N
CALL NCHVAR(IV,I)
2 V(I)=VN*(1/1.414214+SIGN(UC,RANDCM(Y)*0.5))
RETURN
END

```

### 4. SUBROUTINE PCTR XV

```

SUBROUTINE PCTR XV(X,V,N,XMIN,XMAX,HX,LMIN,LPAX,VN)
DIMENSION X(N),V(N),STRING(101)
DIMENSION MESH(250,100)
REAL BL/' ','SN1/' ','SN2/' ','SN3/' ','SN4/' ','SN5/' ','
REAL SN6/' ','SN7/' ','SN8/' ','
N1=1
N2=3
IXMAX=(XMAX-XMIN)/HX+0.1
DO 25 I=1,IXMAX
DO 25 K=1,100
25 MESH(I,K)=0
HV=(LPAX-UMIN)*VN/100.
VMIN=UMIN*VN
DO 1 P=1,N
IX=(X(P)-XMIN)/HX+0.5
IV=(V(P)-VMIN)/HV+0.5
IF((IX.LE.0.OR.IX.GE.IXMAX).OR.(IV.LE.0.OR.IV.GE.100))
*GO TO 1
MESH(IX,IV)=MESH(IX,IV)+1
1 CONTINUE
WRITE(6,2)(LMIN+(I-1)*10*HV/VN,I=1,11)
2 FORMAT(/,14X,F4.1,10(6X,F4.1)/)
WRITE(6,3)
3 FORMAT(16X,10('I-----'),',',)
WRITE(6,14) XMIN
MARK=20
DO 4 IX=1,IXMAX=1
DO 5 K=2,100
IF(MESH(IX,K=1)) 7,6,7
6 STRING(K)=BL
GO TO 5
7 IF(MESH(IX,K=1)=N1) A,A,9
8 STRING(K)=SN1
GO TO 5

```

```

9 IF(MESH(IX,K-1)-N2) 10,10,11
10 STRING(K)=SA2
   GO TO 5
11 STRING(K)=SA3
5 CONTINUE
   BS=SN4
   IF(IX.EQ.PARK) BS=SN5
   STRING(1)=BS
   STRING(101)=BS
17 FORMAT(16X,101A1)
   WRITE(6,12) STRING
   DO 22 K=1,101
   BS=BL
   IF(STRING(K).EQ.SA2) BS=SA4
   IF(STRING(K).EQ.SA3) BS=SA7
22 STRING(K)=BS
   WRITE(6,23) STRING
23 FORMAT(' ',15X,101A1)
   DO 24 K=1,101
   BS=BL
   IF(STRING(K).EQ.SA7) BS=SA9
24 STRING(K)=BS
   WRITE(6,23) STRING
   IF(MARK=IX) 13,13,4

13 WRITE(6,14) X*IN+IX*MX
14 FORMAT(' ',7X,F6.1)
   MARK=MARK+20
4 CONTINUE
   WRITE(6,3)
   WRITE(6,14) XMAX
   DO 16 IV=1,100
   DO 16 IX=2,IXMAX+1
16 MESH(1,IV)=MESH(1,IV)+MESH(IX,IV)
   NMAX=0
   DO 17 IV=1,100
   IF(MESH(1,IV).GT.NMAX)NMAX=MESH(1,IV)
17 CONTINUE
   BS=SN4
   WRITE(6,19) BS,BS
19 FORMAT(' ',7(/),16X,A1/16X,A1)
   DO 21 IX=1,30
   STRING(1)=SA4
   IF(IX.EQ.1.OR.IX.EQ.16) STRING(1)=SN5
   DO 20 K=2,100
   STRING(K)=BL
   IF(MESH(1,K-1).GE.(31-IX)*NMAX*0.03333) STRING(K)=SNJ
20 CONTINUE
   WRITE(6,12) STRING
   IF(IX.EQ.1.OR.IX.EQ.16) WRITE(6,14)(31-IX)/30.
21 CONTINUE
   WRITE(6,3)
   WRITE(6,2)(LMIN+(I-1)+10+MV/VN,I=1,11)
   RETURN
   END

```

## 5. FUNCTION EXTPOT

```
FUNCTION EXTPOT(X,T)
COMMON /EXTP/ PE(11)
PI=3.141592654
T1=T-PE(7)
IF(T1.LE.0.0) GO TO 10
IF(T1-PE(8)) 3,3,4
3 AT=0.5*(1-COS(PI*(T1/PE(9))))
GO TO 9
4 T1=T1-PE(8)
IF(T1-PE(9)) 5,5,6
5 AT=1.0
GO TO 4
6 T1=T1-PE(9)
IF(T1-PE(10)) 7,7,8
7 AT=0.5*(1-PE(11))*COS(PI*T1/PE(10))+1.0*PE(11)
GO TO 9
8 AT=PE(11)
9 IF(ABS(AT).LT.1.E-7) GO TO 10
X1=X-PE(2)
IF(X1) 10,10,11
10 EXTPOT=-PE(1)*AT
GO TO 14
11 IF(X1-PE(3)) 12,12,13
12 IF(PE(1).EQ.0) GO TO 10
EXTPOT=-PE(1)*0.5*AT*(COS(PI*X1/PE(3))+1.0)
GO TO 14
13 X1=X1-PE(3)
IF(X1-PE(5)) 10,10,15
10 EXTPOT=0.
GO TO 14
15 X1=X1-PE(5)
IF(X1-PE(6)) 16,16,17
16 IF(PE(4).EQ.0) GO TO 10
EXTPOT=PE(4)*0.5*AT*(COS(PI*X1/PE(6))-1.0)
GO TO 14
17 EXTPOT=-PE(4)*AT
10 RETURN
END
```